



Heriot-Watt University  
Research Gateway

## Porosity gradients as a means of driving lateral flows at cometary surfaces

### Citation for published version:

Christou, C, Dadzie, SK, Marschall, R & Thomas, N 2020, 'Porosity gradients as a means of driving lateral flows at cometary surfaces', *Planetary and Space Science*, vol. 180, 104752.  
<https://doi.org/10.1016/j.pss.2019.104752>

### Digital Object Identifier (DOI):

[10.1016/j.pss.2019.104752](https://doi.org/10.1016/j.pss.2019.104752)

### Link:

[Link to publication record in Heriot-Watt Research Portal](#)

### Document Version:

Peer reviewed version

### Published In:

Planetary and Space Science

### Publisher Rights Statement:

© 2019 Elsevier Ltd. All rights reserved.

### General rights

Copyright for the publications made accessible via Heriot-Watt Research Portal is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

### Take down policy

Heriot-Watt University has made every reasonable effort to ensure that the content in Heriot-Watt Research Portal complies with UK legislation. If you believe that the public display of this file breaches copyright please contact [open.access@hw.ac.uk](mailto:open.access@hw.ac.uk) providing details, and we will remove access to the work immediately and investigate your claim.

# Porosity gradients as a means of driving lateral flows at cometary surfaces

Chariton Christou<sup>1,2\*</sup>, S Kokou Dadzie<sup>2</sup>, Raphael Marschall<sup>3,4</sup>, Nicolas Thomas<sup>3</sup>

<sup>1</sup> Tickmill Group, EC2Y 9DT, London, UK

<sup>2</sup> Heriot-Watt University, Institute of Mechanical Process and Energy Engineering, EH14 4AS Edinburgh, Scotland, UK

\*Corresponding Author E-mail: [c.christou@tickmill.com](mailto:c.christou@tickmill.com)

<sup>3</sup>Physikalisches Institut, University of Bern, Switzerland

<sup>4</sup>International Space Science Institute, Hallerstrasse 6, CH-3012 Bern, Switzerland.

## Abstract

The Rosetta spacecraft has provided invaluable and unexpected information about cometary outgassing. The on-board instruments ROSINA, MIRO, and VIRTIS showed non-uniform outgassing of H<sub>2</sub>O over the surface of the nucleus. Rarefied gas flows display remarkable flow phenomena that may help explain diverse physical observations and models have been used in various engineering applications. Typical examples are flows generated by temperature variations such as thermal creeping flows that may have applications in cometary research. In this paper, we demonstrate how porosity variations in a porous medium combined with rarefaction can also lead to high tangential speed flow at a surface. Outgassing through a porous surface layer has been postulated as one possible mechanism at the cometary surfaces. We use the Direct Simulation Monte Carlo (DSMC) method to investigate such flows. The porous media structure consists of micro computed tomography (micro-CT) image of real Earth rock samples with high resolution. We show that lateral gas density gradients can be produced in a straightforward manner leading to substantial non-radial flow.

## 1. Introduction

Comets are widely assumed to be the least processed remaining bodies from the formation of the Solar System. The Rosetta mission was designed to investigate a comet in detail and understanding sublimation processes, surface characteristics, and coma formation of comet 67P/Churyumov-Gerasimenko (hereafter 67P) was one of the main objectives. After a ten-year cruise, the Rosetta spacecraft reached 67P on 6 August 2014. The spacecraft was orbiting and observing the comet for two years and has provided a large amount of data. The nucleus of 67P consists of two large lobes joined at a “neck”[1].

The Rosetta orbiter housed a suite of eleven instruments along with the Philae lander in order to observe the nucleus and the coma. The millimetre-submillimeter wave instrument, Microwave Instrument for the Rosetta Orbiter (MIRO), supported analysis of the water vapour distribution in the comet’s inner coma [2]. The Rosetta Orbiter Spectrometer for Ion and Neutral Analysis (ROSINA) determined composition and density in situ. The Comet Pressure Sensor (COPS) of ROSINA provides the total gas density [3]. Hassig *et al.* observed heterogeneities

in the coma of 67P with large fluctuations using ROSINA data [4] and Marschall *et al.* modelled the COPS pressure variations with time in the November 2014 period using an inhomogeneous outgassing distribution [5]. Similar results were also reported by other studies [6, 7]. The Visual InfraRed Thermal Imaging Spectrometer (VIRTIS) instrument made measurements of the coma and showed that the neck of 67P is the region where H<sub>2</sub>O mostly emanates during northern summer [8]. This was in agreement with the modelling of the COPS data. However, CO<sub>2</sub> emanates from the two lobes [9]. Differences in the outgassing distributions of different species had been seen during flybys of other comets [10].

The inhomogeneities in gas outgassing could arise from variations in the surface layer porous structure. But a non-uniform dust layer composed of particles with varying size could also be realistic scenario [11, 12]. Non-uniform outgassing from the nucleus surface might give rise to the inhomogeneities [12]. Coma inhomogeneities can also be attributed to non-uniformly distributed sources of H<sub>2</sub>O and CO<sub>2</sub> over the surface of the nucleus [13].

