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**A comparative study of methanol adsorption and dissociation over WO<sub>3</sub>(001)  
and ReO<sub>3</sub>(001)**

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## Abstract

Tungsten ( $5d^46s^2$ ) and rhenium ( $5d^56s^2$ ) form  $MO_3$  oxides ( $M = W$  or  $Re$ ) with similar structures. The adsorption and dissociation of methanol on these oxide surfaces, often used to probe the surface redox centers, have been analyzed using periodic density functional calculations. Molecular adsorption of methanol at the metal site on both surfaces with 0.5 ML oxygen coverage was found to be exothermic with adsorption energies of -74 and -106 kJ/mol on  $WO_3(001)$  and  $ReO_3(001)$ , respectively. In contrast, heterolytic dissociation of methanol to adsorbed methoxide species at the metal site and H at the surface oxygen site is not energetically favored on  $WO_3(001)$  but favored on  $ReO_3(001)$ . The dissociation energies to form coadsorbed methoxide and hydrogen adatom are 35 kJ/mol on  $WO_3$  and -112 kJ/mol on  $ReO_3$ , respectively. The activation barrier for dissociating the molecularly adsorbed methanol on  $ReO_3(001)$  was determined to be 19 kJ/mol. Dehydrogenation to form coadsorbed hydroxymethyl and hydrogen adatom is not energetically favorable on both surfaces with respect to the molecularly adsorbed methanol. However, the dehydrogenation path is exothermic with respect to the gas phase methanol on  $ReO_3$ , with the heats of reaction of -25 kJ/mol, but highly endothermic on  $WO_3$ , with the heats of reaction of 114 kJ/mol.

**Keywords:** Transition metal oxides; Methanol; Dissociation; Dehydrogenation; DFT

## 1. Introduction

Transition metal oxides are widely used as both catalysts and supports of catalyst.[1] For example, supported  $\text{WO}_x$  catalysts have been used in hydrocarbon cracking,[2] isomerization of alkenes and alkanes,[3-6] olefin metathesis,[7] and selective catalytic reduction of  $\text{NO}_x$ .[8] Supported rhenium oxides were found to catalyze the oxidation of methanol to dimethoxymethane and selective oxidation of propane.[9, 10] The performance of these oxides as either catalysts or supports of catalyst is strongly dependent on their intrinsic properties and structure of the exposing surfaces. Understanding the nature of transition metal oxides and the molecular interactions on the surfaces of oxides is therefore important to optimize the reactivity and selectivity of existing catalysts and to design new catalytic systems with improved performance.

Methanol is a commodity and has been proposed as a primary target molecule for  $\text{CO}_2$  hydrogenation as well as  $\text{CO}_2$  electrochemical reduction.[11, 12] It can be selectively oxidized over many transition metal oxides to products such as formaldehyde and methyl formate which are important primary reagents for organic synthesis.[13] The product distribution of the selective oxidation reaction depends on the nature of oxides. Methanol can also be used as an alternative fuel or an on-board carrier of hydrogen for hydrogen fuel cell powered vehicles.[14] In addition, methanol has been proposed as a probe molecule to analyze the number as well as the nature of surface active sites on many transition metal oxide catalysts.[15] Wachs and coworkers have suggested that the oxygen bridging the support metal atom to the  $\text{V}=\text{O}$  center over supported vanadia catalysts plays an important role in the selective oxidation of methanol based on the dependence of turnover frequency on the nature of the support.[16-18] These authors proposed the following simple mechanism:



Reaction (R1) represents protonation of the  $\text{V}-\text{O}-\text{M}$  bond followed by  $\text{M}-\text{O}-\text{V}$  bond scission and  $\text{CH}_3\text{O}^-$  addition to vanadium. The reactions that follow include methoxide dissociation (R2) and disproportionation reaction of surface hydroxyls (R3). A number of mechanistic models have been developed, including attempts to evoke the role of the electronegativity of the cations in the support and the modes in which various bridging O atoms are activated. This mechanism has

been expanded to explicitly incorporate oxygen vacancy formation and filling, and alteration of charge on the metal cations.

The structurally similar oxides of tungsten and rhenium offered an opportunity to probe the effect of the charges on the cations to their reactivity toward methanol. A difference by one 5d electron between rhenium ( $5d^56s^2$ ) and tungsten ( $5d^46s^2$ ) has significant consequences on the properties of their trioxides  $MO_3$  ( $M = W$  or  $Re$ ). For example,  $ReO_3$  is cubic with a perovskite-like structure,[19] but the structure of  $WO_3$  is distorted through tilting the  $WO_6$  octahedra and displacing the W atoms from the octahedral centers.[19, 20] The stoichiometric  $WO_3$  is a semiconductor with a band gap of 2.6 eV.[21]  $ReO_3$ , however, has one extra electron in the valence shell of each metal atom and the resulting band is only partially filled. Indeed, the conductivity of  $ReO_3$  reaches  $10^7 \Omega\text{cm}^{-1}$  at low temperatures.[22] Due to their unique electronic/redox properties, nanostructured tungsten and rhenium oxides have been found applications in gas sensing and catalysis.[23] We studied the chemistries of methanol on the surfaces of  $CeO_2$ ,[24, 25]  $Pt/TiO_2$ [26] and  $PdIn$ . [27] In the present paper we focus on whether differences in electronic and crystal structure of  $WO_3$  and  $ReO_3$  contribute to differences in chemical properties of the (001) surfaces of these trioxides. In particular we investigate adsorption and dissociation of methanol on these two surfaces.

