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Open topological defects and boundary RG flows

Anatoly Konechny 

Department of Mathematics, Heriot-Watt University, Edinburgh EH14 4AS,
United Kingdom
Maxwell Institute for Mathematical Sciences, Edinburgh, United Kingdom

E-mail: a.konechny@hw.ac.uk

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Abstract

In the context of two-dimensional rational conformal field theories we consider topological junctions of topological defect lines with boundary conditions. We refer to such junctions as open topological defects. For a relevant boundary operator on a conformal boundary condition we consider a commutation relation with an open defect obtained by passing the junction point through the boundary operator. We show that when there is an open defect that commutes or anti-commutes with the boundary operator there are interesting implications for the boundary RG flows triggered by this operator. The end points of the flow must satisfy certain constraints which, in essence, require the end points to admit junctions with the same open defects. Furthermore, the open defects in the infrared must generate a subring under fusion that is isomorphic to the analogous subring of the original boundary condition. We illustrate these constraints by a number of explicit examples in Virasoro minimal models.

Keywords: renormalisation group flows, topological defect lines, boundary conformal field theories

(Some figures may appear in colour only in the online journal)

1. Introduction

In this paper we consider Euclidean two-dimensional quantum field theories on a half-plane which are described by a unitary conformal field theory (CFT) in the bulk, and on the boundary by a perturbed conformal boundary condition. We will assume that the CFT at hand is rational and thus possesses some non-trivial topological defect lines. Such defects were first considered



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in [1] and then more extensively in the context of general rational CFTs in [2]. They describe symmetries and dualities of the critical system described by the given CFT [2, 3]. Topological defect lines can be moved around and, if they do not pass through any other observables any correlation function is independent of their position. When they pass through a local bulk operator, generically we obtain a collection of defect segments attached to the main defect and ending on a disorder field located at the insertion. In certain special situations the additional defect segments may be absent and passing the defect through results in an operator with the same Virasoro representation labels but multiplied by some factor. In the simplest situation the original insertion remains intact and we can say that the defect commutes with this bulk insertion. As shown in [4], if we have some defects that commute with a bulk relevant operator then there are interesting consequences for the bulk renormalisation group (RG) flows triggered by this operator. The fusion algebra of such commuting defects between themselves must be robust under the fusion and this places constraints on the end points of the flows (triggered by the same operator with positive or negative coupling) particularly when the flows are massive and the end points may be described by non-trivial topological theories.

For a CFT on a half plane with a conformal boundary condition, if there is a non-trivial boundary relevant operator, we can perturb the boundary condition by this operator triggering a boundary RG flow. Unlike bulk flows boundary RG flows always end up in a non-trivial conformal boundary condition that at least has the Virasoro identity tower in the boundary spectrum. In the presence of topological defect lines in the bulk CFT we can fuse them with any conformal boundary condition to obtain a new conformal boundary condition which may in general be a superposition of elementary boundary conditions. Based on this construction, an interesting interplay between boundary RG flows and topological defect lines was discussed in [5]. The following general theorem was proved in that paper: given a boundary RG flow from a maximally symmetric conformal boundary condition with label a that ends in a maximally symmetric conformal boundary condition with label b , for any topological defect d there is an RG flow from $d \times a$ to $d \times b$ where the cross stands for fusion. By maximally symmetric we mean here a boundary condition preserving the complete chiral algebra of our rational CFT. We will refer to this result as Graham–Watts theorem in the rest of the paper. The perturbing field for the new flow must have the same Virasoro representation properties (and scaling dimension in particular) as the perturbing field in the original flow. As the new starting point may be a direct sum of elementary boundary conditions there may be many such fields. The precise form of the perturbing field of the new flow has been worked out in [5] for the case of a being elementary and for the general situation it was worked out in [6]. For diagonal modular invariants the action of a defect on boundary fields can be also obtained from the action on chiral defect fields which was worked out in [7].

If we know the end-point for a particular boundary flow, using Graham–Watts theorem we can find the end-points for other flows obtained via fusion. Thus, in [5], using the results of [8], an extension of perturbative flows triggered by boundary $\psi_{1,3}$ fields in minimal models [9] to all Cardy boundary conditions was obtained. It is interesting to note that the g -factors change under fusion according to

$$\sum_{i \in d \times a} g_i = g_a \left(\frac{g_d}{g_1} \right) \quad (1.1)$$

where g_1 is the g -factor associated with the Cardy boundary condition that has only the identity tower in its spectrum. It is not hard to show that in unitary rational CFTs g_1 is the smallest possible g -factor (see e.g. [10]). Thus, (1.1) implies that fusion with non-trivial topological defect always increases the g -factor. A useful strategy in applying the Graham–Watts

theorem may then be to start with a UV boundary condition with a small value of the g -factor, use the g -theorem [11, 12] and symmetries to constrain the end point as much as possible then use fusion to obtain possible end points for flows that start with larger values of g .

In this paper we look at a different usage of topological defects for constraining boundary flows that not merely relates two different flows but directly constrains the possible end points for a given flow. We consider topological junctions of topological defects with a conformal boundary condition. This means that not only the part of the defect line that extends into the bulk but the junction point as well can be moved along the boundary, not changing any correlation functions as long as no boundary insertions are encountered. Such junctions and their properties were considered at length in [6] and we will use the results of that paper extensively. Following [6] we call a topological defect attached to a conformal boundary via a topological junction an open topological defect. When we move such a defect along the boundary with an insertion of a boundary operator present, passing the open defect through the insertion typically results in a configuration with the original insertion replaced by several boundary condition changing fields and new boundary conditions between the insertions and the open defects. But sometimes, for certain defects and boundary fields, no additional fields or boundary conditions arise, the open defect just passes through. In the operator language the defect and the boundary operator commute. In such cases, which are similar to the bulk case considered in [4], we can argue that the end point of the boundary flow must admit a topological junction with the same defect. Moreover, the ring obtained by fusion of such open defects between themselves must be isomorphic to some subring in the infrared boundary condition. This potentially can lead to additional constraints on the end points of RG flows as in the boundary case the fusion ring for open defects in general depends on the boundary condition [6]. Even if the bulk labels are the same the fusion rules may be different. We illustrate this on an explicit example in section 4.3.

Another interesting case is when an open defect just multiplies the operator by minus one when passing through it, or in other words when the open defect anti-commutes with the boundary operator. This situation demands that there must be a topological junction with the same defect and the two boundary conditions describing the infrared end points for the two signs of the perturbation. The fusion rules again must be robust (up to isomorphism) and persist into the infrared fixed points. If both commuting and anti-commuting open defects are present they form a \mathbb{Z}_2 -graded subring under fusion.

The main goal of this paper is to point to the existence of such constraints on boundary RG flows, to explain how to look for commuting and anti-commuting open defects and to illustrate the resulting constraints on concrete examples. To this end we choose to restrict our constructions to Virasoro minimal models with diagonal modular invariant. Moreover, our main examples of boundary flows will be the flows triggered by boundary $\psi_{1,3}$ operators. These flows are integrable and the end points of the flows are known. This allows us to check that the constraints we derive from open defects are satisfied. In addition we consider two flows: one triggered by a boundary $\psi_{2,1}$ operator in the tetracritical Ising model and another triggered by a $\psi_{1,2}$ operator in the pentacritical Ising model. These flows are believed to be integrable but, to the best of our knowledge, have not been investigated before. We derive a number of analytic constraints on the possible end points in these flows.

The main body of the paper is organised as follows. In section 2 we discuss generalities about topological defects and their junctions with boundary conditions. We fix our normalisation conventions and derive a commutation relation for an open defect and a boundary operator. In section 3 we discuss the constraints on RG flows arising from open defects commuting or

anti-commuting with the perturbing operator. In section 4 we work out explicit examples in the tetracritical and the pentacritical Ising models. In section 5 we discuss some specifics for flows triggered by boundary condition changing operators. For such flows there may be special linear combinations of different open defects (with the same Virasoro labels) that commute with the perturbing field. We give some explicit examples of this. Section 6 contains some concluding remarks. The appendix A contains some useful relations between the diagonal minimal model fusion matrices.

2. Open topological defects

Throughout the paper, except for section 3, we restrict ourselves to the case of unitary Virasoro minimal models with diagonal modular invariant. For the minimal models both topological defects [1] and the elementary conformal boundary conditions [13] are labelled by the same pairs of integers from the Kac table as the chiral operators. In this section we will just use the letters: a, b, c, \dots for such labels. Boundary fields linking a boundary condition a on the left with a boundary condition b on the right are built on Virasoro representations $i \in a \times b$. We denote such fields as $\psi_i^{[a,b]}$. On the diagrams below we will omit the upper indices of boundary operators as those can be read off from the boundaries.

Three elementary topological defects labelled by a, b, c can be joined together if that is permitted by the fusion, that is if $a \in b \times c$. Defect networks can be simplified via a sequence of elementary moves. The latter equates two networks as depicted on figure 1. This was shown in [14] using the topological field theory approach of [2].

As emphasised in [6] one does not have to choose the F -matrices appearing on figure 1 to be the same as the conformal block F -matrices. The latter are fixed if we canonically normalise the conformal blocks. To do concrete calculations we are going to use the conformal block F -matrices calculated in [15, 16], so we are going to assume that the defect junctions are normalised in such a way that the defect F -matrices are those of the conformal blocks. We also assume that the identity defect can be attached at any point and can be moved freely without changing anything.

