



Heriot-Watt University
Research Gateway

Extended Riemannian Geometry III

Citation for published version:

Deser, A & Saemann, C 2019, 'Extended Riemannian Geometry III: Global Double Field Theory with Nilmanifolds', *Journal of High Energy Physics*, vol. 2019, no. 5, 209.
[https://doi.org/10.1007/JHEP05\(2019\)209](https://doi.org/10.1007/JHEP05(2019)209)

Digital Object Identifier (DOI):

[10.1007/JHEP05\(2019\)209](https://doi.org/10.1007/JHEP05(2019)209)

Link:

[Link to publication record in Heriot-Watt Research Portal](#)

Document Version:

Publisher's PDF, also known as Version of record

Published In:

Journal of High Energy Physics

General rights

Copyright for the publications made accessible via Heriot-Watt Research Portal is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

Take down policy

Heriot-Watt University has made every reasonable effort to ensure that the content in Heriot-Watt Research Portal complies with UK legislation. If you believe that the public display of this file breaches copyright please contact open.access@hw.ac.uk providing details, and we will remove access to the work immediately and investigate your claim.

Extended Riemannian geometry III: global Double Field Theory with nilmanifolds

Andreas Deser^a and Christian Sämann^{b,c,d}

^a*Faculty of Mathematics and Physics, Charles University,
Sokolovská 83, 186 75 Praha 8, Czech Republic*

^b*Department of Mathematics, Heriot-Watt University,
Colin Maclaurin Building, Riccarton, Edinburgh EH14 4AS, U.K.*

^c*Maxwell Institute for Mathematical Sciences,
Edinburgh, U.K.*

^d*Higgs Centre for Theoretical Physics,
Edinburgh, U.K.*

E-mail: andreas3deser@gmail.com, c.saemann@hw.ac.uk

ABSTRACT: We describe the global geometry, symmetries and tensors for Double Field Theory over pairs of nilmanifolds with fluxes or gerbes. This is achieved by a rather straightforward application of a formalism we developed previously. This formalism constructs the analogue of a Courant algebroid over the correspondence space of a T-duality, using the language of graded manifolds, derived brackets and we use the description of nilmanifolds in terms of periodicity conditions rather than local patches. The strong section condition arises purely algebraically, and we show that for a particularly symmetric solution of this condition, we recover the Courant algebroids of both nilmanifolds with fluxes. We also discuss the finite, global symmetries of general local Double Field Theory and explain how this specializes to the case of T-duality between nilmanifolds.

KEYWORDS: Differential and Algebraic Geometry, Flux compactifications, String Duality

ARXIV EPRINT: [1812.00026](https://arxiv.org/abs/1812.00026)

Contents

1	Introduction and results	1
2	Geometry of gerbes on three-dimensional nilmanifolds	6
2.1	Three-dimensional nilmanifolds	6
2.2	Gerbes on three-dimensional nilmanifolds	7
2.3	Exact Courant algebroids and generalized metrics	8
2.4	Exact Courant algebroids as symplectic Lie 2-algebroids	9
2.5	The Courant algebroid of a gerbe	10
2.6	The Courant algebroid of $\mathcal{G}_{j,k}$	12
3	Finite global symmetries of Generalized Geometry	14
3.1	Lie 2-algebra associated to $\mathcal{C}_{j,k}$	15
3.2	Action of the associated Lie algebra and its action Lie algebra	16
3.3	Finite symmetries of a gerbe	18
4	Global Double Field Theory	18
4.1	Topological T-duality	18
4.2	Symplectic pre-NQ-manifold on the correspondence space	19
4.3	Symmetries, L_∞ -structures and the section condition	21
4.4	L_∞ -algebra structures: solutions to the section condition	23
4.5	T-duality between $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$	25
4.6	Differential geometry on \mathcal{E}_K	25
5	Finite symmetries of Double Field Theory	26
5.1	Generalized Lie derivative and infinitesimal symmetries	27
5.2	Finite symmetries	29
5.3	Specialization to \mathcal{E}_K	31
6	Discussion	31

1 Introduction and results

The goal of Double Field Theory is to provide a Lagrangian which is manifestly symmetric under the hidden $O(d, d; \mathbb{Z})$ T-duality symmetry of string theory and governs the dynamics of the massless sector of superstring theory, see [1–3] and references therein. Currently, the available Lagrangian and its gauge structure are given only on a local patch of target space¹ and it is suggested that some form of gluing has to be performed to recover the

¹With minor exceptions; see e.g. [4, 5], [6–8] or also [9] for alternative approaches to ours.

global picture. Since T-duality is fundamentally linked to compact target space directions, a clean understanding of this gluing procedure is important.