## 2. Computational Details

Periodic DFT calculations reported herein were performed using the Vienna Ab initio Simulation Package (VASP).[28, 29] VASP is a plane wave code with ultrasoft pseudopotential[30] or projector augmented wave (PAW)[31, 32] to describe nuclei and core electrons. Gradient-corrected exchange-correlation functionals, PW91[33] and PBE[34] have been tested and the differences in the energetics and structures of using the two types of functionals were found to be negligible if the calculations were done self-consistently, i.e. using the PAW or pseudo-potentials generated with the same exchange-correlation functionals. As such, PAW potentials and PBE functionals were chosen in the present calculations. The 5d and 6s valence states of W and Re were treated explicitly in the calculations. The effect of treating the semi-core 5p state explicitly for W and Re was found negligible for the properties considered here. Spin-polarization was included in calculations when it became necessary. Test calculations using a cutoff energy up to 600 eV showed that the plane wave basis is converged to within 0.02

Å for relaxed bond distances and 0.05 eV for adsorption energies at 400 eV. A Gaussian type of electronic smearing with a width of 0.1 eV was used to improve convergence of electronic self-consistent cycles. For bulk and surface calculations, Monkhorst-Pack[35] scheme was used to generate k-points to sample the Brillouin zone. A space of less than  $0.05 \text{ \AA}^{-1}$  between two adjacent k-points in any directions was found to be appropriate for these oxides. In the studies of oxygen coverage effect, two-sided slabs with 6 layers of metal atoms were adopted. In the simulation of methanol adsorption and dissociation, we used a one-sided approach, i.e., placing methanol or the dissociative products on one side of the slab. The calculated adsorption energy of methanol on  $\text{WO}_3(001)$  is only different by 0.04 eV from that calculated with a two-sided model. The vacuum space between the adjacent slabs was kept larger than  $12 \text{ \AA}$ . The adsorbate atoms as well as the W (or Re) and O atoms in the top 3 triple layers were optimized according to the forces calculated quantum mechanically. The geometries were relaxed using a conjugate-gradient or quasi-Newton scheme as implemented in VASP. A geometry optimization was considered converged when the maximum force on the movable atoms reaches  $0.05 \text{ eV/\AA}$  or less. Transition state search was carried out by using the nudged elastic band (NEB) method and its improvements developed by Jonsson and co-workers.[36-38] Transition states determined with NEB were further optimized by quasi-Newton algorithm to minimize forces on the movable atoms.[39, 40] Normal mode harmonic vibrational analysis has been performed for the transition states and for selected adsorption geometries using the finite difference method as implemented in the latest version of VASP.[41]

### 3. Results and Discussion

Bulk  $\text{ReO}_3$  has a perovskite cubic structure with an experimental lattice constant of  $3.747 \text{ \AA}$  at 100 K and  $3.748 \text{ \AA}$  at 300 K.[42] On the other hand, bulk  $\text{WO}_3$  exhibits a series of polymorphs which can be considered as distortion of the ideal cubic  $\text{ReO}_3$  structure. The monoclinic unit cell of  $\epsilon\text{-WO}_3$  can be approximated by a reconstructed orthorhombic cell based on a cubic  $\text{ReO}_3$  cell scaled by  $(\sqrt{2} \times \sqrt{2} \times 2)$ . The low temperature  $\epsilon\text{-WO}_3$  phase has a lattice constant of  $a=5.278 \text{ \AA}$ ,  $b=5.156 \text{ \AA}$ ,  $c=7.664 \text{ \AA}$  and  $\beta=91.762^\circ$  as measured in neutron powder diffraction experiment.[43] Using these experimentally determined structures as input, we optimized the bulk structures of both  $\text{ReO}_3$  and  $\text{WO}_3$ .  $\text{ReO}_3$  remains a perfect cubic structure after relaxation and the equilibrium lattice constant is  $3.767 \text{ \AA}$ , in good agreement with the experimental results.[42] The optimized

lattice parameters of the  $\epsilon$ - $\text{WO}_3$  phase are  $a= 5.274 \text{ \AA}$ ,  $b= 5.130 \text{ \AA}$ ,  $c=7.638 \text{ \AA}$  and  $\beta=92.243^\circ$ , also in good agreement with the experiment.[43] The two W-O bonds to the same W atom along c-axis are  $1.777$  and  $2.149 \text{ \AA}$ , respectively. The relaxed bulk  $\text{WO}_3$  and  $\text{ReO}_3$  structures were shown in fig. 1 in octahedron representations. In order to compare with  $\text{WO}_3$ , we constructed a  $(\sqrt{2}\times\sqrt{2}\times 2)$  unit cell for  $\text{ReO}_3$  on the basis of the cubic structure. The distortion of the  $\epsilon$ - $\text{WO}_3$  structure from the perfect cubic structure as represented by  $\text{ReO}_3$  can be clearly identified by comparing the structures shown in fig. 1(a) for  $\text{WO}_3$  and fig. 1(b) for  $\text{ReO}_3$ . These relaxed bulk structures were subsequently used to create the slabs to simulate the surface as well methanol adsorption and dissociation.

There are different ways to build the (001) surface of  $\text{WO}_3$  and  $\text{ReO}_3$ . The surface may be terminated by either O or  $\text{WO}_2$  ( $\text{ReO}_2$ ). The surface may also be partly covered with oxygen. We analyzed different possibilities for the surfaces of both oxides and present the results for  $\text{WO}_3(001)$  and  $\text{ReO}_3(001)$  in turn in the following sections.

### 3.1 $\text{WO}_3$ (001)

We started from a slab with a  $c(2\times 2)$  surface unit cell, terminated by  $\text{WO}_2$  on both sides. The slab so-formed is non-polar and also non-stoichiometric. The surface with such a composition can only exist in a highly reducing environment. If we add an oxygen atom in the unit cell on each side of the slab, we can make the slab stoichiometric and keep it non-polar. Consequently, the W atoms of the half surface  $\text{WO}_2$  units will be covered by the oxygen atoms. The oxygen atoms form a  $c(2\times 2)$  super structure as shown by the blue box in figure 2(a). The formation of the  $c(2\times 2)$  oxygen super structure on  $\text{WO}_3(001)$  is highly exothermic. The integral heats of dissociative adsorption is  $-603.8 \text{ kJ/mol}$ . Increasing the oxygen coverage to a full monolayer, i.e., covering all the W atoms of the slab with oxygen, the integral heats of dissociative adsorption are still negative with respect to the  $\text{WO}_2$  terminated slab and isolated  $\text{O}_2$  molecules, with a value of  $-78.2 \text{ kJ/mol}$ . We note, however, much of the contribution to the negative integral heats at 1 ML oxygen coverage came from occupying the first half of surface W atoms. In fact, the  $\text{O}_2$  dissociative adsorption energy becomes  $448.5 \text{ kJ/mol}$  if the 0.5 ML oxygen covered stoichiometric slab was used as a reference. These results indicate that  $\text{WO}_3(001)$  surfaces with either pure O or pure  $\text{WO}_2$  termination are not stable. The stable configuration corresponds to a surface with 0.5 ML oxygen coverage.[44] This is in general agreement with the experimental

findings from LEED and STM measurements.[45, 46] The striking difference of oxygen binding energies on two surface  $\text{WO}_2$  units can be attributed to the nonequivalent surface W atoms to which the oxygen atoms bind directly. In general, the W-O bond lengths along the surface normal alternate between long and short. This effect was propagated to the surface: the half strongly bonded oxygen happens to occupy the W site, which will form a shorter W-O bond on the surface. Terminating the surface with  $\text{WO}_2$  makes the W-O bonds to the subsurface oxygen almost uniform. Adding 0.5 ML oxygen on the W site corresponding to the shorter W-O bond recovers the “long-short” pattern of W-O bond lengths in bulk  $\text{WO}_3$  along the c-axis and the terminal O-W bond is 1.71 Å. In fact, if oxygen occupies the “wrong” W site, the adsorption energy will be reduced to -162.1 kJ/mol even at 0.5 ML oxygen coverage.