The Optical Spectroscopic and Infrared Imaging System (OSIRIS) has yielded unprecedented views of 67P. Images have shown that the surface is morphologically complex. Thomas *et al.* reported the presence of large dune-like ripples on 67P [14]. The presence of such dunes on a comet challenged the existing knowledge and understanding of comet surface processes as it indicated non-radial gas flow. The suggestion that water vapor is emitted from a subsurface ice and travel through a porous medium was considered realistically in a detailed study [15].

The Philae lander contacted the surface of 67P on the northern hemisphere of the nucleus [16]. The Multipurpose Sensors for Surface and Sub-Surface Science (MUPUS) was designed to measure thermomechanical properties of the surface and the near-surface layer and to monitor the subsurface temperature [17]. However, MUPUS was unable to be inserted and thermomechanical properties were not measured. Hammering was increased to the highest energy level but it was not able to be inserted more than 27mm [17]. Spohn *et al.* interpreted this failure to a near-surface layer of a strength that MUPUS was not capable of penetrating [17]. With a data calibration, the uniaxial (or unconfined) compressive strength of the material was calculated to be about 2MPa. This is at odds with measurements of the tensile and compressive strengths elsewhere and has led to the idea that the surface layer may be locally inhomogeneous in its physical properties. Furthermore, Brouet *et al.* reported the presence of an upper layer with different physical properties to the interior and suggested the existence of a porosity gradient with increasing porosity with depth [18].

Comets are now known to be highly porous and could have a porosity around 85% [11] or in the range of 75 – 85% at the “head” [19]. Similar porosity values have also been calculated but with lower boundaries corresponding to a porosity 70 -75 % [20].

The above observations suggest that inhomogeneity in outgassing is a highly probable property of cometary nuclei, that surface morphology is variable over many scales, and that porosity gradients exist with depth and possibly with lateral position. We expect these properties to influence the near-surface gas flow field. Here we investigate one hypothesis, namely, whether porosity gradients in a surface layer can influence the gas flow from a sublimating sub-surface ice layer. We do this using physical analogues.

We have used micro computed tomography (micro-CT) of terrestrial rock samples to demonstrate some possible rarefied gas flow phenomena. The direct simulation Monte Carlo (DSMC) method is adopted. DSMC has been previously used for simulating the activity around comets and it has been able to reproduce ROSINA observations [7]. The porous medium sample consists of porosity between 45-92 %. In the next section, we discuss the porous material analogues. In section 3, we discuss the model. In section 4, we describe the main results and then conclude.

## 2. X-Ray Micro Computed Tomography (micro-CT) Images

A typical method to study gas transport in near-surface comet porous layers is to generate porous media with monodisperse spheres using either random ballistic deposition (RBD) or random sequential packing (RSP) [11]. While this method can provide us with any targeted porosity it does not represent real porous medium (rock) samples. Feldkamp *et al.* first reported micro-CT X-Ray technology in 1980's and it was used to study bones [21]. Micro-CT technology can represent any object with high resolution in three dimensions (3D) and its applications are found in several areas of engineering including medical applications and petroleum industries. Micro-CT allows scientists to analyze the medium in three-dimensions (3D). The method enables a direct measurement of the sample microarchitecture. The explosion of unconventional gas extraction in the oil and gas sectors has boosted micro-CT scan of porous media. The fact that unconventional gas reservoirs consist of low porosity media indicates also the need for high resolution images in this area. Resolutions of micro-CT images are now in the scale of  $\mu\text{m}$  and  $\text{nm}$  [22]. Micro-CT is still though a relatively new technique. The high resolution of this method allows us to calculate the porosity of a sample through the reconstructed image. It should be noted that micro-CT is a non-destructive method, hence samples that have been scanned and are used in simulations can be re-evaluated at any time [23]. Figure 1 shows the binary image pore network of three rock types.

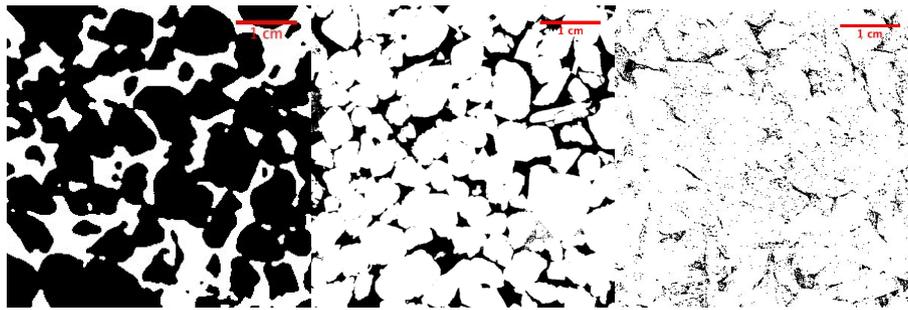


Figure 1: Binary porosity gradient X-Ray Micro-CT image cross section of porous medium samples used with a resolution of  $5\mu\text{m}$ . White colour is the material and the black colour is the void area. From left to right ( $\phi=45\%$ ,  $\phi=80\%$ ,  $\phi=92\%$ )