We are further interested in topological defects that can end topologically on a given conformal boundary condition. This means that the ending should behave as a local operator of dimension zero. It is not hard to see that for an elementary boundary condition with a label a and an elementary defect with label d the junction is topological if the fusion $d \times a$ contains a . To see this we can deform the defect keeping the junction pinned down so that the defect fuses with the boundary on one side of the junction. The junction then looks like a boundary condition-changing operator between a and $d \times a$. There is a dimension zero such operator if $d \times a$ contains a . Equivalently $a \times a$ should contain d and therefore the set of admissible defect labels d is the same as the set labelling boundary operators. This observation generalises to junctions which have two different elementary boundary conditions on either side of the junction: a and b . The junction is topological if there is a fusion vertex linking a, b and d . Such a junction is depicted on figure 2.

Each junction of an elementary defect and an elementary boundary comes with a choice of coupling that can be thought of as a choice of normalisation of the corresponding junction field. We are going to choose the normalisation for the open defects and the boundary fields as described in [6]. The conventions of [6] include additional factors for the junctions of defects with boundaries which arise from taking a defect stretched parallel to the boundary and partially fusing its right or left half with the boundary. To distinguish between the two types of fusion it will be convenient to orient our defects assuming that the defect outgoing from the boundary was fused on the left and the defect coming into the boundary was fused on the right.

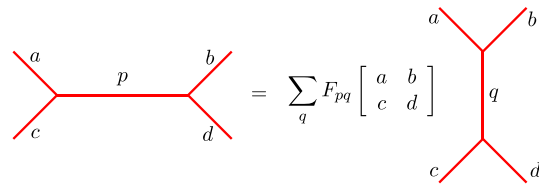


Figure 1. Elementary move in a defect network.

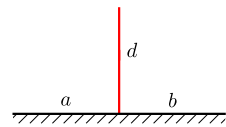


Figure 2. A defect junction with a conformal boundary.

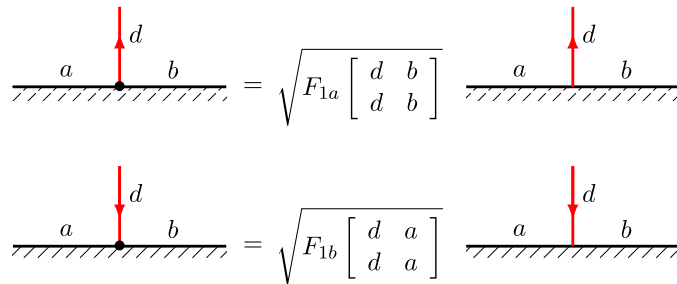


Figure 3. Normalisation factors for oriented defect junctions.

The orientation will be marked by arrows on the diagrams. Furthermore, to signify the presence of these additional factors we will add a bullet on the junction when depicting it. The factors themselves are presented on figure 3.

We will denote a closed elementary defect located in the bulk as \mathcal{D}_a while the open defect corresponding to the left hand side of the first diagram on figure 3 will be denoted as $\mathcal{D}_d^{[a,b]}$ and, for brevity, we will write $\mathcal{D}_d^{[a]}$ instead of $\mathcal{D}_d^{[a,a]}$. We will write $\mathcal{D}_d^{[a]}(t)$ to denote an insertion of such a defect ending at point t on the boundary inside correlation functions of boundary operators. Also we will denote as $\mathcal{D}_d^{[a]}$ the corresponding operator acting on the radial quantisation states on a half plane with the boundary condition a .

With the factors given on figure 3 two simple relations hold. Firstly, a defect arc attached to the boundary with no insertions can be shrunk leaving no additional factors. This is illustrated on figure 4.

Secondly, when we partially fuse a portion of a defect with the boundary we obtain a sum over elementary boundary conditions appearing in the fusion and two junctions with the boundary. This is illustrated on the following diagram (figure 5).

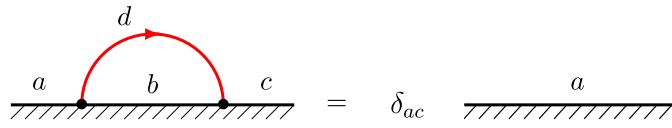


Figure 4. Shrinking an open defect bubble.

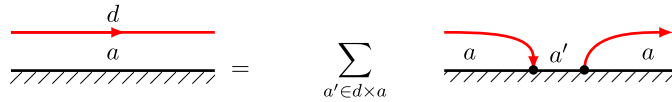


Figure 5. Partial fusion of defect with a boundary.

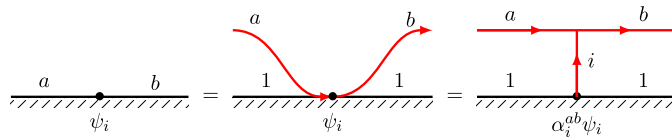


Figure 6. Trading a boundary field for a defect ending field.

Manipulations with junctions of defects with a boundary can be lifted to junctions between topological defects by representing the boundary conditions with label a as fusions between \mathcal{D}_a and the identity boundary condition. A boundary operator with Virasoro label i can be traded for the defect labelled by i ending with a defect ending field located on the identity boundary condition. This is shown on figure 6.

The general expression for coefficients α_i^{ab} has been calculated in [6] (see equation (B.7) of that paper). Once the F -matrices appearing in defect junctions have been fixed, these coefficients can be explicitly calculated. In this paper we use the conformal block F -matrices so in principle α_i^{ab} are fixed but at no point in our calculations we need to use their explicit form. Using figure 6 we calculate, following [6], the action of an open defect on boundary operators. The latter is obtained by surrounding a boundary operator by the defect arc and shrinking the arc onto the operator. This can be calculated by a sequence of moves shown on figure 7 where we consider the most general boundary condition changing operator.

The final factors $X_{i,a\tilde{a}}^{a'a''}$ that appear on figure 7 are

$$X_{i,a\tilde{a}}^{a'a''} = F_{a''a} \begin{bmatrix} d & a' \\ \tilde{a} & i \end{bmatrix} \sqrt{\frac{F_{1\tilde{a}} \begin{bmatrix} d & a'' \\ d & a'' \end{bmatrix}}{F_{1a} \begin{bmatrix} d & a' \\ d & a' \end{bmatrix}}} \left(\frac{\alpha_i^{a'a''}}{\alpha_i^{a\tilde{a}}} \right). \tag{2.1}$$

An alternative derivation of this result can be done using the three-dimensional topological quantum field theory representation developed in [17, 18].

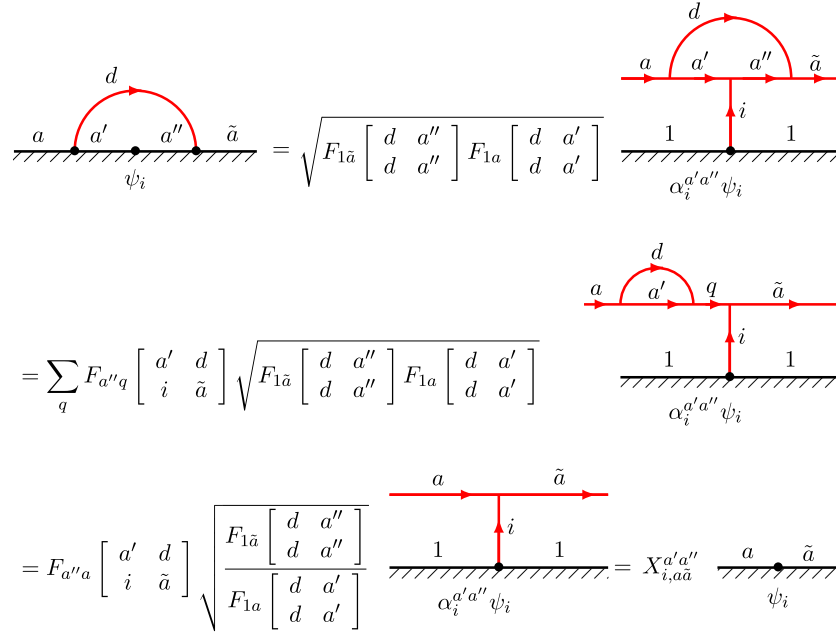


Figure 7. Action of an open defect on a boundary field.

Using figure 7 we can derive a commutation relation between an open defect and an insertion of a boundary operator ψ_i . To that end we need to pass the defect junction through ψ_i from left to right. This can be done by creating an arc around the insertion of ψ_i , partially fusing a portion of the defect to the right of the insertion and finally shrinking the arc onto the boundary field. This is depicted on figure 8.

For the particular case of a boundary operator on an elementary boundary we have $a = a' = a''$ and the factors in the commutation relations become

$$X_{i,a\tilde{a}}^{aa} = F_{a\tilde{a}} \begin{bmatrix} d & a \\ a & i \end{bmatrix} \sqrt{\frac{F_{1a} \begin{bmatrix} d & a \\ d & a \end{bmatrix}}{F_{1\tilde{a}} \begin{bmatrix} d & a \\ d & a \end{bmatrix}}} \left(\frac{\alpha_i^{aa}}{\alpha_i^{a\tilde{a}}} \right). \tag{2.2}$$

We see from this expression that the defect commutes with ψ_i if

$$F_{a\tilde{a}} \begin{bmatrix} d & a \\ a & i \end{bmatrix} = \delta_{a\tilde{a}} \tag{2.3}$$

and it anti-commutes if

$$F_{a\tilde{a}} \begin{bmatrix} d & a \\ a & i \end{bmatrix} = -\delta_{a\tilde{a}}. \tag{2.4}$$

We can also conclude from the orthogonality relation (A.6) that these are the only interesting situations for RG flows originating from an elementary boundary condition, there cannot be a commutation up to a non-trivial rescaling of ψ_i . The latter however are possible when boundary condition changing fields are involved (see section 5).