The results of spin-polarized calculations for all three cases:  $\text{WO}_2$ -terminated, 0.5 and 1 ML oxygen covered slabs, showed that only the system with 1 ML oxygen coverage has local magnetic moments. Atom resolved electronic density analysis revealed that the unpaired electrons are primarily localized on the oxygen atom which forms a longer O-W bond, and this surface oxygen atom carries about 1.5 unpaired electrons. Two of the four O atoms in the same plane as the W atoms are only slightly magnetized, each account for 0.25 electrons.

We also found that molecular oxygen may be formed and adsorbed on the surface when the oxygen coverage is larger than 0.5 ML, forming superoxide-like species. This finding is consistent with the photoelectron spectroscopic characterization of  $\text{WO}_x$  clusters[47] and may explain the structural disorder observed in the STM study of oxygen adsorption on  $\text{WO}_3(001)$  when oxygen coverage is larger than 0.5 monolayer.[45]

### 3.2 $\text{ReO}_3(001)$

We then performed a similar analysis for oxygen on  $\text{ReO}_3(001)$  by starting with a nonstoichiometric  $\text{ReO}_2$  terminated slab. The formation of the  $c(2 \times 2)$  oxygen super structure, as shown by the blue square in fig. 2(b), on the basis of  $\text{ReO}_2$  terminated surface is also exothermic, with integral heats of dissociative adsorption of -439.7 kJ/mol. At this oxygen coverage (0.5 ML), the terminal O-Re bond is 1.74 Å, similar to that on  $\text{WO}_3(001)$ . The creation of the surface also makes the Re-O bonds along the surface normal alternate between long and short, although the extent of the oscillation is not as big as that on  $\text{WO}_3(001)$ . Further increasing the oxygen coverage to 1 ML, i.e., covering all the surface Re atoms of the  $\text{ReO}_2$  terminated surface, the



integral heats of dissociative adsorption became -361.1 kJ/mol with respect to the ReO<sub>2</sub> terminated slab and isolated O<sub>2</sub> molecules. This value is significantly larger than the corresponding value on WO<sub>3</sub>(001). Consequently, dissociative adsorption of oxygen on ReO<sub>3</sub>(001) is exothermic even if the 0.5 ML oxygen covered stoichiometric slab was used as a reference (-282.5 kJ/mol). The strong exothermicity of O<sub>2</sub> dissociative adsorption at 1 ML oxygen coverage indicates that it is feasible to maintain a high oxygen coverage on ReO<sub>3</sub>(001). Furthermore, the relative small variation of oxygen adsorption energies upto 1 ML coverage suggests that ReO<sub>3</sub> may be used an oxygen carrier in oxidation reactions. Molecular O<sub>2</sub> adsorption on the surface with 0.5 ML O forming a c(2×2) oxygen super structure was also examined. The results showed that dissociative adsorption is significantly favorable over the molecular adsorption even at the full monolayer oxygen coverage. To the best of our knowledge, there have been no experimental studies of the surface structure based on well-defined ReO<sub>3</sub> single crystal as well as the dependence of the structure on oxygen partial pressures.

### 3.3 Methanol adsorption and dissociation on WO<sub>3</sub> (001)

Methanol adsorption and dissociation have been studied in both the c(2×2) and p(2×2) unit cells with 0.5 ML oxygen coverage on WO<sub>3</sub>(001). The corresponding methanol coverages are 0.5 and 0.25 ML, respectively. Methanol adsorption energies were calculated as  $\Delta E_{\text{ads}} = E_{\text{Me/slab}} - E_{\text{slab}} - E_{\text{methanol}}$ , where  $E_{\text{Me/slab}}$  is the energy of the slab with either molecularly or dissociatively adsorbed methanol,  $E_{\text{slab}}$  represents the energy of WO<sub>3</sub> (or ReO<sub>3</sub> in the following section) slab terminated with 0.5 ML oxygen, and  $E_{\text{methanol}}$  is the energy of an isolated methanol molecule. The results were summarized in table 1.

The top and side views of a molecularly adsorbed methanol in a p(2×2) surface unit cell were shown in fig. 3. In this configuration, the oxygen atom of the methanol hydroxyl binds the surface W atom. The C-O and O-H bond distances in adsorbed methanol are 1.45 and 0.99 Å, respectively, and are only slightly longer than those calculated with an isolated methanol molecule (1.43 and 0.97 Å, respectively). The distance between of the methanol O to the surface W is 2.33 Å. This configuration also allows the formation of an O-H...O structure, with a H...O distance of 1.80 Å. The hydrogen bond-like structure helps to stabilize the molecularly adsorbed methanol. As shown in table 1, the adsorption energy for methanol calculated in a p(2×2) unit

cell is slightly larger than that in a c(2×2) unit cell, indicating a weak repulsive interactions between the molecularly adsorbed methanol as the coverage is increased.

Two dissociation channels: forming  $\text{CH}_3\text{O}_{(a)}+\text{H}_{(a)}$  and  $\text{CH}_2\text{OH}_{(a)}+\text{H}_{(a)}$  were analyzed. The dissociative adsorption energies with respect to an isolated methanol molecule were also listed in Table 1. Obviously, both dissociation channels are endothermic. In  $\text{CH}_3\text{O}_{(a)}+\text{H}_{(a)}$ , the hydrogen adatom occupies the neighboring surface oxygen site and forms a surface hydroxyl group, as shown in figure 4. This configuration with coadsorbed methoxide and H, circled in figure 4, can only be maintained by rotating the newly formed O-H bond away from the original plane. In fact, relaxation from adsorbed methoxide on W site and hydrogen adatom on the surface oxygen site with the newly formed O-H bond in the same plane as  $\text{O-H}\cdots\text{O}$  led to H moving back to the O atom of the adsorbed methoxide and forming molecularly adsorbed methanol. The O-H bond in the surface hydroxyl is 0.97 Å. The C-O and O-W bond distances in the adsorbed methoxide are 1.40 and 1.93Å, respectively, shorter than the corresponding values in the molecularly adsorbed methanol.