Micro-CT technology captures the pore space in detail. Images can then be implemented in an unstructured mesh after the three-dimensional image reconstruction. The capability to use an unstructured mesh directly on the images allows us to use different Computational Fluid Dynamics codes (here the DSMC [24]) to model the fluid flow and extract flow properties through the medium. As DSMC is a discrete method it is convenient for simulating the rarefied gas flow in these complex geometries.

## 3. Model

Porous structures are complex geometries and networks consisting of pores with different size. Gas flow in porous medium is usually modelled by Darcy's law which is derived from the standard Navier-Stokes equations [25]. Here our outgassing processes through the comet surface from the subsurface are purely rarefied. The degree of rarefaction is characterized by the

Knudsen number,  $Kn$  (i.e., the ratio of the mean free path,  $\lambda$ , to the characteristic length of porous medium,  $H$ )

$$Kn = \lambda H \quad (1)$$

Fluid flow is often classified into four regimes [26]: continuum flow ( $Kn < 0.001$ ) in which Navier-Stokes equations are applicable; slip flow regime ( $0.001 < Kn < 0.1$ ) where Navier-Stokes with velocity slip and temperature jump can be used to describe the flow; transition regime ( $0.1 < Kn < 10$ ) and free molecular regime ( $Kn > 10$ ). In the transition and free molecular regimes Navier-Stokes are inapplicable since the flow is far from equilibrium and discrete methods such as DSMC are usually adopted [27]. DSMC is a kinetic approach that provides solutions to the Boltzmann equation. Despite its computational demand, DSMC is valid for all collisional regimes encountered in the cometary coma. Depending on the heliocentric distance, the outgassing from the comet surface is found in the transition (close to perihelion) and free molecular (near aphelion) regimes and consequently the gas is in a non-equilibrium state because of the very few intermolecular collisions. In the case where the Knudsen number is greater than unity then the mean free path is higher than the porous media thickness and as a result more collisions with pore walls take place.

Particles are inserted below the porous gradient structure, where we assume the sublimation is taking place. Intermolecular collisions in the current study are treated using the Variable Soft Sphere (VSS) model which has advantages over other models [28]. The outer and lateral boundaries are considered to be a vacuum like boundaries. Wall reflection boundaries for the porous media are set to diffuse. For the gas-surface interaction model, diffuse model considered to be more appropriate in comparison with specular model especially for porous media surfaces [24].

Driven by outgassing inhomogeneity observed on 67P, we use Earth basin rock samples of different porosities to construct a porosity gradient geometry. DSMC is applied on the 3D rock images to investigate the outgassing process.

## 4. A Porosity Gradient Driving Flow

We start our investigation by evaluating terrestrial rock samples individually (Figure 2) in order to illustrate how porosity influences the rarefied flow outgassing pressure. We evaluate six different samples with porosities ranging from 24 – 92 % with dimensions of the rock equal to 5cmx5cmx5cm. We assume gas emission from the interior with flow through the porous samples to the surface. The gas production rate is  $Z = 3 \times 10^{22}$  molecules  $m^{-2} s^{-1}$  in the interior. Table 1 reports the pressures at the surfaces of the porous rock samples. The surface outgassing pressure ratio between rocks with porosity 92 % and 45 % is nearly 50. Rock samples with relatively low porosity (24 % and 45 %) are acting almost as seal rocks; that is, preventing migration of gas molecules to the surface.

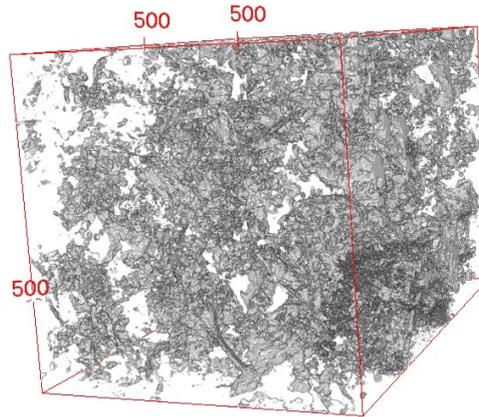


Figure 2: X-Ray Micro-CT reconstructed 3D digital image sample

Table 1: Surface gas pressure as a function of porosity for production rate,  $Z = 3 \times 10^{22}$  molecules  $m^{-2} s^{-1}$

Porosity (%)	Pressure (Pa $\times 10^{-4}$ ) at the surface
92	150
90	110
88	112
80	54.4
45	3.6
24	3.1