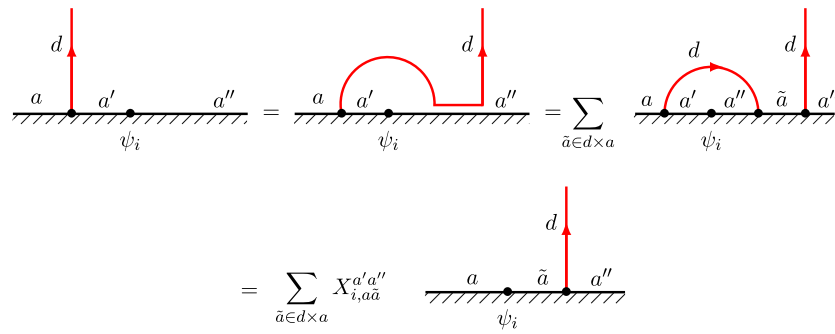


Figure 8. Commutator of open defect with a boundary operator.

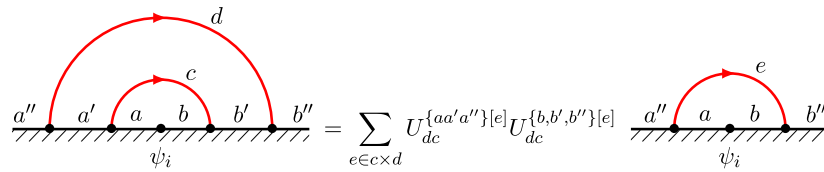


Figure 9. Fusion of two open defects.

Two arcs of open defects surrounding a boundary operator can be fused into a combination of open defects according to figure 9.

The coefficients on the right hand side of figure 9 were worked out in [6]. They are given by the following combinations of the fusion matrices

$$U_{dc}^{\{aa'a''\}[e]} = F_{a'e} \begin{bmatrix} d & c \\ a'' & a \end{bmatrix} \sqrt{\frac{F_{1a''} \begin{bmatrix} d & a' \\ d & a' \end{bmatrix} F_{1a'} \begin{bmatrix} c & a \\ c & a \end{bmatrix}}{F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix} F_{1a''} \begin{bmatrix} e & a \\ e & a \end{bmatrix}}}. \quad (2.5)$$

The open defects ending on a fixed elementary conformal boundary condition a are closed under fusion. Curiously, as noted in [6], the corresponding fusion algebra is not given by the usual bulk fusion rule but depends on the boundary a . In our notation we can write the deformed fusion product as

$$\mathcal{D}_c^{[a]} \mathcal{D}_d^{[a]} = \sum_{e \in c \times d} N_{cd}^{[a]e} \mathcal{D}_e^{[a]} \quad (2.6)$$

where the coefficients $N_{cd}^{[a]e}$ can be obtained from figure 9 and formula (2.5) by specialising to the case $a = a' = a'' = b = b' = b''$. The corresponding expression can be recast into¹.

$$N_{cd}^{[a]e} = F_{da} \begin{bmatrix} e & a \\ c & a \end{bmatrix} F_{ca} \begin{bmatrix} e & a \\ d & a \end{bmatrix} \frac{F_{1e} \begin{bmatrix} c & d \\ c & d \end{bmatrix}}{F_{1a} \begin{bmatrix} e & a \\ e & a \end{bmatrix}}. \quad (2.7)$$

¹This expression is slightly more compact than the one following from (2.5) but is equivalent to it by virtue of F -matrices' identities as we show in the appendix A.

This expression is valid when the defect labelled by e can end topologically on the conformal boundary labelled by a . It may happen that e appears in the bulk fusion $c \times d$ but the defect labelled by e cannot end on a topologically. In this case $N_{cd}^{[a]e}$ vanishes. In general the coefficients $N_{cd}^{[a]e}$ are non-negative and symmetric under the interchange of c and d . The associativity of the open defect fusion was proven in [6]. Among other general properties of (2.6) and (2.7) we note the following identities

$$N_{1d}^{[a]d} = 1, \quad N_{dd}^{[a]1} = F_{11} \begin{bmatrix} d & d \\ d & d \end{bmatrix} = \frac{S_1^1}{S_1^d}, \quad \sum_{e \in c \times d} N_{cd}^{[a]e} = 1 \quad (2.8)$$

where S_j^i is the modular S -matrix.

3. Constraints on boundary RG flows

Suppose now that we take an elementary conformal boundary condition labelled by a and perturb it by a relevant operator $\psi(t)$ with a coupling λ . Let ψ_i stand for a complete basis of local boundary operators in the UV theory. A renormalised boundary correlator in the perturbed theory can be written as

$$\langle [\psi_{i_k}](x_k) \cdots [\psi_{i_1}](x_1) \rangle_\lambda = Z^{-1} \left\langle e^{-\lambda \int_{-\infty}^{\infty} \psi(t) dt - S_{ct}} [\psi_{i_k}](x_k) \cdots [\psi_{i_1}](x_1) \right\rangle \quad (3.1)$$

where S_{ct} stands for the counterterms action, $[\psi_i]$ denote renormalised boundary operators, λ is the renormalised coupling constant and Z is the normalisation factor:

$$Z = \left\langle e^{-\lambda \int_{-\infty}^{\infty} \psi(t) dt - S_c} \right\rangle. \quad (3.2)$$

More explicitly we have

$$S_{ct} = \int dt \sum_i M^i \psi_i(t), \quad [\psi_i](x) = \psi_i(x) + \sum_j M_i^j \psi_j(x) \quad (3.3)$$

where ψ_i stand for the (bare) operators at the UV fixed point and the terms with M^i and M_i^j stand for the coefficients of counterterms renormalising the action and the local operators respectively. For brevity we are not explicitly writing the dependence on a regulator but assume that the regulator is point splitting and the minimal subtraction scheme is employed.

Consider a correlation function at the UV fixed point in the presence of an open topological defect $\mathcal{D}_d^{[a]}$. It can be expressed in terms of correlators of local operators by sliding the defect to infinity along the boundary. Pulling the defect to the right and using the moves depicted on figure 8 we obtain

$$\begin{aligned} & \langle \psi_{i_k}(x_k) \cdots \psi_{i_{p+1}}(x_{p+1}) \mathcal{D}_d^{[a]}(s) \psi_{i_p}(x_p) \cdots \psi_{i_1}(x_1) \rangle \\ &= \langle (\hat{\mathcal{D}}_d \psi_{i_k})(x_k) \cdots (\hat{\mathcal{D}}_d \psi_{i_{p+1}})(x_{p+1}) \psi_{i_p} \cdots \psi_{i_1}(x_1) \rangle \end{aligned} \quad (3.4)$$

where $x_k > \cdots > x_{p+1} > s > x_p > \cdots > x_1$,

$$\begin{aligned}
 (\hat{\mathcal{D}}_d \psi_{i_{p+1}})(x_{p+1}) &= \sum_{\tilde{a} \in d \times a} \sum_j X_{i_{p+1}, \tilde{a} \tilde{a}}^{j, aa} \psi_j^{[a, \tilde{a}]}(x_{p+1}), \\
 (\hat{\mathcal{D}}_d \psi_{i_l})(x_l) &= \sum_{a', a'' \in d \times a} \sum_j X_{i_l, a' a''}^{j, aa} \psi_j^{[a', a'']}(x_l), \quad l = p + 1, \dots, k \quad (3.5)
 \end{aligned}$$

and the boundary condition \tilde{a} is assumed to appear between products of consecutive operators, like $\psi_j^{[a', \tilde{a}]}(x_{q+1}) \psi_l^{[\tilde{a}, a'']}(x_q)$, where the neighbouring boundary conditions match, while we get zero when they do not match. The coefficients $X_{l, a' a''}^{j, aa}$ represent the shrinking bubble of the defect $\mathcal{D}_d^{[a]}$ surrounding the operator $\psi_l^{[a, a']}$ with j labelling possible degeneracies of the Virasoro representation. For CFTs with minimal models type fusion² we have $X_{l, a' a''}^{j, aa} = \delta_l^j X_{l, a' a''}^{aa}$ where $X_{l, a' a''}^{aa}$ are given in (2.1).