An alternative pathway involves extracting hydrogen from the methyl group of methanol to form coadsorbed hydroxymethyl and hydrogen adatom, i.e.,  $\text{CH}_2\text{OH}_{(a)}+\text{H}_{(a)}$ . In this configuration,  $\text{CH}_2\text{OH}$  binds to surface W atom through the C atom, forming a C-W bond of 2.35 Å. Similarly, the hydrogen adatom from the dehydrogenation forms a surface hydroxyl group with the terminal oxygen. The calculated dissociative adsorption energy for this product state shows that this channel is even more endothermic, with heats of reaction being 94.4 and 114.4 kJ/mol at 0.5 and 0.25 ML methanol coverage, respectively. Due to the high endothermicity of both reaction channels on  $\text{WO}_3(001)$ , we did not perform transition state calculations for dissociation of adsorbed methanol.

### 3.4 Methanol adsorption and dissociation on $\text{ReO}_3(001)$

In order to compare the reactivity of  $\text{ReO}_3$  with  $\text{WO}_3$ , we examined methanol adsorption and dissociation only in a p(2×2) surface unit cell with 0.5 ML oxygen coverage. The top and side views of a molecularly adsorbed methanol in the p(2×2) surface unit cell were shown in fig. 5. Similar to methanol on  $\text{WO}_3(001)$ , the oxygen atom of the methanol hydroxyl binds the surface Re atom. The C-O and O-H bond distances of adsorbed methanol on  $\text{ReO}_3(001)$  are almost the same as those on  $\text{WO}_3(001)$ . The distance between of the methanol O to the surface Re is 2.21 Å,

shorter by 0.12 Å than the corresponding O-W distance. The H···O distance in the O-H···O structure on ReO<sub>3</sub>(001) is 2.04 Å., significantly longer than that on WO<sub>3</sub>(001). Molecular adsorption was found to be exothermic, with an adsorption energy of -106.4 kJ/mol.

We also examined two dissociation channels, i.e., forming CH<sub>3</sub>O<sub>(a)</sub>+H<sub>(a)</sub> and CH<sub>2</sub>OH<sub>(a)</sub>+H<sub>(a)</sub>, on ReO<sub>3</sub>(001). The dissociative adsorption energies with respect to an isolated methanol molecule were also listed in Table 1. Obviously, both dissociation channels are exothermic. In CH<sub>3</sub>O<sub>(a)</sub>+H<sub>(a)</sub>, the hydrogen adatom occupies the neighboring surface oxygen site and forms a surface hydroxyl group, as shown in figure 6. In contrast to methanol dissociation on WO<sub>3</sub>(001), the coadsorbed methoxide and H on ReO<sub>3</sub>(001), circled in figure 6, is a minimum on the potential energy surface. The O-H bond in the surface hydroxyl is 0.98 Å. The C-O and O-Re bond distances in the adsorbed methoxide are 1.42 and 1.93 Å. The H···O distance in the O-H···O structure is 2.42 Å.

We further explored the reaction pathways to form CH<sub>3</sub>O<sub>(a)</sub>+H<sub>(a)</sub> and searched for transition state along this path on ReO<sub>3</sub>(001). Five images were interpolated between the molecularly adsorbed initial state (fig. 5) and the dissociative state (fig. 6). The transition state determined from these images by using NEB method was further optimized using quasi-Newton method. The final structure of the transition state is shown in figure 7 with the key structure parameters labeled. The H atom in the transition state is located almost in the middle of the surface and methanol O atoms, with distances being 1.195 and 1.267 Å, respectively. The barrier is only 19.1 kJ/mol with respect to the molecularly adsorbed methanol. This transition state is 87.3 kJ/mol lower than the energy of the gas phase methanol, suggesting that the molecularly and dissociatively adsorbed methanol must be in equilibrium on the surface. Therefore, operating conditions such as water vapor pressure will affect the surface coverage of hydroxyls, and in turn will determine relative abundance of surface molecular methanol and methoxide.

The pathway leading to the formation of coadsorbed hydroxymethyl and hydrogen adatom, i.e., CH<sub>2</sub>OH<sub>(a)</sub>+H<sub>(a)</sub>, by extracting hydrogen from the methyl group of methanol, was found to be energetically favorable with respect to gas phase methanol molecule. CH<sub>2</sub>OH binds to surface Re atom through the C atom and the C-Re bond is 2.18 Å. The hydrogen adatom from the dehydrogenation forms a surface hydroxyl group with the terminal oxygen and the O-H bond length is 1.01 Å. The dissociative adsorption energy for this product state is -23.7 kJ/mol.

Therefore, breaking C-H bond directly is possible on  $\text{ReO}_3(001)$  and will lead to a different product distribution for further oxidation of methanol.

### 3.5 General Discussion

Clearly, the reactivity of  $\text{WO}_3$  and  $\text{ReO}_3$  towards O-H and C-H bond cleavage in methanol is very different. On  $\text{WO}_3(001)$  with 0.5 ML oxygen, the heterolytic breaking of methanol O-H bond is clearly not favorable. In contrast, breaking O-H bond is energetically favorable on  $\text{ReO}_3(001)$ . Even breaking C-H bond in methanol becomes feasible on  $\text{ReO}_3(001)$ . In order to understand the difference between these two metal oxides surfaces, we analyzed the detailed local density of states of these adsorption systems. In figure 8, we plotted the local density of states of W and O in molecularly adsorbed methanol and adsorbed methoxide on  $\text{WO}_3(001)$ . For comparison, we plotted in figure 9 the LDOS of Re and O for molecularly adsorbed methanol and adsorbed methoxide on  $\text{ReO}_3(001)$  as well as at the transition state of methanol dissociation in figure 9. The LDOS were plotted by aligning at the Fermi level of the corresponding system. Detailed analysis showed that the O 2p states at energies of -2.5 eV and -7 eV for  $\text{CH}_3\text{OH}/\text{WO}_3(001)$  and -4.5 eV and -9 eV for  $\text{CH}_3\text{OH}/\text{ReO}_3(001)$  are associated with the O-H bond of the adsorbed methanol. These states disappeared in the adsorbed methoxide on both oxide surfaces. We can see that the 2p levels of O are located at the top of the valence band for methanol on  $\text{WO}_3(001)$ . The contribution of W at this level is not pronounced. Consequently, the coupling between the O 2p levels of methanol (methoxide) and W 5d is not strong. On  $\text{ReO}_3(001)$ , the contribution to the states close to the Fermi level comes mainly from the 5d states of Re allowing for a better coupling with the energy levels of the O-H bond of methanol. The O 2p contribution is located at  $\sim -2$  eV below Fermi level. Of course, the coverage of oxygen is expected to affect the surface chemistry of methanol. Increasing the oxygen coverage beyond 0.5 ML will likely reduce the sites available for methanol adsorption although the barrier for O—H bond breaking is not expected to change significantly as the process only involves the Re-bound methanol and the neighbouring terminal oxygen. On the other hand, other reaction channels such as C—O bond breaking may become accessible at lower oxygen coverages.