In this paper, we study this problem in the case of the prominent example of 3-dimensional nilmanifolds. These are circle bundles over the torus T^2 and we equip them with a three form flux, which mathematically corresponds to the curvature of a gerbe. The advantage of this class of examples is that the non-trivial patchings can be understood in terms of twisted periodicity conditions on the local data. We will use the general formalism we presented previously in [10] and extend the local picture developed there to a global one by demanding compatibility with the periodicity conditions.

In physics terms, we are studying the NS-NS sector of supergravity without dilaton,² consisting of a vielbein (or metric) on a topological manifold M carrying an abelian gerbe with connective structure (i.e. Kalb-Ramond field) B and curvature (i.e. flux) $H = dB$. The symmetries of this theory consist of diffeomorphisms and gauge transformations of the B -field together with the gauge-for-gauge transformations of the latter. At the infinitesimal level, these are elegantly described by Hitchin's *Generalized Geometry*, cf. e.g. [12]. The vector field and the 1-form parameterizing these symmetries are sections of the *generalized tangent bundle*, which has some non-trivial gluing properties if the gerbe is topologically non-trivial.

The perspective of the generalized tangent bundle can be further improved by considering the corresponding Courant algebroid in the language of symplectic differential graded manifolds or Q -manifolds as pioneered in [13]. The infinitesimal symmetry structure of NS-NS gravity is then simply obtained in terms of derived brackets, cf. [14, 15]. From a more mathematical perspective, abelian gerbes are categorified principal bundles or principal 2-bundles for short. For each topological class of abelian gerbes, there is a corresponding class of Courant algebroids (and Courant algebroids are best regarded as symplectic Lie 2-algebroids). The Lie 2-algebra³ of symmetries of a gerbe is then the associated Lie 2-algebra of the corresponding Courant algebroid, obtained from the derived bracket construction of [16, 17].

The aim is now to lift this rather well-understood global picture from Generalized Geometry to Double Field Theory. The latter associates additional winding degrees of freedom to the compact circle fibers of the nilmanifold. Locally, T-duality is understood as an exchange of the coordinates Fourier-dual to the momentum and winding degrees of freedom but globally, we have topological features complicating the picture.

One approach to a global picture of Double Field Theory is to use the finite symmetry actions to glue together local descriptions [18]. Recall that in the case of Generalized Geometry, the infinitesimal symmetries are governed by the Dorfman bracket. While the symmetries themselves form a Lie 2-algebra which is awkward to integrate, the *actions* of the infinitesimal symmetries by the Dorfman bracket form a Lie algebra, which integrates to the semidirect product of diffeomorphisms and closed 2-forms on the base manifold. In the case of Double Field Theory, the infinitesimal symmetries are governed by an analogue

²A generalization to heterotic NS-NS gravity along the lines discussed in [11] should be straightforward.

³By this term, we shall always mean a 2-term L_∞ -algebra, which is essentially the same.

of the Dorfman bracket, often called the generalized Lie derivative or D-bracket. Again, the symmetries form a Lie 2-algebra, and the actions by the D-bracket form an ordinary Lie algebra which can be integrated to finite transformations [19], see also [20] for further discussions and [21] for a slightly different perspective. These finite transformations were then used as gluing transformations for local descriptions of Double Field Theory [18].