Suppose now an open defect $\mathcal{D}_d^{[a]}$ commutes with ψ . We note that $\mathcal{D}_d^{[a]}$ also commutes with the operators that appear in the counterterms action S_{ct} . This follows from the fact that the counterterms are put in to subtract the divergences arising when several perturbing operators ψ collide. Such collisions can be represented by an operator product expansion of a group of operators:

$$\psi(t_n) \psi(t_{n-1}) \cdots \psi(t_1) = \sum_i C^i(t_1, t_2, \dots, t_n) \psi_i(t_1) \quad (3.6)$$

where $C^i(t_1, t_2, \dots, t_n)$ are some functions. Since $\mathcal{D}_d^{[a]}$ commutes with each operator $\psi(t_i)$ on the left-hand side, it commutes with each operator $\psi_i(t_1)$ on the right-hand side and thus with all operators appearing in S_{ct} . This means that an insertion of $\mathcal{D}_d^{[a]}$ into a perturbed correlation function (3.1) with the junction located at a point s can be moved freely inside the perturbed correlation functions as long as it does not pass through insertions of additional boundary operators, that is the correlation function

$$Z^{-1} \left\langle e^{-\lambda \int_{-\infty}^{\infty} \psi(t) dt - S_{ct}} [\psi_{i_k}](x_k) \cdots \mathcal{D}_d^{[a]}(s) \cdots [\psi_{i_1}](x_1) \right\rangle \quad (3.7)$$

is independent of s as long as it does not cross any of x_1, \dots, x_k . Moreover, passing through any of $[\psi_{i_j}](x_j)$ is given by exactly the same formulae (3.5) as in the UV theory (with the bare operators ψ_i replaced by $[\psi_i]$). To show this we note that the renormalised operators $[\psi_{i_j}]$ are given by the UV operators ψ_{i_j} plus counterterms. The latter are taken to cancel divergencies arising when some number of perturbing operators collide at the insertion point x_j . Again, the counterterm operators are contained in the operator product expansion of a group of operators containing the operators ψ and the UV operator ψ_{i_j} . Since $\mathcal{D}_d^{[a]}$ commutes with all ψ 's it acts on the counterterms in the same way as it acts on ψ_{i_j} . This means that an insertion of $\mathcal{D}_d^{[a]}$ into a perturbed theory correlator can be traded for a linear combination of renormalised local correlation functions with coefficients given by those of the UV theory. We should also note that besides the linear combinations (3.5) moving the defect also results in replacing the boundary condition a between the insertions by those arising in the fusion $d \times a$ of the defect with the UV boundary condition. Due to the Graham–Watts theorem at the end of the flow such

²For other chiral algebras these coefficients can be computed by a sequence of moves depicted in figure 7 but the answer will be different from (2.1). It would have to take into account possible degeneracies in Virasoro representations, different fusion vertices and a charge conjugation matrix.

segments have the conformal boundary condition given by the fusion of d with the infrared BCFT.

To summarise, the above means that $\mathcal{D}_d^{[a]}$ descends to a topological open defect in the perturbed theory and, consequently, at the infrared fixed point at the end of the flow. This places a constraint on the end point of the flow—the end point must be given by a conformal boundary condition that admits a topological junction with the defect labelled by d . Moreover, the action of $\mathcal{D}_d^{[a]}$ on the boundary operators of the perturbed theory is independent of the coupling λ —it is given by the action in the UV BCFT. Since all open defects that commute with ψ_i form a closed algebra under fusion, generated by elementary defects $\mathcal{D}_d^{[a]}$, the same fusion rules will be valid also in the deformed theory. Thus, in addition to admitting topological junctions with defects labelled by the same d 's, the corresponding open defects at the end point of the RG flow must form a subring under fusion that is isomorphic to that of the UV boundary condition. Given that in general the fusion algebra depends on the boundary condition this may place some additional constraints on the IR BCFT.

Consider next an open defect $\mathcal{D}_d^{[a]}$ that anti-commutes with ψ . Let us place the corresponding junction at a point s on the boundary and consider a perturbation with a coupling λ to the left of s and with a coupling $-\lambda$ to the right of s . A deformed correlation function in such a configuration can be written as

$$Z^{-1} \left\langle [\psi_{i_k}]_{-\lambda}(x_k) \cdots e^{\lambda \int_s^\infty \psi(\tau) d\tau - S_{\text{ct}}^+} \mathcal{D}_d^{[a]}(s) e^{-\lambda \int_{-\infty}^s \psi(\tau) d\tau - S_{\text{ct}}^-} \cdots [\psi_{i_1}]_\lambda(x_1) \right\rangle \tag{3.8}$$

where $x_k > \cdots > x_{p+1} > s > x_p > \cdots > x_1$, and

$$S_{\text{ct}}^- = \int_{-\infty}^s M^i(\lambda) \psi_i(\tau) d\tau, \quad S_{\text{ct}}^+ = \int_s^\infty M^i(-\lambda) \psi_i(\tau) d\tau, \tag{3.9}$$

that is S^- contains counterterms for the theory specified by λ and integrated to the left of the defect and S^+ contains the counterterms for the theory with the coupling $-\lambda$ integrated to the right of the defect. Furthermore, in (3.8) we have

$$[\psi_i]_\lambda(x) = \psi_i(x) + \sum_j M_i^j(\lambda) \psi_j(x), \quad [\psi_i]_{-\lambda}(x) = \psi_i(x) + \sum_j M_i^j(-\lambda) \psi_j(x) \tag{3.10}$$

so that the renormalised operators inserted to the left of the defect are defined with counterterm coefficients $M_i^j(\lambda)$ corresponding to the coupling λ while those inserted to the right have counterterms specified by $-\lambda$.

Since $\mathcal{D}_d^{[a]}$ anti-commutes with ψ it commutes with the counterterms that come from collisions of even numbers of ψ 's and anti-commutes with those coming from collisions of an odd number of ψ 's. This means that

$$\mathcal{D}_d^{[a]}(\tau + \epsilon) M^i(\lambda) \psi_i(\tau) = M^i(-\lambda) \psi_i(\tau) \mathcal{D}_d^{[a]}(\tau - \epsilon), \quad \epsilon > 0. \tag{3.11}$$

Hence the correlation function in (3.8) is independent of s as long as s does not cross any of the insertion points³ x_1, \dots, x_k . Moreover, for the same reasons as in the commuting case, when $\mathcal{D}_d^{[a]}$ is being passed from right to left through any of the insertions $[\psi_{i_j}](x_j)$ it acts on them via the UV theory coefficients (3.5) and changes the counterterms to those of the theory with the opposite coupling:

$$[\psi_i]_{-\lambda}(\tau)\mathcal{D}_d^{[a]}(\tau - \epsilon) = \mathcal{D}_d^{[a]}(\tau + \epsilon) \sum_{a', a'' \in d \times a} \sum_j X_{i_j, a' a''}^{j, aa} [\psi_j^{[a', a'']}]_{\lambda}(\tau), \quad \epsilon > 0 \quad (3.12)$$

and the boundary condition \tilde{a} is assumed to appear between the insertion and the new position of the defect. We finally comment on the normalisation factor Z in (3.8). It can be taken as in (3.2) to be given by the λ -deformed partition function however it is the same if we change in (3.2) λ to $-\lambda$ as we can insert $\mathcal{D}_d^{[a]}$ at minus infinity and move it through to plus infinity changing the sign of λ .

Taking λ to the infrared fixed point, it follows from the above that we get a topological junction of the defect labelled by d and the two conformal boundary conditions that describe the IR endpoints of the flow in the positive and negative λ directions. Thus, for each anti-commuting defect there must exist a topological junction between the two end-points of the flows in the positive and negative direction and the same bulk defect. If we take all open defects ending on a that either commute or anti-commute with ψ_i they form a \mathbb{Z}_2 -graded algebra with respect to fusion. Since the action (3.12) is independent of λ the corresponding fusion subring of the defects at and between the infrared fixed points must be isomorphic to the one at the UV theory.

There is one other interesting constraint arising from the presence of an anti-commuting defect: the g -factors of the two infrared fixed points must be the same. To explain why this is the case recall that the boundary entropy of the perturbed boundary condition with the coupling λ arises from the perturbed partition function on a disc (see e.g. [11] or [12]). With a point splitting regulator the value of the disc partition function remains the same if we insert into it an arc with two junctions of $\mathcal{D}_d^{[a]}$ between a pair of neighbouring insertions of ψ . We can then move one junction around the circle, anti-commuting with the insertions of ψ and counterterms, until it reaches the other junction at which point the arc can be removed. This implies that the disc partition function for the coupling λ is the same as the one with $-\lambda$ and hence the same goes for their boundary entropies and the g -factors in the infrared fixed points.

It should be noted that all of the above constraints generalise in a straightforward manner to the case when the UV boundary condition is a direct sum of elementary boundary conditions.

Before we finish this section we would like to comment briefly on the Hamiltonian description of the above situations. For simplicity we will not consider here the effects of possible divergences in the Hamiltonian formalism. Consider an infinite strip of width L with the boundary condition a put on both ends. Let $0 \leq \sigma \leq L$ be the coordinate across the strip and $-\infty < \tau < \infty$ be the coordinate along the strip. For τ being Euclidean time the Hilbert space can be decomposed into Virasoro irreducible representations V^i as

$$\mathcal{H}^{[a, a]} = \bigoplus_{i \in a \times a} V^i. \quad (3.13)$$

³This implies in particular that the counterterms in S^+ and S^- are sufficient to renormalise the theory with the defect inserted in s . In particular no additional counterterms are needed to be inserted at s . Such additional counterterms would be needed if no anti-commuting topological defect was inserted at s while perturbing with different couplings to the left and to the right of s .

Open defects that end topologically on both ends of the strip act on the states in $\mathcal{H}^{[a,a]}$ by the action of operator $\mathcal{D}_d^{[a]}$ described in the previous section. Consider next perturbing the boundary condition a on one or both ends of the strip by a relevant operator ψ . For a perturbation on the lower end the perturbed Hamiltonian acting on $\mathcal{H}^{[a,a]}$ can be written as

$$H_\lambda = \frac{\pi}{L} \left[L_0^{\text{UV}} - \frac{c}{24} + \lambda L\psi(0,0) \right] \tag{3.14}$$

where c is the central charge and L_0^{UV} is the dilation operator acting on $\mathcal{H}^{[a,a]}$. If an open defect $\mathcal{D}_d^{[a]}$ commutes with ψ then for any λ

$$\mathcal{D}_d^{[a]} H_\lambda = H_\lambda \mathcal{D}_d^{[a]} \tag{3.15}$$

and in particular at the IR fixed point $\mathcal{D}_d^{[a]}$ should commute with L_0 and thus $\mathcal{D}_d^{[a]}$ gives a symmetry of the infrared spectrum.