The molecular structures of the active metal oxide species are expected to play important roles as the heterogeneous catalytic reactions occur at the surfaces of catalysts.[48] There have been many studies attempted to relate the molecular structures of the supported metal oxide

species to their catalytic properties. Many physical characterization techniques have been developed to probe different aspects of the catalysts “off-situ” or “in-situ”. For example, Wachs and coworkers have employed Raman and IR to characterize a wide range of supported metal oxides including  $\text{WO}_x$  and  $\text{ReO}_x$ . [49] Their studies provided some useful insight into the molecular structures under different operating conditions. The degree of hydration of the metal oxide surfaces has strong influence on the reactions that require the acidic/basic sites on the surface. [50, 51] In fact, water adsorption and dissociation on these surfaces follow a similar trend to that of the methanol, i.e.  $\text{H}_2\text{O}$  favors molecular adsorption strongly on the perfect  $\text{WO}_3(001)$  whereas the molecularly and dissociatively adsorbed  $\text{H}_2\text{O}$  on  $\text{ReO}_3$  have a similar stability. The affinities of the oxide surfaces to water and their ability to dissociatively adsorb water have been shown to have a significant on the oxides as catalysts or support/promoters for a catalyst. [39, 52-55]

#### 4. Conclusions

First principles density functional theory was used to analyze the adsorption and dissociation of methanol on the surfaces of two structurally similar transition metal oxides,  $\text{WO}_3(001)$  and  $\text{ReO}_3(001)$ . The adsorption energies of molecularly adsorbed methanol on  $\text{WO}_3(001)$  and  $\text{ReO}_3(001)$  are -74.0 and -106.4 kJ/mol, respectively. On the other hand, dissociative adsorption of methanol to adsorbed methoxide species is not energetically favored on  $\text{WO}_3(001)$  but favored on  $\text{ReO}_3(001)$ . The dissociation to form coadsorbed methoxide and hydrogen adatom has been found 34.8 kJ/mol endothermic on  $\text{WO}_3$  and -111.2 kJ/mol exothermic on  $\text{ReO}_3$ , respectively. The activation barrier for the dissociation of molecularly adsorbed methanol on  $\text{ReO}_3(001)$  was determined to be 19.1 kJ/mol. Dehydrogenations to form coadsorbed hydroxymethyl and hydrogen adatom are not energetically favorable on both surfaces with respect to the molecularly adsorbed methanol. However, the reaction with respect to gas phase methanol is exothermic on  $\text{ReO}_3$  with the heats of reaction of -23.7 kJ/mol but endothermic on  $\text{WO}_3$  with the heat of reaction of 114.4 kJ/mol. We conclude that the differences between electronic and crystal structure of  $\text{ReO}_3$  and  $\text{WO}_3$  contribute to differences in chemical properties of the (001) surfaces of these trioxides.