There is, however, a problem with this construction, which is already seen in the context of Generalized Geometry. As stated above, each topological class of gerbes comes with its own class of Courant algebroids. Switching from a topologically trivial H -flux (i.e. $H = dB$ globally) to a non-trivial one therefore requires either to modify the Dorfman bracket or to modify the gluing prescriptions for the corresponding generalized tangent bundle. Both changes are related by a coordinate change, as we shall see. Since the generalized Lie derivative in Double Field Theory is almost exclusively discussed without taking this modification into account,⁴ it is not surprising that the above mentioned integrated infinitesimal symmetries, as well as the gluing transformations, are only suitable for topologically trivial gerbes. The same conclusion was reached by explicit computations in [22].

In our previous paper [15], we presented a formalism in which the necessary modifications of the D-bracket due to topologically non-trivial gerbes become natural and can be discussed in detail, see also [23]. The key observation was that the Q -manifold picture of the Courant algebroid over space-time can be extended to an analogous picture over the target space of Double Field Theory. We showed that this requires to drop the condition $Q^2 = 0$, leading us to *symplectic pre-NQ-manifolds*. On these, we recover the generalized Lie derivative of Double Field Theory as a derived bracket, and the strong section condition (or rather a slight weakening thereof) is recovered algebraically by demanding that the derived brackets form a Lie 2-algebra of symmetries, as on the Courant algebroid.

The modification of the Dorfman bracket of Generalized Geometry for topologically non-trivial gerbes with curvature H is obtained by adding H to the Hamiltonian Q of the differential given by the homological vector field $Q = \{Q, -\}$. This modification has a clear analogue for the symplectic pre-NQ-manifolds underlying the structure of Double Field Theory.

The above was already claimed in [15]. In this paper, we show that this sketch can indeed be completed to a full coherent picture. We study the simple example of T-duality between 3-dimensional nilmanifolds carrying topologically non-trivial gerbes. This example is still simple enough to allow for a transparent and detailed discussion in coordinates. In particular, we can use the fact that nilmanifolds are the orbit spaces of particular \mathbb{Z}_3 -actions on \mathbb{R}^3 . Therefore, we do not have to use cumbersome local patches, but we can work with flat space and a number of periodicity conditions to describe functions and global sections of the various bundles. Note also that for more general examples, one needs to leave the realm of smooth manifolds and generalize to noncommutative and nonassociative spaces, which we plan to do in future work.

Recall that three dimensional nilmanifolds N_j are principal circle bundles over the 2-torus T^2 , and thus they are characterized by an integer $j \in \mathbb{Z}$, the Chern class of the bundle.

⁴We are not aware of any discussion of this point before [15].

Similarly, gerbes over N_j are also given by an integer, their Dixmier-Douady class $[H] = k$, which is an element in $H^3(N_j, \mathbb{Z}) \cong \mathbb{Z}$. Under T-duality along the fiber direction of the principal bundle, one obtains the nilmanifold N_k endowed with a gerbe of Dixmier-Douady class j :

$$\mathcal{G}_{j,k} = (N_j, [H] = k) \xleftarrow{\text{T-duality}} \mathcal{G}_{k,j} = (N_k, [H] = j), \tag{1.1}$$

which follows from the Buscher rules and which was explained in detail in [24].

A classical way to understand T-duality is to use correspondence spaces, which are fibered products of original and T-dual manifold over a common base manifold. *Locally*, we associate the original- and dual fiber coordinates to momentum and winding coordinates. In the case at hand, the correspondence space K between N_j and N_k is the fiber product $K = N_j \times_{T^2} N_k$. The gerbes $\mathcal{G}_{j,k}$ and $\mathcal{G}_{k,j}$ come with associated Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$ (by which we always mean the corresponding Q -manifolds as given in [13]), which encode the relevant Lie 2-algebra of symmetries.⁵ Both gerbes and both associated Courant algebroids can be pulled back to the correspondence space K . Tensoring the gerbes, we obtain the gerbe \mathcal{G}_K with associated Courant algebroid \mathcal{C}_K . The space relevant for Double Field Theory is now a graded submanifold \mathcal{E}_K of \mathcal{C}_K . Altogether, we have the following diagram:

$$\tag{1.2}$$

The symplectic pre-NQ-manifold \mathcal{E}_K plays the same role in Double Field Theory as the Courant algebroid for Generalized Geometry. Its sections parameterize the symmetries whose structures and actions are governed by derived brackets. Our first important result is therefore the detailed description of \mathcal{E}_K in terms of coordinates. Explicitly, we lift the periodicity conditions of the coordinates on the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$ to periodicity conditions of the various coordinates on \mathcal{E}_K and we show that the ingredients to construct the C-bracket via a derived bracket are invariant under this periodicity and therefore they are indeed global.