If $\mathcal{D}_d^{[a]}$ anti-commutes with ψ then we have

$$\mathcal{D}_d^{[a]} H_\lambda = H_{-\lambda} \mathcal{D}_d^{[a]}. \tag{3.16}$$

Taking λ to the fixed point (which is typically at infinity) we obtain

$$\mathcal{D}_d^{[a]} L_0^{\text{IR},1} = L_0^{\text{IR},2} \mathcal{D}_d^{[a]} \tag{3.17}$$

where $L_0^{\text{IR},1}$ and $L_0^{\text{IR},2}$ are the dilation operators for the IR endpoints corresponding to the negative and positive λ respectively. Thus, $\mathcal{D}_d^{[a]}$ intertwines the spectra of the two end-points.

4. Examples

4.1. Diagonal unitary minimal models

Here we remind the reader some basic facts about the unitary Virasoro minimal models with diagonal modular invariant. Such models are labelled by an integer m and have the central charge

$$c_m = 1 - \frac{6}{m(m+1)}. \tag{4.1}$$

The primary fields $\phi_{r,s}$ are labelled by two integers $1 \leq r \leq m-1$, $1 \leq s \leq m$ from the Kac table with the identification

$$\phi_{r,s} \equiv \phi_{m-r,m+1-s}. \tag{4.2}$$

The same set of integers label the bulk defects $\mathcal{D}_{r,s}$ as well as the elementary conformal boundary conditions which we will denote as (r,s) .

The fusion rules are summarised in the following equation

$$\phi_{r,s} \times \phi_{r',s'} = \sum_{p,q} \mathcal{N}_{(r,s),(r',s')}^{(p,q)} \phi_{p,q}, \quad \mathcal{N}_{(r,s),(r',s')}^{(p,q)} = \mathcal{N}_{r,r'}^p(m) \mathcal{N}_{s,s'}^q(m+1) \tag{4.3}$$

where

$$\mathcal{N}_{a,b}^c(m) = \begin{cases} 1 & \text{if } |a-b| + 1 \leq c \leq \min(a+b-1, 2m-a-b-1) \\ & \text{and } a+b+c \text{ is odd} \\ 0 & \text{otherwise} \end{cases} \tag{4.4}$$

The fusion ring contains two subrings generated by fields of the form $\phi_{1,s}$ and $\phi_{r,1}$ respectively. The two subrings intersect over a subring generated by the identity field and the operator $\phi_{1,m} \equiv \phi_{m-1,1}$. The bulk defects satisfy the same fusion rule. The defect

$$\mathcal{S} \equiv \mathcal{D}_{1,m} \equiv \mathcal{D}_{m-1,1} \tag{4.5}$$

describes the spin reversal symmetry. It satisfies the group property

$$\mathcal{S} \circ \mathcal{S} = \mathcal{D}_{1,1}, \tag{4.6}$$

fuses with the other defects according to

$$\mathcal{S} \circ \mathcal{D}_{r,s} = \mathcal{D}_{r,m+1-s} \equiv \mathcal{D}_{m-r,s} \tag{4.7}$$

and acts on Cardy boundary conditions as

$$\mathcal{S} \cdot (r, s) = (r, m + 1 - s) \equiv (m - r, s). \tag{4.8}$$

The spin reversal invariant Cardy boundary conditions are of the form $(\frac{m}{2}, s)$, $s = 1, \dots, \frac{m}{2}$ if m is even and of the form $(r, \frac{m+1}{2})$, $r = 1, \dots, \frac{m+1}{2}$ if m is odd. For such boundary conditions we can introduce the \mathcal{S} -charge for the boundary fields that according to (2.2) is given by

$$\mathcal{S}\psi_i^{[a,a]} = F_{aa} \begin{bmatrix} (m-1, 1) & a \\ a & i \end{bmatrix} \psi_i^{[a,a]} = \pm \psi_i^{[a,a]} \tag{4.9}$$

where the boundary label a is $(\frac{m}{2}, s)$ or $(r, \frac{m+1}{2})$ depending on the parity of m and $i \in a \times a$. This charge is equal to ± 1 due to the orthogonality relation (A.6) and the fusion rule $(m-1, 1) \times a = a$.

If we are perturbing an \mathcal{S} -invariant boundary condition by a charge 1 boundary field then, by virtue of the Graham–Watts theorem, we expect each end point of the flow to be \mathcal{S} -invariant. If we perturb by a charge -1 field then the end points of the flow are interchanged by the action of \mathcal{S} . For example in the tricritical Ising model, that corresponds to $m = 4$, we have two spin reversal invariant Cardy boundary conditions: $(2, 2)$ and $(2, 1)$. The latter boundary condition is stable while the $(2, 2)$ boundary condition, also known as the disordered boundary condition, admits two relevant boundary fields: $\psi_{1,2}$ and $\psi_{1,3}$. The first field has the \mathcal{S} -charge -1 while the second one has charge 1. The boundary RG flows in the tricritical Ising model that start from the elementary boundary conditions were described in [19]. Both $\psi_{1,2}$ and $\psi_{1,3}$ perturbations of the disordered boundary condition are integrable and their end points are given on the following diagrams:

$$\begin{aligned} (2, 1) &\xleftarrow{\psi_{1,3}} (2, 2) \xrightarrow{-\psi_{1,3}} (1, 1) \oplus (3, 1), \\ (1, 1) &\xleftarrow{\psi_{1,2}} (2, 2) \xrightarrow{-\psi_{1,2}} (3, 1). \end{aligned} \tag{4.10}$$

It is straightforward to check that the endpoints satisfy the requirements for the action of \mathcal{S} .

Below we will be particularly interested in boundary flows triggered by perturbing the boundary condition (r, s) by the boundary field $\psi_{1,3}^{[rs,rs]}$. For large values of m these flows were studied in [9] where the end points were identified using the g -theorem. The end points in the non-perturbative regime were found in [5] with the help of Graham–Watts theorem, which was put forward in that paper, and using the results of [8]. The general rule for the end points of the

Table 1. Open defects on boundary conditions in tetracritical Ising model.

b.c.	Defects commuting with $\psi_{1,3}$	Defects anti-commuting with $\psi_{1,3}$
(1, 3)	$\mathcal{D}_{1,1}^{[1,3]}$	$\mathcal{S}^{[1,3]}$
(3, 3)	$\mathcal{D}_{1,1}^{[3,3]}, \mathcal{D}_{3,1}^{[3,3]}$	$\mathcal{S}^{[3,3]}, \mathcal{D}_{2,1}^{[3,3]}$
(2, 2)	$\mathcal{D}_{1,1}^{[2,2]}, \mathcal{D}_{3,1}^{[2,2]}$	None
(3, 2)	$\mathcal{D}_{1,1}^{[3,2]}, \mathcal{D}_{3,1}^{[3,2]}$	None

$\psi_{1,3}$ flows that start from elementary boundary conditions can be summarised in the following two expressions

$$(r, s) \longrightarrow \bigoplus_{i=1}^{\min(r,s,m-r,m-s)} (|r-s| + 2i - 1, 1), \tag{4.11}$$

$$(r, s) \longrightarrow \bigoplus_{i=1}^{\min(r,s-1,m-r,m-s+1)} (|r-s+1| + 2i - 1, 1) \tag{4.12}$$

where one expression corresponds to a positive choice of the coupling and the other to the negative choice. To the best of our knowledge it has not been fixed in general which answer corresponds to which sign. The expressions (4.11) and (4.12) are interchanged under the action of the field identification (4.2).

Commutators of boundary fields with open defects can be computed using the general expression (2.1). The fusion matrices for the diagonal minimal models can be calculated recursively following [15] (see also [16] for a closed expression).

4.2. Tetracritical Ising model

The first example of a non-trivial open defect that is different from \mathcal{S} and commutes with a relevant operator on an elementary boundary condition appears in the tetracritical Ising model that is the unitary minimal model with $m = 5$. This model has 10 primary fields and thus the same number of topological defects and elementary conformal boundary conditions. We focus on the $\psi_{1,3}$ boundary field where we know the end points of the flows. All elementary boundary conditions have a $\psi_{1,3}$ boundary field except for the 4 boundary conditions of the form $(r, 1)$, $1 \leq r \leq 4$. Table 1 shows the open defects that have a topological junction with a given boundary condition and that commute or anti-commute with $\psi_{1,3}$.

We note that (1, 3) and (3, 3) boundary conditions are stable under fusion with \mathcal{S} and thus are spin reversal symmetric while (2, 2) and (3, 2) form a doublet. The two boundary conditions omitted from the table: (1, 2), (1, 4), have no non-trivial defects commuting or anti-commuting with $\psi_{1,3}$.

In view of the general discussion in section 3 for the $\psi_{1,3}$ flows that start from (3, 3), (2, 2) or (3, 2) boundary condition the end points must admit a topological junction with $\mathcal{D}_{3,1}$. Examining the fusion rules we find that this implies that they must contain one of the following 5 elementary boundary conditions: (3, 1), (2, 1), (3, 3), (2, 2), (3, 2).