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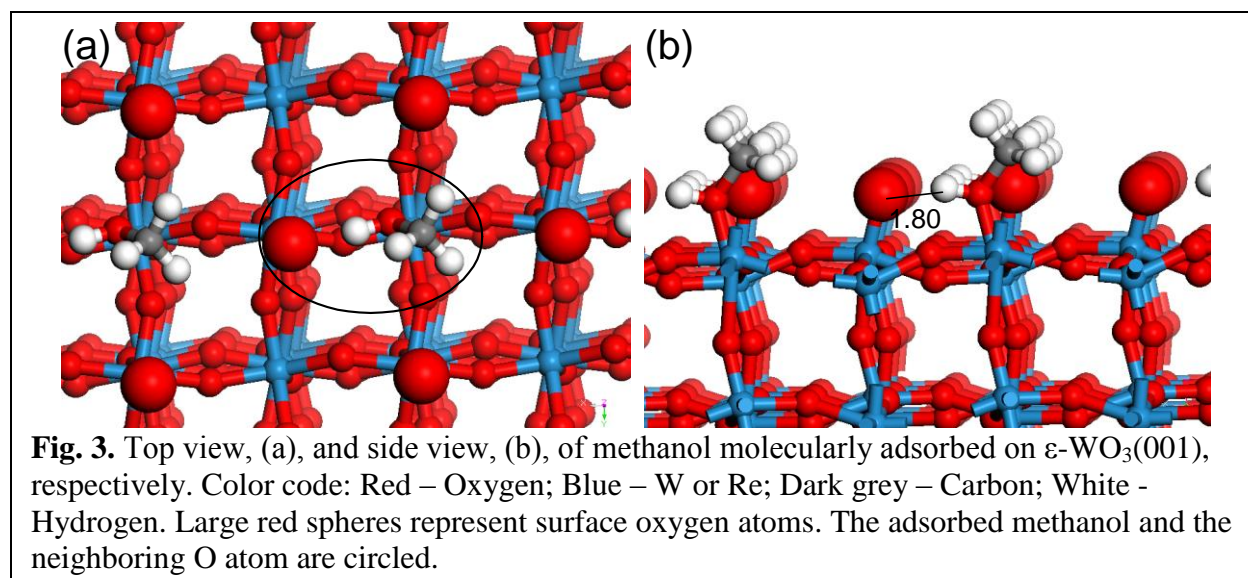
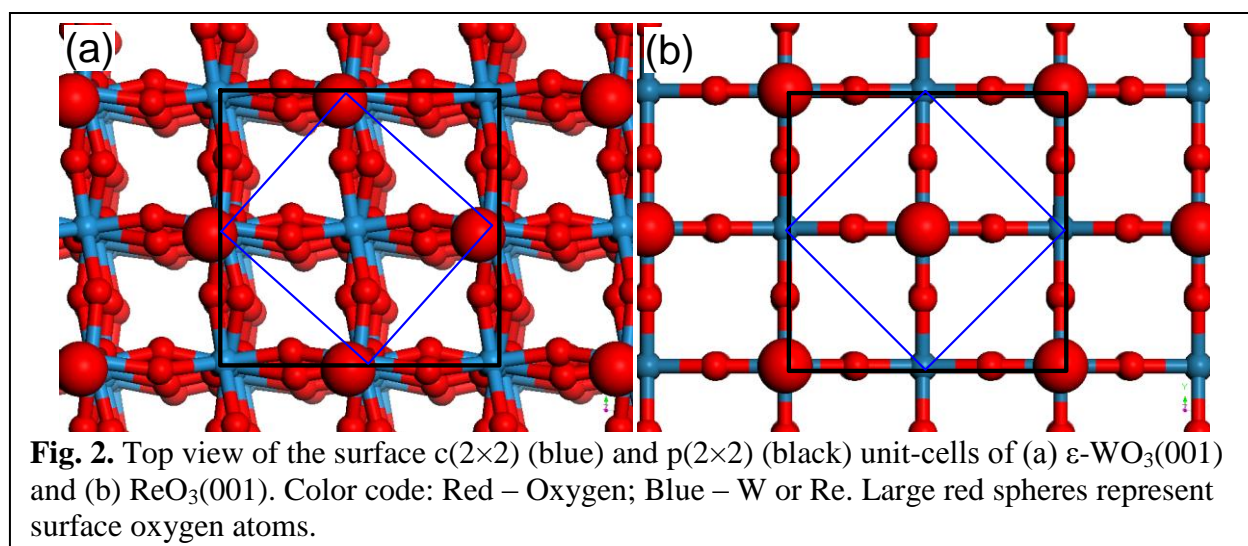
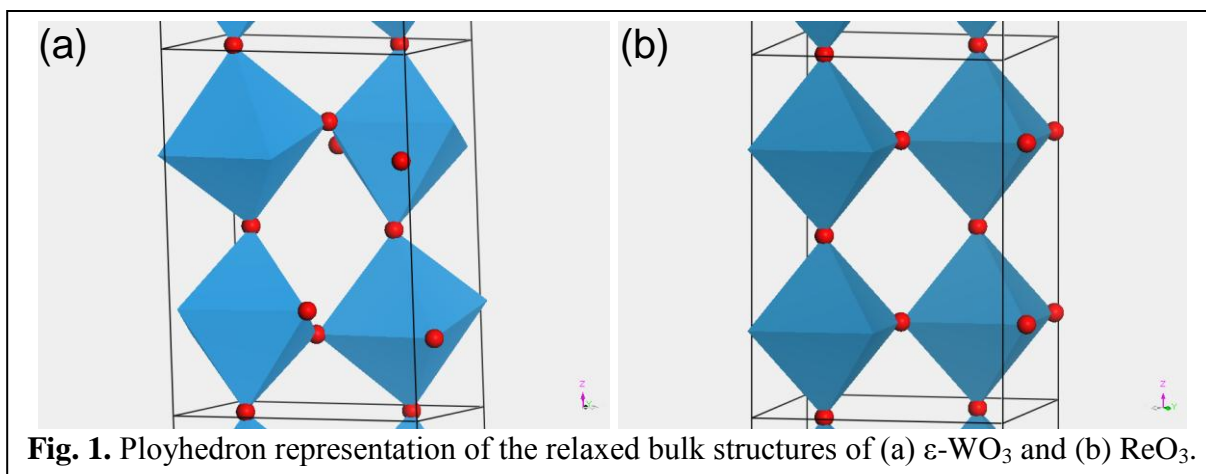
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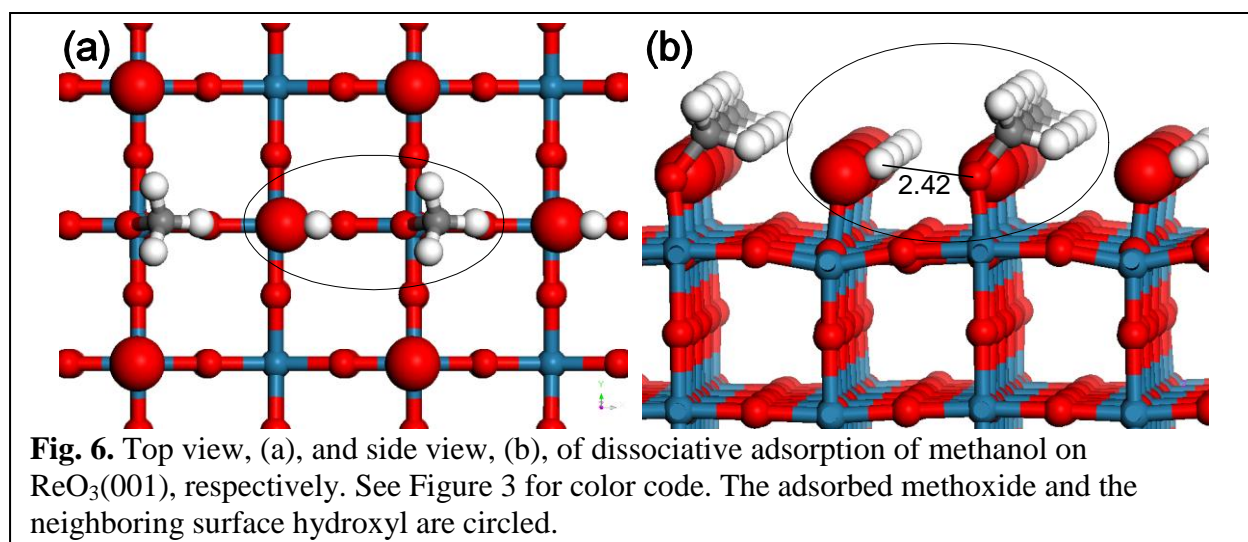
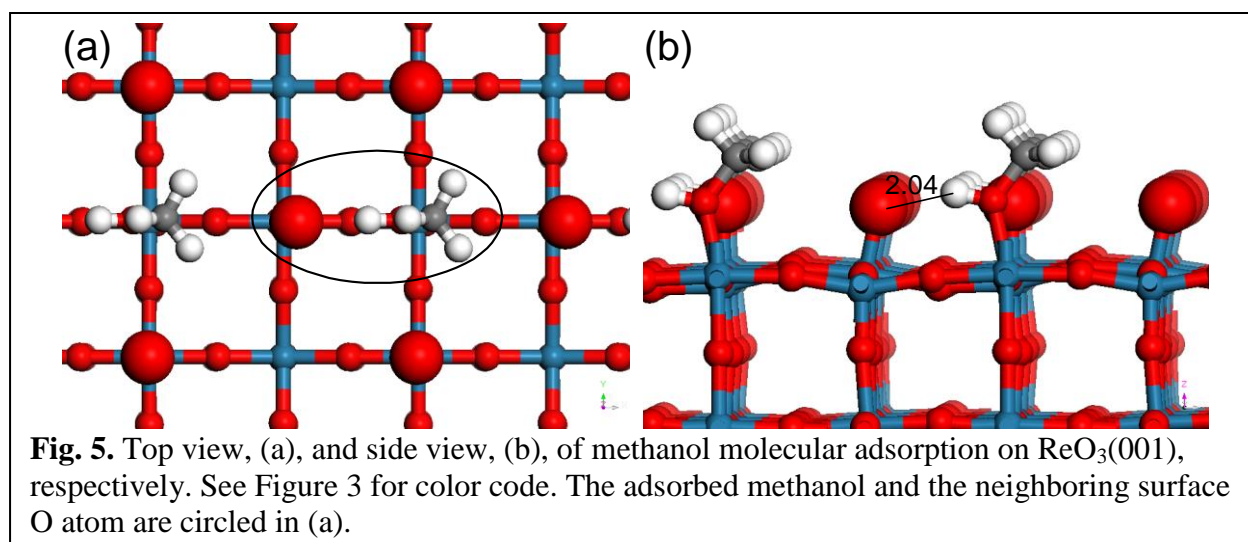
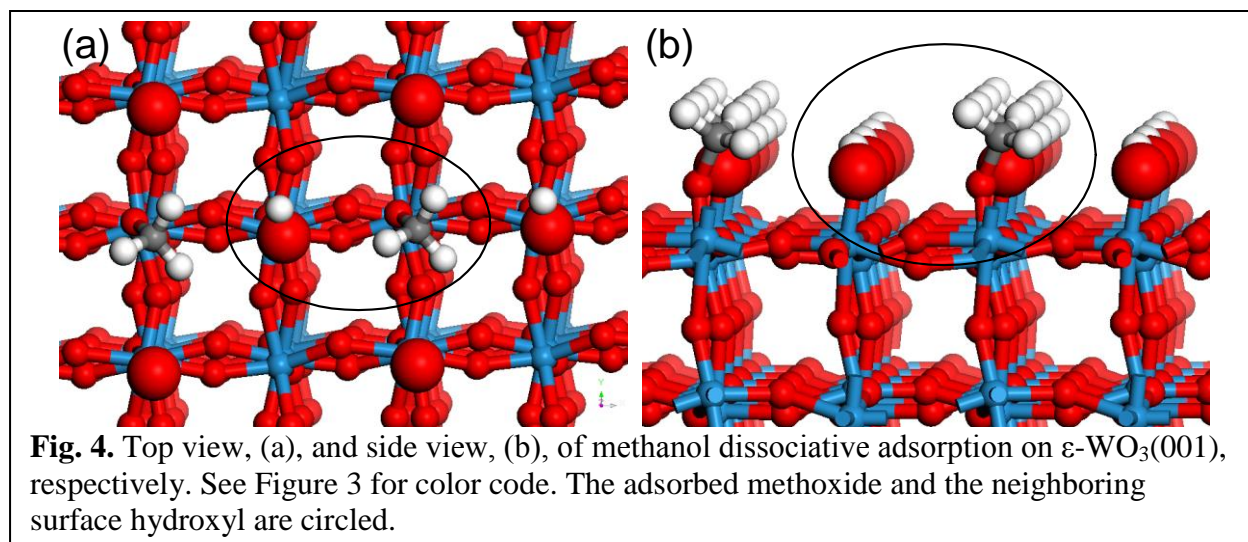