From Table 1, we hypothesize that the existence of a lateral porosity gradient along the surface in these rarefaction conditions results in a steep lateral pressure gradient and thus later flow of gas across the surface. We therefore construct a porosity gradient flow configuration combining three rock samples (Figure 3). Rock samples selected correspond to porosities: 45%, 80% and 92 %. The combined three different rock samples represents our non-uniform porous comet structure. Two production/sublimation rates are investigated. The effects of a porosity gradient orthogonal to this, i.e., a porosity increasing or decreasing with depth has previously been investigated [18]. Here, we focus on the lateral porosity gradient effects on the flows in the near surface boundary layer assuming water vapour as the outgassing gas. For the DSMC simulations, binary collisions are treated using the variable soft sphere model (VSS). Collisions of gas molecules with the boundaries are assumed to be fully diffusive [28]. Such a model has been previously reported to have advantages over a specular reflection model. Each simulation case consists of more than five million computational cells. All cells meet the size criterion, in order to be less than the mean free path ( $\lambda$ ). Each computational cell contains approximately thirty particles. The DSMC code implemented in the OpenFOAM framework is used in the current study [29].

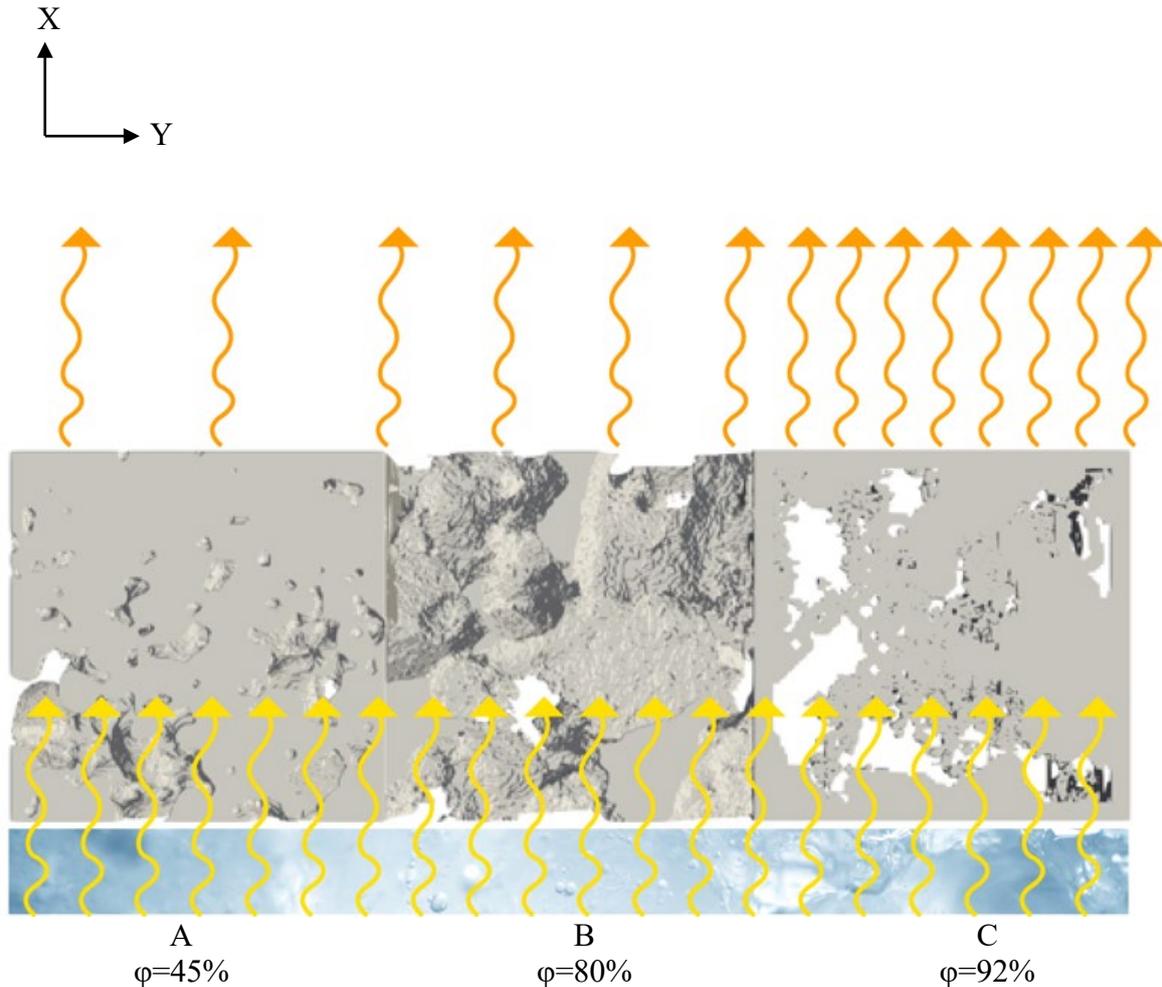


Figure 3: Problem configuration (porosity gradient structure). Rock temperature,  $T_{rock} = 300K$  and gas sublimation temperature  $T_{gas} = 200K$