Moreover, for the flows from (3, 3) and (1, 3) the end points are exchangeable by the fusion with \mathcal{S} . Also for the flows from (3, 3) there is a topological junction between the two end points and $\mathcal{D}_{2,1}$. These conditions become even more restrictive if we add constraints from the g -theorem. The end points of the flows given by (4.11) and (4.12) for the flows at hand are

$$(2, 1) \longleftarrow (1, 3) \longrightarrow (3, 1) \tag{4.13}$$

$$(4, 1) \oplus (2, 1) \longleftarrow (3, 3) \longrightarrow (1, 1) \oplus (3, 1) \tag{4.14}$$

$$(1, 1) \oplus (3, 1) \longleftarrow (2, 2) \longrightarrow (2, 1) \tag{4.15}$$

$$(4, 1) \oplus (2, 1) \longleftarrow (3, 2) \longrightarrow (3, 1). \tag{4.16}$$

We check that these flows satisfy all of the constraints following from table 1.

It is interesting to calculate the boundary fusion rings formed by the defects in table 1. The (2, 2), (3, 3) and (3, 2) boundary conditions have the open defect ring consisting of defects commuting with $\psi_{1,3}$ generated by $\mathcal{D}_{3,1}^{[a]}$ with a single relation given by

$$\mathcal{D}_{3,1}^{[a]} \circ \mathcal{D}_{3,1}^{[a]} = f\mathcal{D}_{1,1}^{[a]} + (1 - f)\mathcal{D}_{3,1}^{[a]}, \quad f = \frac{1}{2}(\sqrt{5} - 1). \tag{4.17}$$

In fact (4.17) holds for any boundary condition a admitting a topological junction with $\mathcal{D}_{3,1}$. This fact is a simple consequence of the bulk fusion rule and the general identities (2.8). Thus, if the end point of a $\psi_{1,3}$ flow contains an elementary boundary condition admitting a topological junction with $\mathcal{D}_{3,1}$ then it will satisfy the same composition rule (4.17) as in the UV boundary condition.

The (3, 3) boundary condition has additional open defects that anti-commute with $\psi_{1,3}$ that satisfy the following relations under fusion

$$\mathcal{D}_{2,1}^{[3,3]} \circ \mathcal{D}_{3,1}^{[3,3]} = f\mathcal{S}^{[3,3]} + (1 - f)\mathcal{D}_{2,1}^{[3,3]}, \tag{4.18}$$

$$\mathcal{D}_{2,1}^{[3,3]} \circ \mathcal{D}_{2,1}^{[3,3]} = f\mathcal{D}_{1,1}^{[3,3]} + (1 - f)\mathcal{D}_{3,1}^{[3,3]} \tag{4.19}$$

and $\mathcal{S}^{[3,3]}$ fuses with the other open defects according to the bulk fusion rule.

The (3, 3) boundary condition also has a boundary $\psi_{2,1}$ field which is relevant. To the best of our knowledge these flows have not been investigated before and the end points have not been identified. We find that this perturbation commutes with $\mathcal{S}^{[3,3]}$ and anti-commutes with $\mathcal{D}_{1,3}^{[3,3]}$. The commutation with the spin reversal defect implies that each of the end points must be invariant under the spin reversal. Together with the constraints from the g -theorem this gives us two possible infrared end points: (1, 3) and $(1, 1) \oplus (1, 5)$. The existence of a junction with $\mathcal{D}_{1,3}$ gives us two possible pairs of fixed points: either they are both (1, 3) or one of them is (1, 3) and the other is $(1, 1) \oplus (1, 5)$. Interestingly the condition on the g -factors being the same is satisfied for the second pair due to the identity $g_{1,3} = 2g_{1,1} = 2g_{1,5}$. Moreover, we calculate the UV fusion of the anti-commuting defect to be given by

$$\mathcal{D}_{1,3}^{[3,3]} \circ \mathcal{D}_{1,3}^{[3,3]} = \frac{1}{2}\mathcal{D}_{1,1}^{[3,3]} + \frac{1}{2}\mathcal{S}^{[3,3]}. \tag{4.20}$$

The same fusion rule must be satisfied by the $\mathcal{D}_{1,3}$ defect between the two infrared fixed points. It is straightforward to check that $\mathcal{D}_{1,3}^{[1,3]}$ satisfies the same rule. For the second pair we find that there is a unique combination

$$\mathcal{D}_{13}^{\text{IR}} = \frac{1}{\sqrt{2}}(\mathcal{D}_{1,3}^{13,11} + \mathcal{D}_{1,3}^{13,15}) \tag{4.21}$$

Table 2. Open defects on boundary conditions in pentacritical Ising model.

b.c.	Defects commuting with $\psi_{1,3}$	Defects anti-commuting with $\psi_{1,3}$
(3, 3)	$\mathcal{D}_{1,1}^{[3,3]}, \mathcal{D}_{3,1}^{[3,3]}, \mathcal{S}^{[3,3]}$	None
(2, 2)	$\mathcal{D}_{1,1}^{[2,2]}, \mathcal{D}_{3,1}^{[2,2]}$	None
(2, 3)	$\mathcal{D}_{1,1}^{[2,3]}, \mathcal{D}_{3,1}^{[2,3]}$	None
(3, 2)	$\mathcal{D}_{1,1}^{[3,2]}, \mathcal{D}_{3,1}^{[3,2]}, \mathcal{S}^{[3,2]}$	None

that satisfies⁴

$$\begin{aligned} \mathcal{D}_{13}^{\text{IR}} \circ \mathcal{D}_{13}^{\text{IR}\dagger} &= \frac{1}{2}(\mathcal{D}_{1,1}^{[1,3]} + \mathcal{S}^{[1,3]}), \\ \mathcal{D}_{13}^{\text{IR}\dagger} \circ \mathcal{D}_{13}^{\text{IR}} &= \frac{1}{2}[(\mathcal{D}_{1,1}^{[1,1]} + \mathcal{D}_{1,1}^{[1,5]}) + (\mathcal{S}^{[11,15]} + \mathcal{S}^{[15,11])}] \end{aligned} \quad (4.22)$$

where $\mathcal{D}_{13}^{\text{IR}\dagger} = (\mathcal{D}_{1,3}^{11,13} + \mathcal{D}_{1,3}^{15,13})/\sqrt{2}$ is the conjugate defect. Thus, all constraints from the commuting and anti-commuting open defects are satisfied by each of the two pairs. We did check numerically⁵, using the truncated boundary conformal space approach of [20], that for positive λ the flow at hand ends up at (1, 3) while for negative λ it flows to $(1, 1) \oplus (1, 5)$.

4.3. Pentacritical Ising model

Pentacritical Ising model corresponds to the minimal model with $m = 6$. This model has 15 primary states and the same number of topological defects and conformal boundary conditions. Up to the action of the spin reversal generator we have 6 representatives of elementary boundary conditions admitting a boundary $\psi_{1,3}$ field: (1, 2), (1, 3), (2, 2), (3, 3), (2, 3), (3, 2). The boundary conditions (3, 3) and (3, 2) are spin-reversal invariant. The elementary boundary conditions that have non-trivial open defects commuting or anti commuting with $\psi_{1,3}$ are tabulated in table 2. We see that the end points of the flows that start with the 4 boundary conditions in table 2 (and their spin reverses) must admit a topological junction with $\mathcal{D}_{3,1}$. For the spin reversal invariant boundary conditions: (3, 3), (3, 2), the end points must be also spin-reversal invariant. Noting that the g -factors satisfy $g_{3,3} > g_{3,2} > g_{3,1}$ we see that each end point of the $\psi_{1,3}$ flows from (3, 2) is either degenerate or is given by the (3, 1) boundary condition that is spin-reversal invariant.

The expressions (4.11) and (4.12) give the flows

$$(1, 1) \oplus (3, 1) \longleftarrow (2, 3) \longrightarrow (2, 1) \oplus (4, 1) \quad (4.23)$$

$$(4, 1) \oplus (2, 1) \longleftarrow (3, 3) \longrightarrow (1, 1) \oplus (3, 1) \oplus (5, 1) \quad (4.24)$$

$$(1, 1) \oplus (3, 1) \longleftarrow (2, 2) \longrightarrow (2, 1) \quad (4.25)$$

$$(4, 1) \oplus (2, 1) \longleftarrow (3, 2) \longrightarrow (3, 1). \quad (4.26)$$

It is straightforward to check that these flows satisfy the above constraints.

⁴In checking these relations it is important to allow a' and b' on figure 9 each to take the values (1, 1) and (1, 5) independently of each other.

⁵The author found an analytic argument based on RG interfaces that excludes the possibility that both end points are (1, 3), but this is outside the scope of the present paper and will be reported elsewhere.