Having found \mathcal{E}_K , we can follow our formalism of [15] further and derive the conditions for the antisymmetrized derived brackets to form a Lie 2-algebra of symmetries. This Lie 2-algebra of symmetries will always appear in the form of a 2-term L_∞ -algebra.⁶ The resulting

⁵A picture of T-duality as an isomorphism of invariant Courant algebroids was given in [25].

⁶Note that this 2-term L_∞ -algebra gives rise to, but is *not* identical to the L_∞ -algebra of Double Field Theory as partially given in [26] which arises from the symmetry Lie 2-algebra via the classical BV formalism, as explained in [27].

conditions are the appropriate global form of (a slight weakening of) the strong section condition for Double Field Theory making the T-duality between nilmanifolds carrying gerbes manifest.

Solutions to this section conditions are found in a very natural way. Given a Poisson structure with compatible vector field Q and corresponding Hamiltonian \mathcal{Q} , we have $Q^2 F = \frac{1}{2} \{\{\mathcal{Q}, \mathcal{Q}\}, F\}$ for functions F on \mathcal{E}_K (actually on any pre-NQ-manifold of even degree). The expression $\{\mathcal{Q}, \mathcal{Q}\}$ now factorizes as

$$\{\mathcal{Q}, \mathcal{Q}\} = p_3 p_4, \tag{1.3}$$

where p_3 and p_4 are the momentum coordinates in the T-duality directions. They form Hamiltonians of the vector fields $\frac{\partial}{\partial x^3} = \{p_3, -\}$ and $\frac{\partial}{\partial x^4} = \{p_4, -\}$ along the lifts of the projections pr_1 and pr_2 to the Courant algebroids,

$$\hat{\text{pr}}_1 : \mathcal{E}_K \rightarrow \mathcal{C}_{j,k} \quad \text{and} \quad \hat{\text{pr}}_2 : \mathcal{E}_K \rightarrow \mathcal{C}_{k,j}. \tag{1.4}$$

Complementing these by global vector fields V_3 and V_4 , we arrive at two integrable distributions, $\langle \frac{\partial}{\partial x^4}, V_4 \rangle$ and $\langle \frac{\partial}{\partial x^3}, V_3 \rangle$. The first defines $\hat{\text{pr}}_1$ and the second one defines $\hat{\text{pr}}_2$. Both projections admit particularly simple sections, which are embeddings $e_1 : \mathcal{C}_{j,k} \hookrightarrow \mathcal{E}_K$ and $e_2 : \mathcal{C}_{k,j} \hookrightarrow \mathcal{E}_K$ given by morphisms of pre-NQ-manifolds. That is, they are smooth maps of graded manifolds and the Hamiltonian of Q as well as the symplectic form are pullbacks of the data on \mathcal{E}_K along e_1 and e_2 .

We note that some aspects of the specialization of our example in which $k \in \mathbb{Z}$ but $j = 0$ were discussed earlier in [28] and the overlapping results agree.

Altogether, we conclude that there is no geometric obstacle to a global description of Double Field Theory, at least for those cases in which the manifest T-duality connects smooth manifolds carrying gerbes.

This paper is structured as follows. In section 2, we review nilmanifolds, gerbes over these and the corresponding Courant algebroids. The symmetry structures (both for infinitesimal and finite symmetries) of these Courant algebroids are then explained in detail in terms of derived brackets in section 3. In section 4, we present the Double Field Theory lift of this picture. We start with a very brief review of topological T-duality for nilmanifolds, as found in [24] before we construct the symplectic pre-NQ-manifold \mathcal{E}_K which encodes the geometry and symmetries of Double Field Theory. We discuss the strong section condition as well as the most relevant solutions, restricting \mathcal{E}_K to the Courant algebroids of the two nilmanifolds with gerbes. We also make some detailed comments on extended differential geometry. In [15], we had introduced the concepts of extended tensors, extended covariant derivatives, extended torsion and Riemann tensors. All this is readily specialized to \mathcal{E}_K , and the extended tensors of \mathcal{E}_K reduce nicely to those of the two Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$ for the corresponding solutions of the strong section condition. Section 5 contains a discussion of the infinitesimal and finite symmetries of general Double Field Theory before we specialize to the case discussed in section 4. We close in section 6 with a brief discussion of generalizations of our constructions.