Figure 3 depicts the configuration under investigation. It is suggested that water ice is present very close (some cm's) below the surface in several places of 67P [30]. Thus, we use a micro-CT digital rock with total length of 15 cm and a height equal to 5 cm. Sublimation is assumed to take place 5 cm below the surface at a uniform rate and temperature  $T_{gas} = 200K$ . Two gas production rates are assumed, namely,  $Z = 3 \times 10^{19} \& 22$  molecules  $m^{-2} s^{-1}$  [31]. The porous rock sample has a constant temperature  $T_{rock} = 300K$ . Above the surface is considered an open-domain. Next, we look into the flow properties at a distance of few centimeters above the surface, on the Y-axis. This configuration which is on the centimetre scale may be viewed as a pore scale modelling rather than at the full 67P scale. An upscaling method may be used to translate results into the macroscopic or kilometre scale [32]. Then gas properties obtained in the few cm's scale above the surface can serve as input in models describing the gas flow in the scale of the km [33].

## 5. Results and Discussion

Figure 4 illustrates the lateral pressure profiles at 1cm above surface from DSMC simulations data for the two heliocentric distances. High surface pressure is observed at the high porosity side and low at the low porosity side. We observe therefore a steep pressure gradient in the direction of increasing porosity generated by the existence of a lateral porosity gradient despite a constant production rate at source. The outgassing pressure shows a rapid increase with a small change in porous media porosity. In more detail, Figure 4a and 4b shows that outgassing pressure can be up to seven times higher by increasing the porosity 47 %.

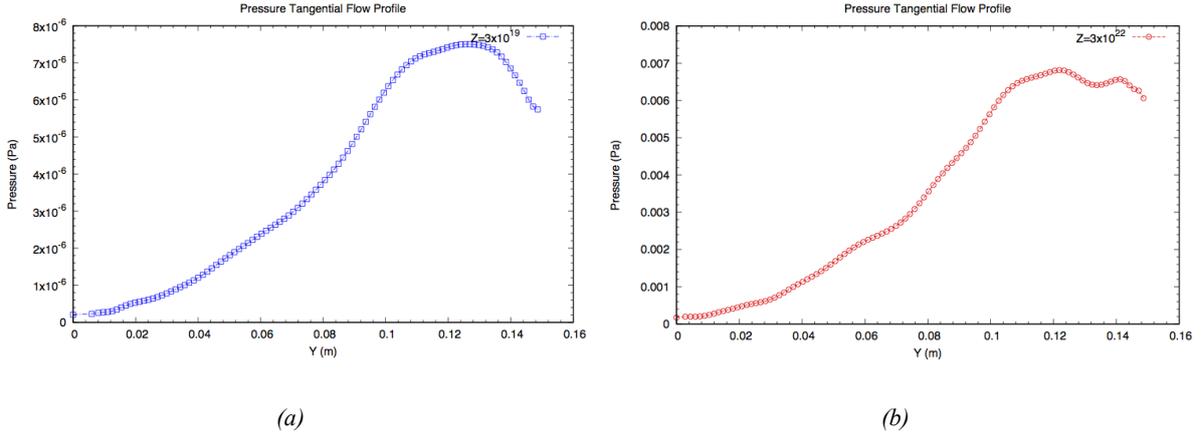


Figure 4: Outgassing pressure profiles for two different production rates at 1cm above surface. a)  $Z=3 \times 10^{19}$  and b)  $Z=3 \times 10^{22}$

In Figure 5 we look into the lateral velocity profile in the negative pressure gradient direction corresponding to porosity gradient set-up in the previous section. We observe a lateral velocity of over 550 m/s which expands into the low pressure region. The fraction of the lateral component to the speed of the gas is 0.85. This indicates that the flow is indeed mainly parallel to the surface. These lateral flow velocities are difficult to compare with Rosetta data but these values are similar to those seen for radial expansion. Gulkis *et al.* reported from MIRO observations, a radial expansion velocity of 680 m/s which is comparable to the range of lateral velocity obtained in the high porosity region from our simulations [34]. Lower values than the one reported by Gulkis *et al.* were reported by Bieler *et al.* [6]. Biver *et al.* investigated the expansion velocity for a heliocentric distance of 3.4 au and reported values of 470 – 590 m/s [13]. Marschall *et al.* hypothesized that the red emission wing in the HO line shape of some MIRO spectra could indicate an almost stagnant layer close to the nucleus surface. Due to their observed viewing geometry such a stagnant layer would not imply zero gas velocity but rather a lateral gas flow [35].

On the modelling side, Thomas *et al.* by assuming a gas number density of  $5 \times 10^{16} \text{ m}^{-3}$  calculated a shear velocity at a fluid threshold for  $d = 1$  cm diameter particles that gave 528 m/s and concluded that particles of that size and larger might be lifted by such a flow [36]. The gas pressure used in that study was comparable to the one used in our DSMC simulations. It is worth noting that shear velocity at velocity threshold is proportional to gas pressure. Recently Jia *et al.* compared bedforms observed on 67P to giant ripples [19].