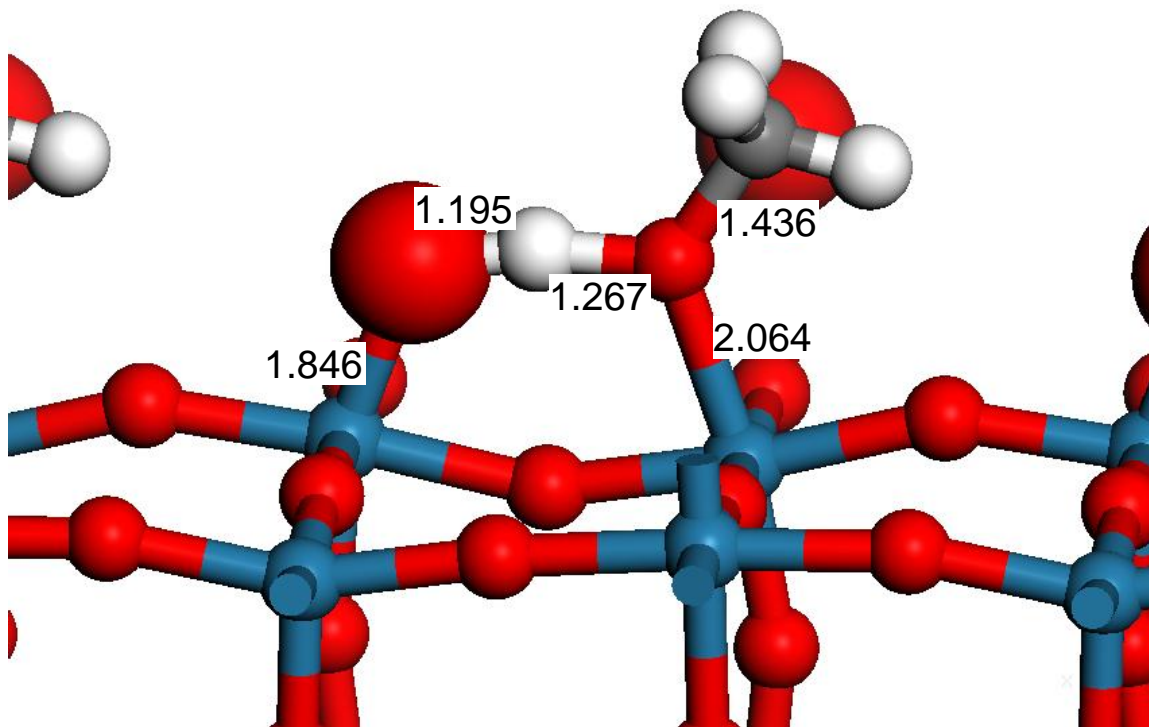
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**Table 1.** DFT GGA calculated adsorption energies of molecular and dissociative adsorption on  $\text{WO}_3(001)$  and  $\text{ReO}_3(001)$