2 Geometry of gerbes on three-dimensional nilmanifolds

2.1 Three-dimensional nilmanifolds

A *nilmanifold* M is a compact quotient manifold $M = N/\Gamma$, where N is a nilpotent Lie group and Γ a discrete subgroup in N . Recall that a *nilpotent Lie group* is connected and has Lie algebra \mathfrak{g} such that the sequence

$$\mathfrak{g}_1 = [\mathfrak{g}, \mathfrak{g}], \quad \mathfrak{g}_2 = [\mathfrak{g}_1, \mathfrak{g}], \quad \dots \tag{2.1}$$

becomes trivial at some term \mathfrak{g}_k .

We shall be interested in examples of 3-dimensional nilmanifolds.⁷ Recall that up to isomorphisms, there is a unique, connected, simply connected nilpotent Lie group, which is the Heisenberg group $H(3)$ of upper-triangular 3×3 -matrices. The cocompact discrete subgroups Γ_j are classified by an integer $j \in \mathbb{N}$. We shall denote the nilmanifold $H(3)/\Gamma_j$ by N_j .

The nilmanifold N_j is a circle bundle over T^2 , and we shall describe it in terms of coordinates x^i on \mathbb{R}^3 with the following periodicity conditions imposed:

$$(x^1, x^2, x^3) \sim (x^1, x^2 + 1, x^3) \sim (x^1, x^2, x^3 + 1) \sim (x^1 + 1, x^2, x^3 - jx^2). \tag{2.2}$$

Here, x^1 and x^2 are coordinates on T^2 , while x^3 is the coordinate along the fiber. Clearly, we can include the case $j = 0$, for which the circle bundle becomes trivial: $N_0 = T^2 \times S^1 = T^3$.

We have a basis of globally defined 1-forms⁸

$$\bar{\xi}^1 = dx^1, \quad \bar{\xi}^2 = dx^2, \quad \bar{\xi}^3 = dx^3 + jx^1 dx^2 \tag{2.3a}$$

together with dual vector fields

$$\bar{\zeta}_1 = \frac{\partial}{\partial x^1}, \quad \bar{\zeta}_2 = \frac{\partial}{\partial x^2} - jx^1 \frac{\partial}{\partial x^3}, \quad \bar{\zeta}_3 = \frac{\partial}{\partial x^3} \tag{2.3b}$$

satisfying $\langle \bar{\xi}^i, \bar{\zeta}_j \rangle = \delta_j^i$.

The circle bundle $\pi_j : N_j \rightarrow T^2$ can be described using a single transition function h specifying the transition $x^1 \rightarrow x^1 + 1$. In its additive form, it reads as $h = -jx^2$, and we can extend this Čech cocycle to a Deligne cocycle (h, A) , encoding a principal circle bundle with connection A . One compatible choice is the connection $A = jx^1 dx^2$, which satisfies the cocycle condition $A(x^1 + 1) - A(x^1) = -dh$. The resulting curvature reads as $F = jdx^1 \wedge dx^2$ and thus j is the first Chern number $j = \int_{T^2} F$ of the principal circle bundle N_j .

Finally, we note that the nilmanifold N_j comes with the natural metric

$$g = (dx^1)^2 + (dx^2)^2 + (dx^3 + jx^1 dx^2)^2. \tag{2.4}$$

⁷These examples play a prominent role in Thurston's geometrization conjecture: every closed 3-manifold can be decomposed into pieces with one of eight geometric structures, one being that of a nilmanifold.

⁸The bar is merely distinguishing these tensors from a slightly different set (ξ^i, ζ_i) used mostly in the later discussion.