In the case where the production rate is equal to  $Z = 3 \times 10^{19} \text{ molecules m}^{-2} \text{ s}^{-1}$  the gas is more rarefied and the Knudsen number is three orders of magnitudes higher. Hence, the outgassing

flow is far from equilibrium and there are more collisions between gas molecules and boundaries than intermolecular collisions. Porosity gradient flows occurs in both reported production rates (Figure 5) however the velocity is lower for the higher Knudsen number (low production rate).

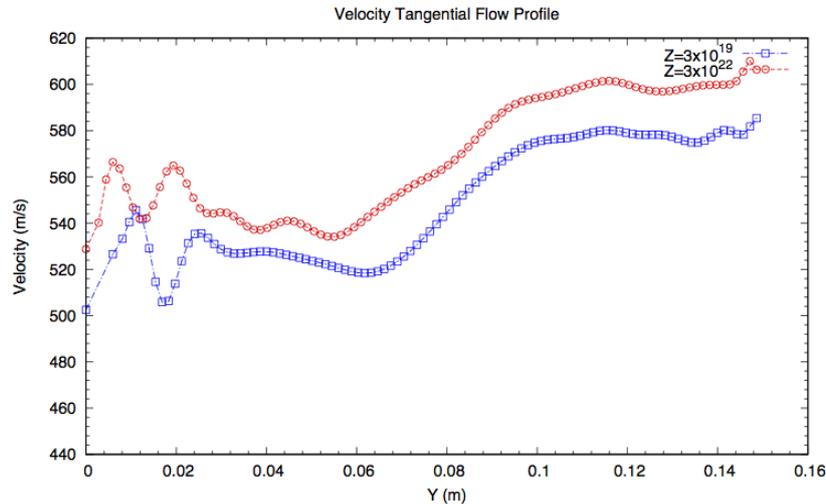


Figure 5: Outgassing lateral velocity profiles for two different production rates

Comparing Figure 4 and 5 we observe a positive correlation between pressure and velocity profiles. Such observations confirm some previous reports [37]. Figure 6 presents the temperature profile corresponding to the porosity gradient driving flow configuration. The non-uniformity in the comet near surface porous media layer results in lateral temperature gradients. The outgassing gas temperature at the surface decreases as the porosity increases. That is, the gas temperature at the surface is inversely proportional to porosity. On average, 170 K is the minimum gas surface temperature obtained with the higher porosity medium. The highest temperature is found for a medium with porosity of 45 % which is on average 215 K [38]. In other words, the temperature jump at the nucleus surface increases with porosity. Taking into account this temperature gradient it may be concluded that the higher speed flow depicted in Figure 5 is a combination of both temperature and pressure gradient effects.

These lateral flows may be capable of generating grain transports. Lara *et al.* reported that dust particles in the scale of mm could be ejected [39]. More recently, Lia *et al.* calculated that during peak production, near perihelion passage, dust grains can be as large as up to 10mm when ejected from the nucleus surface [40]. Based on those observations and the lateral velocity in Figure 5 we can conclude that porosity gradient can generate the transport of grains that are ejected from the surface.

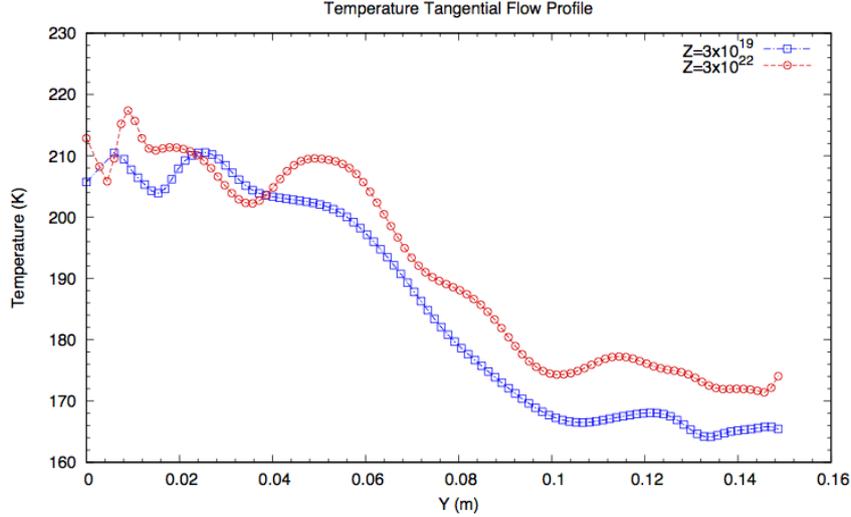


Figure 6: Outgassing temperature above surface for porosity gradient; lower porosity on the left side and higher on the right

Lee *et al.* reported that outgassing activity can be affected by many local environmental factors such as: the present amount of water ice, porosity layer and surface temperature [41]. The current work directly depicts the porosity layer and surface temperature factors. Starting from porosity of the medium in the near-surface non-equilibrium outgassing activity it has been shown that it can be affected by porosity variations and interestingly it can be driven by it. Such flows can also be found in both the active and inactive regions. However, inactive regions are with higher Knudsen number flows and hence with more rarefied conditions [3]. Figure 4 indicates that the present porosity gradient flows may be observed in both regions.