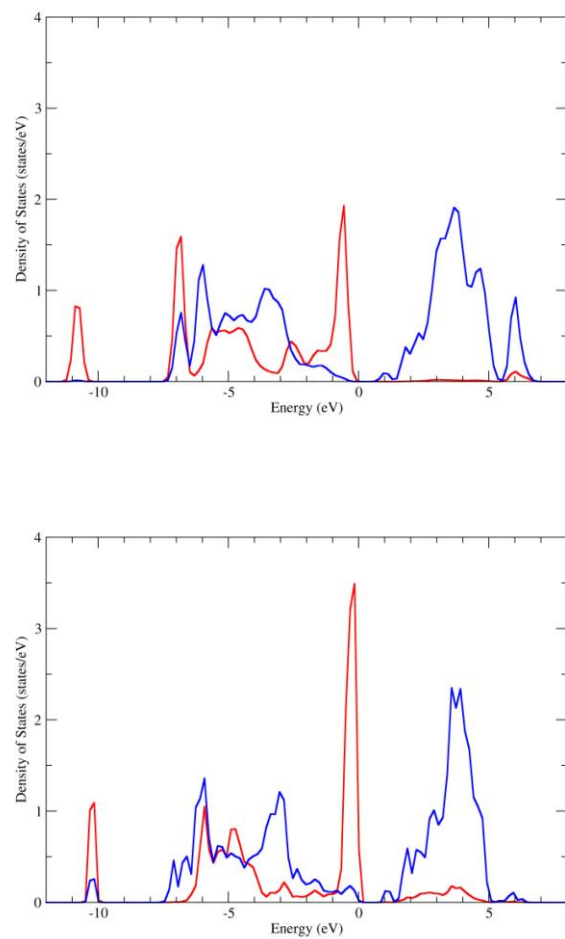
	$\text{WO}_3(001)$		$\text{ReO}_3(001)$
	c(2x2) (kJ/mol)	p(2x2) (kJ/mol)	p(2x2) (kJ/mol)
$\text{CH}_3\text{OH}_{(a)}$	-74.1	-74.0	-106.4
$\text{CH}_3\text{O}_{(a)}+\text{H}_{(a)}$	15.0	34.8	-111.2
$\text{CH}_2\text{OH}_{(a)}+\text{H}_{(a)}$	94.4	114.4	-23.7



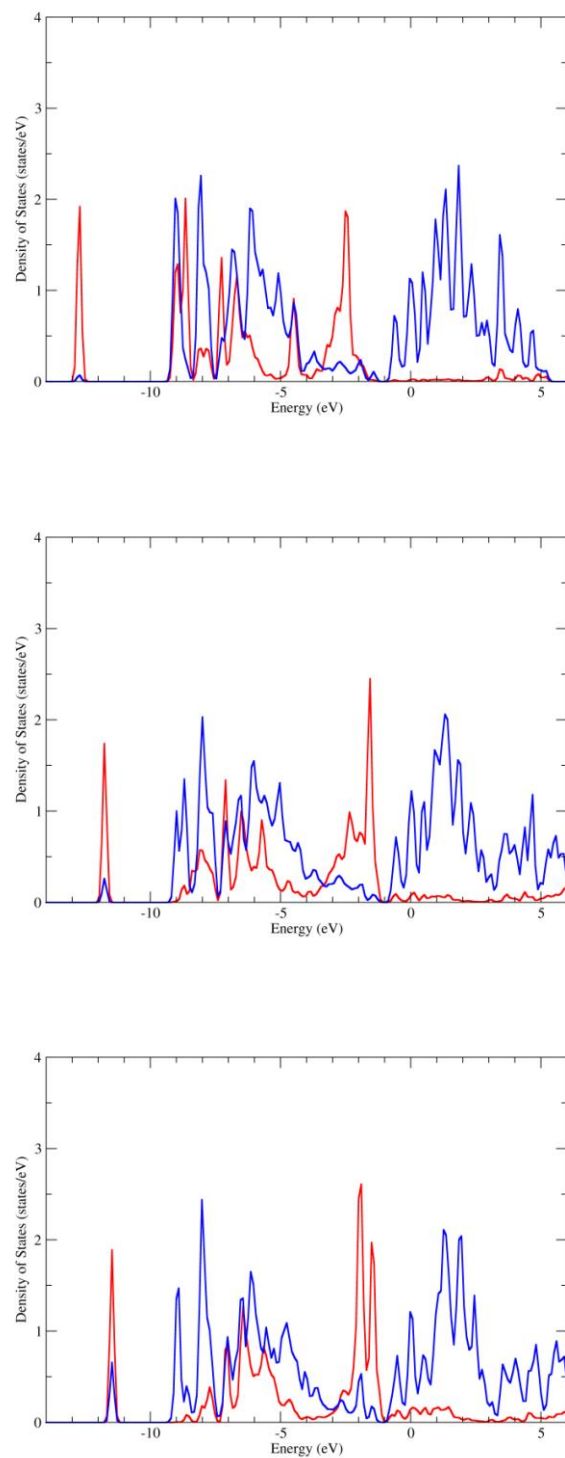




**Fig. 7.** Transition state structure for methanol dissociation on ReO<sub>3</sub>(001). See Figure 3 for color code.



**Fig. 8.** Local density of states plots of O (red) and W (blue) in molecularly adsorbed methanol (upper) and adsorbed methoxide (lower).



**Fig. 9.** Local density of states plots of O (red) and Re (blue) in molecularly adsorbed methanol (top), transition state (middle), and adsorbed methoxide (bottom). Zero on x-axis corresponds to Fermi energy.