2.2 Gerbes on three-dimensional nilmanifolds

At the heart of Generalized Geometry and T-duality is the fact that the metric g is complemented by a 2-form B which is part of the connective structure of a gerbe. For our purposes, the rather simplistic definition of gerbes in the form of *Hitchin-Chatterjee gerbes* [29] will be sufficient.

Consider a manifold M together with a cover $\sqcup_a U_a \rightarrow M$. Let $U_{ab} := U_a \cap U_b$, $U_{abc} := U_a \cap U_b \cap U_c$, etc. Then a *gerbe*, or $U(1)$ -gerbe, or *Hitchin-Chatterjee gerbe*, is a Čech 2-cocycle h , that is a collection of functions

$$h_{(abc)} : U_a \cap U_b \cap U_c \rightarrow \mathbb{R} \quad \text{with} \quad h_{(abc)} - h_{(abd)} + h_{(acd)} - h_{(bcd)} = 0 \quad \text{on} \quad U_{abcd}. \quad (2.5)$$

Just as in the case of principal circle bundles, we can extend the Čech cocycle encoding the gerbe to a Deligne cocycle encoding a gerbe with *connective structure*, the analogue of a connection for a gerbe. Such a Deligne cocycle consists of the Čech cocycle $h_{(abc)}$ on triple overlaps U_{abc} , 1-forms $A_{(ab)}$ on the double overlaps U_{ab} and 2-forms $B_{(a)}$ on the patches U_a . The cocycle conditions (2.5) are then completed by

$$A_{(ab)} - A_{(ac)} + A_{(bc)} = -dh_{(abc)} \quad \text{on} \quad U_{abc} \quad \text{and} \quad B_{(a)} - B_{(b)} = dA_{(ab)} \quad \text{on} \quad U_{ab}. \quad (2.6)$$

The second cocycle condition implies that the curvatures on each U_a ,

$$H|_{U_a} = dB_{(a)}, \quad (2.7)$$

form indeed a global 3-form. Just as the curvature 2-form F of a principal $U(1)$ -bundle captures the first Chern class⁹ $c_1(P) = [F]$, so the curvature 3-form H of a gerbe \mathcal{G} encodes its Dixmier-Douady class $dd(\mathcal{G}) = [H]$.

Geometrically, gerbes should be regarded as central groupoid extensions. That is, they are particular categories in which both the morphisms and the objects form manifolds. A rather detailed review of $U(1)$ -gerbes and the relation between the various definitions is found, e.g., in [30].

Just as principal $U(1)$ -bundles, gerbes can be tensored and restricted to submanifolds. Geometrically, this amounts to tensoring and restricting their categorified fibers. From the cocycle perspective, tensoring amounts to adding the underlying cocycles, yielding a gerbe whose Dixmier-Douady class is the sum of the Dixmier-Douady classes of the individual gerbes.¹⁰ A restriction to some submanifold is also readily performed by restricting the cocycles to that submanifold.¹¹

Let us now come to gerbes on the nilmanifolds N_j . The Dixmier-Douady classes of these are encoded in global closed 3-forms, and since N_j is connected, compact, orientable and 3-dimensional, these forms are unique up to rescaling. The cocycle conditions impose

⁹I.e. up to torsion elements.

¹⁰More explicitly, the two gerbes on M are described by cocycles with respect to two covers of M . There is a common refinement of these two covers, over which we can add the corresponding cocycles.

¹¹Explicitly, given a submanifold $X \hookrightarrow M$ of some manifold M , the cocycles on M are given with respect to some cover U_M . Together with the embedding $X \hookrightarrow M$, this cover induces a cover U_X on X and we can restrict the cocycles on M by pullback.

a quantization condition on the Dixmier-Douady class, just as it is the case for the Chern class for principal U(1)-bundles. Altogether, gerbes on N_j are characterized by an integer k , and we denote the corresponding gerbe by $\mathcal{G}_{j,k}$.