## 6. Concluding Remark

We have presented numerical simulations of gas dynamics within the first few centimetres of comet-like surface layers. We demonstrate how a lateral porosity gradient can generate steep lateral pressure and temperature gradients leading to potentially high speed lateral flows. This follows observations of sediment transports at the surface of comet 67P. These results add a new perspective to the observed inhomogeneities in coma formation. Previously lateral surface flows leading to dune-like structures observed on 67P were explained by temperature differences between highly illuminated and shadowed regions on the nucleus. Our present results indicate that porosity variations by themselves may be responsible for lateral flows and we suggest that these may be capable of generating substantial grain transports. We also observed that a difference in porosity creates a large difference in the difference between gas and comet material temperature at the surface (i.e., in the surface temperature jump).

We found that even small variations in porosity (12%) result in significant changes in outgassing properties. For the reported observations at a distance of a few centimetres above the surface medium, an increase in porosity, results in an increase in surface water vapour pressure. Summarizing, we propose that a structure with varying porosity that might be present in different regions of the comet's surface could also contribute to the observed coma inhomogeneities. We conclude that the coma might have an additional non-time-dependent component which is related to lateral porosity gradients.

## References:

1. Sierks, H., et al., *On the nucleus structure and activity of comet 67P/Churyumov-Gerasimenko*. Science, 2015. **347**(6220): p. aaa1044.
2. Gulkis, S., et al., *MIRO: Microwave instrument for Rosetta orbiter*. Space Science Reviews, 2007. **128**(1-4): p. 561-597.
3. Hässig, M., et al., *Time variability and heterogeneity in the coma of 67P/Churyumov-Gerasimenko*. Science, 2015. **347**(6220): p. aaa0276.
4. Kokorev, V.I., et al., *The Impact of Thermogas Technologies on the Bazhenov Formation Studies Results*. Society of Petroleum Engineers.
5. Marschall, R., et al., *Cliffs versus plains: Can ROSINA/COPS and OSIRIS data of comet 67P/Churyumov-Gerasimenko in autumn 2014 constrain inhomogeneous outgassing?* 2017.
6. Bieler, A., et al., *Comparison of 3D kinetic and hydrodynamic models to ROSINA-COPS measurements of the neutral coma of 67P/Churyumov-Gerasimenko*. Astronomy & Astrophysics, 2015. **583**: p. A7.
7. Fougere, N., et al., *Three-dimensional direct simulation Monte-Carlo modeling of the coma of comet 67P/Churyumov-Gerasimenko observed by the VIRTIS and ROSINA instruments on board Rosetta*. Astronomy & Astrophysics, 2016. **588**: p. A134.
8. Bockelée-Morvan, D., et al., *First observations of H<sub>2</sub>O and CO<sub>2</sub> vapor in comet 67P/Churyumov-Gerasimenko made by VIRTIS onboard Rosetta*. Astronomy & Astrophysics, 2015. **583**: p. A6.
9. Migliorini, A., et al., *Water and carbon dioxide distribution in the 67P/Churyumov-Gerasimenko coma from VIRTIS-M infrared observations*. Astronomy & Astrophysics, 2016. **589**: p. A45.
10. Feaga, L., et al., *Asymmetries in the distribution of H<sub>2</sub>O and CO<sub>2</sub> in the inner coma of Comet 9P/Tempel 1 as observed by Deep Impact*. Icarus, 2007. **191**(2): p. 134-145.
11. Skorov, Y.V., et al., *Activity of comets: Gas transport in the near-surface porous layers of a cometary nucleus*. Icarus, 2011. **212**(2): p. 867-876.
12. Shi, X., et al., *Sunset jets observed on comet 67P/Churyumov-Gerasimenko sustained by subsurface thermal lag*. Astronomy & Astrophysics, 2016. **586**: p. A7.
13. Biver, N., et al., *Distribution of water around the nucleus of comet 67P/Churyumov-Gerasimenko at 3.4 AU from the Sun as seen by the MIRO instrument on Rosetta*. Astronomy & Astrophysics, 2015. **583**: p. A3.
14. Thomas, N., et al., *The morphological diversity of comet 67P/Churyumov-Gerasimenko*. Science, 2015. **347**(6220): p. aaa0440.
15. Jia, P., B. Andreotti, and P. Claudin, *Giant ripples on comet 67P/Churyumov-Gerasimenko sculpted by sunset thermal wind*. Proceedings of the National Academy of Sciences, 2017: p. 201612176.
16. Biele, J., et al., *The landing (s) of Philae and inferences about comet surface mechanical properties*. Science, 2015. **349**(6247): p. aaa9816.
17. Spohn, T., et al., *Thermal and mechanical properties of the near-surface layers of comet 67P/Churyumov-Gerasimenko*. Science, 2015. **349**(6247): p. aab0464.
18. Brouet, Y., et al., *A porosity gradient in 67P/CG nucleus suggested from CONSERT and SESAME-PP results: an interpretation based on new laboratory permittivity measurements of porous icy analogues*. Monthly Notices of the Royal Astronomical Society, 2016. **462**(Suppl\_1): p. S89-S98.