It is interesting to take a look at the fusion rings of the open defects in table 2. Noting that the bulk fusion rule

$$(3, 1) \times (3, 1) = (1, 1) + (3, 1) + (5, 1) \tag{4.27}$$

contains 3 terms, the boundary fusion rule (2.6) now has room for different deformations. Indeed, we find

$$\mathcal{D}_{3,1}^{[a]} \circ \mathcal{D}_{3,1}^{[a]} = \frac{1}{2} \mathcal{D}_{1,1}^{[a]} + \frac{1}{2} \mathcal{D}_{3,1}^{[a]} \tag{4.28}$$

for $a = (2, 2), (4, 4), (4, 6), (2, 6)$ and

$$\mathcal{D}_{3,1}^{[b]} \circ \mathcal{D}_{3,1}^{[b]} = \frac{1}{2} \mathcal{D}_{1,1}^{[b]} + \frac{1}{2} \mathcal{S}^{[b]} \tag{4.29}$$

for $b = (3, 3), (3, 2), (3, 1)$. The $\mathcal{S}^{[b]}$ generator fuses according to the bulk fusion rule. It is interesting to note that the boundary fusion rule (4.29) means that $\mathcal{D}_{3,1}^{[b]}$ is an open duality defect in the sense of [2, 3], that is its fusion with itself contains only group-like open defects.

We also have a boundary field $\psi_{1,2}$ present on the boundary condition (3, 3). The corresponding boundary flows are believed to be integrable but, as in the case of $\psi_{2,1}$ perturbation in the tetracritical model, have not been investigated before. We find that $\mathcal{D}_{3,1}^{[3,3]}$ anti-commutes and $\mathcal{S}^{[3,3]}$ commutes with $\psi_{1,2}$ that makes this case quite similar to the case of $\psi_{2,1}$ perturbation considered at the end of the previous section. It is instructive to see how all of the consequences considered in section 3 can be combined with the constraints from the g -theorem to restrict the choices of the infrared fixed points. We can first list all spin reversal invariant boundary conditions with a g -factor lower than that of the UV value. This gives us two singlets: (3, 2), (3, 1); four doublets: $A = (3, 1) \oplus (3, 1)$, $B = (1, 1) \oplus (5, 1)$, $C = (5, 5) \oplus (1, 5)$, $D = (4, 6) \oplus (2, 6)$; one triplet: $(1, 1) \oplus (5, 1) \oplus (3, 1)$; and one quadruplet: $(1, 1) \oplus (1, 1) \oplus (5, 1) \oplus (5, 1)$. The condition that the g -factors of both IR end points must be equal implies that either both end points are the same (and belong to the above list), or form one of the following 3 pairs:

$$\begin{aligned} &(3, 1) \text{ and } (1, 1) \oplus (5, 1), \quad (3, 1) \oplus (3, 1) \text{ and } (1, 1) \oplus (5, 1) \oplus (3, 1) \\ &(3, 1) \oplus (3, 1) \text{ and } (1, 1) \oplus (1, 1) \oplus (5, 1) \oplus (5, 1) \end{aligned}$$

that are permitted because of the identity: $g_{3,1} = 2g_{1,1} = 2g_{5,1}$. Adding the condition that there must be a topological junction with $\mathcal{D}_{3,1}$ possible between the two end points discards two pairs with equal boundary conditions: B, B and C, C . Finally, requiring that the junction with $\mathcal{D}_{3,1}$ should satisfy the fusion product given in (4.28) we can discard one more pair: D, D . This follows from checking that no combination of 3 available junctions: $\mathcal{D}_{3,1}^{[4,6]}$, $\mathcal{D}_{3,1}^{[2,6]}$, $\mathcal{D}_{3,1}^{[46,26]}$ can be chosen to satisfy (4.28). All of these constraints leave us in the end with 7 distinct pairs of possible infrared fixed points. Thus, in this example we see that *each* of the constraints we derived in section 3 reduces the number of possibilities. We finish this example by reporting that truncated conformal space approach numerics gives the spectra that match with the (3, 1) boundary condition for large positive λ and that of $(1, 1) \oplus (5, 1)$ boundary condition for large negative λ .

5. Flows from direct sums of boundary conditions

So far we have discussed the implication of open defects commuting with the perturbation for the flows originating from an elementary boundary condition. This can be generalised

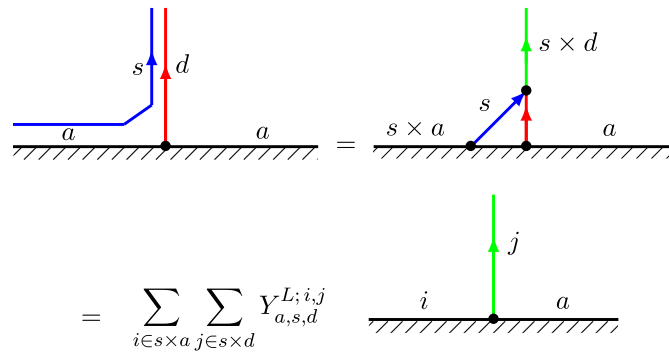


Figure 10. Fusion of a closed defect with a junction on the left.

to flows from direct sums of elementary boundary conditions triggered by boundary condition changing operators. Examples of such flows, including those triggered by $\psi_{1,3}$ operators, were studied in [21]. The commutator with an open defect has been calculated on figure 8.

A new feature of direct sum boundary conditions is that the perturbing field can be a linear combination of several components and a commuting (or anti-commuting) open defect can be a particular linear combination of defects with the same Virasoro label but linking different sets of elementary boundary conditions. One way to generate such linear combinations is by starting with a commuting open defect on an elementary boundary condition and fuse it with a closed defect. Suppose $\mathcal{D}_d^{[a]}$ commutes with $\psi_i^{[a,a]}$. We can fuse the junction with a closed string defect \mathcal{D}_s on either side of the junction. Such a fusion done on the left is illustrated on figure 10, where the coefficients $Y_{a,s,d}^{L;i,j}$ are easily computed using the results of [6] (see figure 8 of [6] in particular). As the final configuration on figure 10 is only an intermediate result, we omit the explicit expression for $Y_{a,s,d}^{L;i,j}$. At this stage it is important for us to note that, as a consequence of the commutation of $\mathcal{D}_d^{[a]}$ with $\psi_i^{[a,a]}$, the open defect combination

$$\mathcal{D}_j^L \equiv \sum_{i \in s \times a} Y_{a,s,d}^{L;i,j} \mathcal{D}_j^{[i,a]} \tag{5.1}$$

satisfies

$$\mathcal{D}_j^L \psi_i^{[a,a]} = \mathcal{D}_s(\psi_i^{[a,a]}) \mathcal{D}_j^L \tag{5.2}$$

where

$$\mathcal{D}_s(\psi_i^{[a,a]}) = \sum_{n,m \in s \times a} X_{i,nm}^{aa} \psi_i^{[n,m]} \tag{5.3}$$

is the action on the boundary field $\psi_i^{[a,a]}$ of the fusion of \mathcal{D}_s with the boundary a . Note that (5.2) is true for any fixed label j . Now, picking a configuration given by (5.1) with a fixed label j we can further fuse it with the closed defect \mathcal{D}_s on the right side of the junctions. Using steps similar to those on figure 10 we arrive at the following open defect combinations

$$\mathcal{D}_i^{j,LR} = \sum_{n,m \in s \times a} Z_{j(nml)}^{asd} \mathcal{D}_i^{[n,m]} \tag{5.4}$$

that, due to the associativity of the fusion operations, commute with the fused boundary field (5.3) for each choice of the label $l \in s \times s \times d$ and $j \in s \times d$. The coefficients $Z_{j(nml)}^{asd}$ are calculated using the results of [6]:

$$Z_{j(nml)}^{asd} = F_{aj} \begin{bmatrix} a & n \\ d & s \end{bmatrix} F_{al} \begin{bmatrix} m & n \\ s & j \end{bmatrix} \sqrt{\frac{F_{1n} \begin{bmatrix} s & a \\ s & a \end{bmatrix} F_{1m} \begin{bmatrix} s & a \\ s & a \end{bmatrix}}{F_{1n} \begin{bmatrix} l & m \\ l & m \end{bmatrix}}} \mathcal{N}_{jl}^{asd} \quad (5.5)$$

where

$$\mathcal{N}_{jl}^{asd} = \sqrt{\frac{F_{1a} \begin{bmatrix} d & a \\ d & a \end{bmatrix}}{F_{1j} \begin{bmatrix} s & d \\ s & d \end{bmatrix} F_{1l} \begin{bmatrix} s & j \\ s & j \end{bmatrix}}} \quad (5.6)$$

is an overall normalisation factor. Similarly, we can do the above fusion in the reversed order, that is first fusing with a closed defect on the right, singling out an elementary component labelled by j , then fusing it on the left and singling out open defects labelled by l . The resulting open defects are given by a combination

$$\mathcal{D}_l^{j,RL} = \sum_{n,m} \tilde{Z}_{asd}^{j(nml)} \mathcal{D}_l^{[n,m]} \quad (5.7)$$

where

$$\tilde{Z}_{asd}^{j(nml)} = Z_{asd}^{j(nml)} \sqrt{\frac{F_{1n} \begin{bmatrix} l & m \\ l & m \end{bmatrix}}{F_{1m} \begin{bmatrix} l & n \\ l & n \end{bmatrix}}} \quad (5.8)$$

These open defects also commute with the fused boundary field $\mathcal{D}_s(\psi_i^{[a,a]})$.