The 3-form representing the Dixmier-Douady class of $\mathcal{G}_{j,k}$ is the global 3-form

$$H = k dx^1 \wedge dx^2 \wedge dx^3, \quad \int_{N_j} H = k, \quad (2.8)$$

where we used again coordinates (2.2). Note that a 2-form B giving rise to this curvature is

$$B = k x^1 dx^2 \wedge dx^3. \quad (2.9)$$

Clearly, this 2-form is not global because $B(x^1 + 1) - B(x^1) = k dx^2 \wedge dx^3 \neq 0$. The cocycle conditions (2.6) then implies that

$$dA = k dx^2 \wedge dx^3. \quad (2.10)$$

2.3 Exact Courant algebroids and generalized metrics

An *exact Courant algebroid* over a manifold M is usually defined as a vector bundle E which fits into the short exact sequence of vector bundles,

$$0 \longrightarrow T^*M \hookrightarrow E \xrightarrow{\rho} TM \longrightarrow 0, \quad (2.11)$$

where $\rho : E \rightarrow TM$ is called the *anchor map*, and which is endowed with a fiberwise inner product $\langle -, - \rangle : E \times E \rightarrow M \times \mathbb{R}$ and a bracket $\mu_2(-, -) : \Gamma E \times \Gamma E \rightarrow \Gamma E$ called *Courant bracket*, subject to a number of axioms.

A *splitting* of the anchor map ρ , i.e. a bundle map $\sigma : TM \hookrightarrow E$ such that $\rho \circ \sigma = \text{id}$, yields the projector $P := \sigma \circ \rho$ and thus an isomorphism

$$E = \text{im}(P) \oplus \text{ker}(P) \cong TM \oplus T^*M. \quad (2.12)$$

Explicitly, consider the anchor map

$$\rho : E \cong TM \oplus T^*M \rightarrow TM, \quad \rho(X + \alpha) = X. \quad (2.13)$$

A general splitting is then of the form

$$\sigma : X \mapsto X + \omega(X) = X + (g + B)(X), \quad (2.14)$$

where we split the rank 2 tensor ω into its symmetric and antisymmetric parts g and B . It also leads to the two graphs

$$C_{\pm} = \{X + (\pm g + B)(X) \mid X \in \Gamma(TM)\}, \quad (2.15)$$

which span the subbundles of E on which the inner product

$$\langle X + \alpha, Y + \beta \rangle = \alpha(Y) + \beta(X) \quad (2.16)$$

is positive and negative definite, respectively. Such a half-dimensional subbundle C_+ of E with positive definite induced metric is called a *generalized metric*. More explicitly, let G be an endomorphism on E such that

$$C_{\pm} = \ker(G \mp \mathbb{1}) \quad \text{and} \quad G^2 = GG^T = \mathbb{1} \quad (2.17)$$

and let η be the $O(d, d)$ -metric

$$\eta = \begin{pmatrix} 0 & \mathbb{1} \\ \mathbb{1} & 0 \end{pmatrix} \quad \text{with} \quad \langle X + \alpha, Y + \beta \rangle = \begin{pmatrix} X \\ \alpha \end{pmatrix}^T \eta \begin{pmatrix} Y \\ \beta \end{pmatrix}. \quad (2.18)$$

Then the map $\mathcal{H} : \Gamma(E) \rightarrow \Gamma(E^*)$ with $\mathcal{H} = \eta G$ is also called the *generalized metric*. With C_{\pm} as given in (2.15) and (2.17), it follows that

$$G = \begin{pmatrix} -g^{-1}B & g^{-1} \\ g - Bg^{-1}B & Bg^{-1} \end{pmatrix} \quad \text{and} \quad \mathcal{H} = \begin{pmatrix} g - Bg^{-1}B & Bg^{-1} \\ -g^{-1}B & g^{-1} \end{pmatrix}. \quad (2.19)$$

2.4 Exact Courant algebroids as symplectic Lie 2-algebroids

A more powerful perspective is obtained by regarding a Courant algebroid as a symplectic Lie 2-algebroid \mathcal{C} or *symplectic NQ-manifold of degree 2* as explained in detail in [13]. That is, \mathcal{C} is an \mathbb{N} -graded manifold with body M , which carries a symplectic form of \mathbb{N} -degree 2 as well as a nilquadratic vector field Q of degree 1 generating a symplectomorphism. For a very detailed discussion of symplectic Lie n -algebroids in our conventions, see [10].