19. Pätzold, M., et al., *A homogeneous nucleus for comet 67P/Churyumov–Gerasimenko from its gravity field*. *Nature*, 2016. **530**(7588): p. 63-65.
20. Jorda, L., et al., *The global shape, density and rotation of Comet 67P/Churyumov-Gerasimenko from preperihelion Rosetta/OSIRIS observations*. *Icarus*, 2016. **277**: p. 257-278.
21. Feldkamp, L.A., et al., *The direct examination of three-dimensional bone architecture in vitro by computed tomography*. *Journal of bone and mineral research*, 1989. **4**(1): p. 3-11.
22. Dong, H., *Micro-CT imaging and pore network extraction*. 2007.
23. Ho, S.T. and D.W. Huttmacher, *A comparison of micro CT with other techniques used in the characterization of scaffolds*. *Biomaterials*, 2006. **27**(8): p. 1362-1376.
24. Christou, C. and S.K. Dadzie, *Direct-Simulation Monte Carlo Investigation of a Berea Porous Structure*. *SPE Journal*, 2015.
25. Darcy, H., *Les fontaines publiques de la ville de Dijon: exposition et application*. 1856: Victor Dalmont.
26. Karniadakis, G., A. Beşkök, and N.R. Aluru, *Microflows and nanoflows : fundamentals and simulation*. *Interdisciplinary applied mathematics*. 2005, New York, NY: Springer. xxi, 817 p.
27. Chapman, S. and T.G. Cowling, *The mathematical theory of non-uniform gases: an account of the kinetic theory of viscosity, thermal conduction and diffusion in gases*. 1970: Cambridge university press.
28. Shen, C., *Rarefied gas dynamics: fundamentals, simulations and micro flows*. 2006: Springer.
29. Scanlon, T., et al., *Open source DSMC chemistry modelling for hypersonic flows*. *AIAA Journal*, 2014.
30. Auger, A.-T., et al., *Meter-scale thermal contraction crack polygons on the nucleus of comet 67P/Churyumov-Gerasimenko*. *Icarus*, 2018. **301**: p. 173-188.
31. Hansen, K.C., et al., *Evolution of water production of 67P/Churyumov–Gerasimenko: an empirical model and a multi-instrument study*. *Monthly Notices of the Royal Astronomical Society*, 2016. **462**(Suppl\_1): p. S491-S506.
32. Valvatne, P.H. and M.J.J.W.r.r. Blunt, *Predictive pore-scale modeling of two-phase flow in mixed wet media*. 2004. **40**(7).
33. Liao, Y., et al., *3D Direct Simulation Monte Carlo modelling of the inner gas coma of comet 67P/Churyumov–Gerasimenko: a parameter study*. *Earth, Moon, and Planets*, 2016. **117**(1): p. 41-64.
34. Gulkis, S., et al., *Subsurface properties and early activity of comet 67P/Churyumov-Gerasimenko*. *Science*, 2015. **347**(6220): p. aaa0709.
35. Marschall, R., et al., *A comparison of multiple Rosetta data sets and 3D model calculations of 67P/Churyumov-Gerasimenko coma around equinox (May 2015)*. *Icarus*, 2019. **328**: p. 104-126.
36. Thomas, N., et al., *Redistribution of particles across the nucleus of comet 67P/Churyumov-Gerasimenko*. *Astronomy & Astrophysics*, 2015. **583**: p. A17.
37. Davidsson, B.J., et al., *Gas kinetics and dust dynamics in low-density comet comae*. *Icarus*, 2010. **210**(1): p. 455-471.
38. Christou, C., et al., *Gas flow in near surface comet like porous structures: Application to 67P/Churyumov-Gerasimenko*. *Planetary and Space Science*, 2018.
39. Lara, L., et al., *Large-scale dust jets in the coma of 67P/Churyumov-Gerasimenko as seen by the OSIRIS instrument onboard Rosetta*. 2015. **583**: p. A9.
40. Lai, I.-L., et al., *Gas outflow and dust transport of comet 67P/Churyumov–Gerasimenko*. 2017. **462**(Suppl\_1): p. S533-S546.

41. Lee, S., et al., *Spatial and diurnal variation of water outgassing on comet 67P/Churyumov-Gerasimenko observed from Rosetta/MIRO in August 2014*. *Astronomy & Astrophysics*, 2015. **583**: p. A5.