We illustrate the constructions (5.3), (5.4), (5.7) on a couple of explicit examples. Consider the Cardy boundary condition (2, 2) in the tetracritical Ising model. The open defect $\mathcal{D}_{31}^{[2,2]}$ commutes with $\psi_{13}^{[22,22]}$. Fusing the boundary with a closed defect $\mathcal{D}_{1,2}$ we obtain the direct sum $(2, 1) \oplus (2, 3)$. Up to an overall factor the boundary field $\psi_{13}^{[22,22]}$ is mapped to the combination

$$\Psi \equiv \tilde{\psi}_{13}^{[23,23]} - 2(\tilde{\psi}_{13}^{[23,21]} + \tilde{\psi}_{13}^{[21,23]}) \quad (5.9)$$

where we use the notation

$$\tilde{\psi}_i^{[a,b]} = \frac{1}{\alpha_i^{ab}} \psi_i^{[a,b]} \quad (5.10)$$

and the boundary fields $\psi_i^{[a,b]}$ are normalised as in [6]. Since $(1, 2) \times (3, 1) = (3, 2)$ we have only one value $j = (3, 2)$ in (5.4), (5.7). Using (5.4), (5.7) we find the open defects combinations

$$\mathcal{D}_{3,3}^{[23,23]} - \mathcal{D}_{3,3}^{[23,21]} - \sqrt{2} \mathcal{D}_{3,3}^{[21,23]}, \quad (5.11)$$

$$\mathcal{D}_{3,1}^{[23,23]} - \mathcal{D}_{3,1}^{[21,21]} \quad (5.12)$$

each of which commutes with (5.9) as can be checked directly. Using (5.7) gives the same combinations. These combinations are fixed by the commutation condition up to an overall factor.

Our second example starts with the same triple: $a = (2, 2)$, $d = (3, 1)$, $\psi_{1,3}^{[22,22]}$, but this time we fuse it with $\mathcal{D}_{2,1}$. This gives the direct sum of boundary conditions: $(1, 2) \oplus (3, 2)$ with a boundary field

$$\Psi' \equiv 7\tilde{\psi}_{1,3}^{[32,32]} - 9\tilde{\psi}_{1,3}^{[12,12]}. \quad (5.13)$$

(Again for brevity we dropped the overall normalisation factor.) We now have two choices for j in (5.4), (5.7): $j \in \{(2, 1), (4, 1)\}$. This gives us 4 particular linear combinations of the following three elementary open defects:

$$\mathcal{D}_{3,1}^{[32,32]}, \quad \mathcal{D}_{3,1}^{[32,12]}, \quad \mathcal{D}_{3,1}^{[12,32]}. \quad (5.14)$$

The coefficients of these linear combinations are quite ugly so we do not present them here, but we checked that they span the linear subspace generated by the elementary open defects listed in (5.14). Indeed, a separate calculation shows that each of the defects in (5.14) commutes with Ψ' .

6. Concluding remarks

Our considerations in section 3 did not depend on any particular choice of a rational CFT. Given an open topological defect on a conformal boundary that either commutes or anti-commutes with a relevant boundary perturbation all the consequences for RG flows derived in that section would apply. By working out a number of explicit examples in the minimal models we showed that all of these constraints can be used to restrict the possible infrared fixed points in RG flows, in particular in situations in which no other analytic arguments are known that would give the same restrictions.

It would be interesting to generalise the calculations done in [6] to a more general chiral algebra and to find other examples of applications of our general results to boundary RG flows in other models. More systematically, one can try to obtain some general results towards classifying all possible pairs—a relevant boundary operator plus a commuting or anti-commuting open topological defect, in given RCFTs or classes of RCFTs. Certainly such situations, when such a pair exists, are special. As we discussed in section 3, boundary operators in such pairs generate a subalgebra under OPE. In the bulk CFT perturbations with this property the Hamiltonian is block diagonal that signals the presence of additional conserved charges. Moreover, like the $\Phi_{1,3}$, $\Phi_{1,2}$ and $\Phi_{2,1}$ bulk perturbations of minimal models (see [22]), such perturbations are known to give integrable models. The integrability aspect of boundary perturbations is still comparatively less studied, particularly for the $\psi_{1,2}$ and $\psi_{2,1}$ perturbations. It would be interesting to investigate possible connections between the presence of commuting or anti-commuting defects and integrability, perhaps one could try to exploit the link between defects and integrability established for bulk perturbations in [14]. We hope to address these questions in future work.

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Appendix A. Some identities for the minimal model fusion matrices

$$F_{pq} \begin{bmatrix} a & b \\ c & d \end{bmatrix} = F_{pq} \begin{bmatrix} c & d \\ a & b \end{bmatrix} = F_{pq} \begin{bmatrix} b & a \\ d & c \end{bmatrix} \tag{A.1}$$

$$F_{11} \begin{bmatrix} a & a \\ a & a \end{bmatrix} = \frac{S_{11}}{S_{1a}} \tag{A.2}$$

where S_{ab} stand for the elements of the modular S -matrix.

$$F_{1a} \begin{bmatrix} b & c \\ b & c \end{bmatrix} F_{a1} \begin{bmatrix} b & b \\ c & c \end{bmatrix} = \frac{S_{11}S_{1a}}{S_{1b}S_{1c}} \tag{A.3}$$

$$F_{a1} \begin{bmatrix} b & b \\ c & c \end{bmatrix} = \frac{S_{1a}}{S_{1c}} F_{c1} \begin{bmatrix} a & a \\ b & b \end{bmatrix} \tag{A.4}$$

$$F_{e1} \begin{bmatrix} a & a \\ d & d \end{bmatrix} F_{fa} \begin{bmatrix} b & c \\ e & d \end{bmatrix} = F_{f1} \begin{bmatrix} c & c \\ d & d \end{bmatrix} F_{ec} \begin{bmatrix} a & b \\ d & f \end{bmatrix} \tag{A.5}$$

$$\sum_s F_{ps} \begin{bmatrix} b & c \\ a & d \end{bmatrix} F_{sr} \begin{bmatrix} c & d \\ b & a \end{bmatrix} = \delta_{pr} \tag{A.6}$$

We next show how one can use the above identities to establish the equivalence of the expression for the coefficients $N_{cd}^{[a]e}$ that follows from (2.5) and the expression presented in formula (2.7). Formula (2.5) gives

$$N_{cd}^{[a]e} = \frac{F_{1a} \begin{bmatrix} d & a \\ d & a \end{bmatrix} F_{1a} \begin{bmatrix} c & a \\ c & a \end{bmatrix}}{F_{1a} \begin{bmatrix} e & a \\ e & a \end{bmatrix} F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix}} F_{ae}^2 \begin{bmatrix} d & c \\ a & a \end{bmatrix}. \tag{A.7}$$

To show that this equals the expression in (2.7) we first use the identities

$$F_{ae} \begin{bmatrix} d & c \\ a & a \end{bmatrix} = \frac{F_{da} \begin{bmatrix} e & a \\ c & a \end{bmatrix} F_{a1} \begin{bmatrix} a & a \\ c & c \end{bmatrix}}{F_{d1} \begin{bmatrix} e & e \\ c & c \end{bmatrix}}, \tag{A.8}$$

$$F_{ae} \begin{bmatrix} d & c \\ a & a \end{bmatrix} = \frac{F_{ca} \begin{bmatrix} e & a \\ d & a \end{bmatrix} F_{a1} \begin{bmatrix} a & a \\ d & d \end{bmatrix}}{F_{c1} \begin{bmatrix} e & e \\ d & d \end{bmatrix}} \tag{A.9}$$

that are particular instances of (A.5). Substituting each of these identities into (A.7) we obtain

$$N_{cd}^{[a]e} = \frac{F_{da} \begin{bmatrix} e & a \\ c & a \end{bmatrix} F_{ca} \begin{bmatrix} e & a \\ d & a \end{bmatrix}}{F_{1a} \begin{bmatrix} e & a \\ e & a \end{bmatrix}} \tilde{N}_{cd}^{[a]e} \tag{A.10}$$

where

$$\tilde{N}_{cd}^{[a]e} = \frac{F_{1a} \begin{bmatrix} d & a \\ d & a \end{bmatrix} F_{a1} \begin{bmatrix} a & a \\ d & d \end{bmatrix} F_{1a} \begin{bmatrix} c & a \\ c & a \end{bmatrix} F_{a1} \begin{bmatrix} a & a \\ c & c \end{bmatrix}}{F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix} F_{d1} \begin{bmatrix} e & e \\ c & c \end{bmatrix} F_{c1} \begin{bmatrix} e & e \\ d & d \end{bmatrix}} \quad (\text{A.11})$$

Using (A.3) in the numerator of (A.11) we rewrite the last expression as

$$\tilde{N}_{cd}^{[a]e} = \frac{S_{11}^2}{S_{1d}S_{1c}F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix} F_{d1} \begin{bmatrix} e & e \\ c & c \end{bmatrix} F_{c1} \begin{bmatrix} e & e \\ d & d \end{bmatrix}} \quad (\text{A.12})$$

Finally we use the two identities

$$F_{d1} \begin{bmatrix} e & e \\ c & c \end{bmatrix} = \frac{S_{11}}{S_{1c}F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix}}, \quad F_{c1} \begin{bmatrix} e & e \\ d & d \end{bmatrix} = \frac{S_{11}}{S_{1d}F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix}} \quad (\text{A.13})$$

that follow from a combination of (A.3) and (A.4). Substituting (A.13) into (A.12) we obtain

$$\tilde{N}_{cd}^{[a]e} = F_{1e} \begin{bmatrix} d & c \\ d & c \end{bmatrix} \quad (\text{A.14})$$

that being combined with (A.10) gives (2.7).

Alternatively, formula (2.7) can be obtained independently by using a sequence of moves on the topological defects involved, different from the ones used in [6].

ORCID iDs

Anatoly Konechny  <https://orcid.org/0000-0003-3261-8262>

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