As shown in [13], the NQ-manifold underlying an exact Courant algebroid over a manifold M is isomorphic to $\mathcal{C} = T^*[2]T[1]M$. This NQ-manifold is essentially the vector bundle T^*TM , but regarded as a graded manifold with local coordinates $x^\mu, \xi^\mu, \zeta_\mu, p_\mu$ of degrees 0, 1, 1 and 2, respectively. Locally, the functions on \mathcal{C} are polynomials in the fiber coordinates with coefficients which are arbitrary functions of the x^μ . There is a canonical symplectic form and a canonical vector field of degree 1 in the above coordinates,

$$\omega = dx^\mu \wedge dp_\mu + d\xi^\mu \wedge d\zeta_\mu, \quad Q = \xi^\mu \frac{\partial}{\partial x^\mu} + p_\mu \frac{\partial}{\partial \zeta_\mu} \quad (2.20)$$

and Q satisfies indeed $Q^2 = 0$ and $\mathcal{L}_Q \omega = 0$. The symplectic form ω induces a Poisson structure,¹²

$$\begin{aligned} \{f, g\} := & \left(\frac{\partial}{\partial p_\mu} f \right) \left(\frac{\partial}{\partial x^\mu} g \right) - \left(\frac{\partial}{\partial x^\mu} f \right) \left(\frac{\partial}{\partial p_\mu} g \right) \\ & - (-1)^{|f|} \left(\frac{\partial}{\partial \zeta_\mu} f \right) \left(\frac{\partial}{\partial \xi^\mu} g \right) - (-1)^{|f|} \left(\frac{\partial}{\partial \xi^\mu} f \right) \left(\frac{\partial}{\partial \zeta_\mu} g \right), \end{aligned} \quad (2.21)$$

and Q is Hamiltonian:

$$Q = \{Q, -\} \quad \text{with} \quad Q = \xi^\mu p_\mu. \quad (2.22)$$

¹²See again [10] for details and our conventions.

The relation between the classical definition of an exact Courant algebroid and the NQ-manifold picture is now as follows. Putting $p_\mu = 0$ on \mathcal{C} yields a vector bundle over M with fibers of degree 1, $E[1] \rightarrow M$. Here, $E[1]$ is the grade-shifted bundle E in (2.11). Functions $\mathcal{C}^\infty(\mathcal{C})$ of degree 1 correspond to sections of E and in terms of our local coordinates, we identify

$$X + \alpha = X^\mu(x)\zeta_\mu + \alpha_\mu(x)\xi^\mu \quad \longleftrightarrow \quad X^\mu(x)\frac{\partial}{\partial x^\mu} + \alpha_\mu(x)dx^\mu \in \Gamma(TM \oplus T^*M) . \quad (2.23)$$

The kernel of the anchor map are the sections spanned by the coefficient of p_μ in \mathcal{Q} , which is here ξ^μ . Explicitly, the anchor map reads as $\rho(X^\mu\zeta_\mu) = X^\mu\frac{\partial}{\partial x^\mu}$. The pairing on E is given by the Poisson bracket (2.21) on functions of degree 1 on \mathcal{C} and the Courant bracket,

$$\mu_2(X_1 + \alpha_1, X_2 + \alpha_2) = \frac{1}{2} (\{ \{ \mathcal{Q}, X_1 + \alpha_1 \}, X_2 + \alpha_2 \} - \{ \{ \mathcal{Q}, X_2 + \alpha_2 \}, X_1 + \alpha_1 \}) , \quad (2.24)$$

is part of an associated Lie 2-algebra that canonically arises on \mathcal{C} from the homological vector field \mathcal{Q} . We shall discuss this Lie 2-algebra in more detail in section 3.

2.5 The Courant algebroid of a gerbe

Consider now a gerbe \mathcal{G} with connective structure $(A_{(ab)}, B_{(a)})$ with respect to a cover $\sqcup_a U_a \rightarrow M$ of M and let $H = dB \in H^3(M)$ be a representative of its Dixmier-Douady class $dd(\mathcal{G})$. This data induces local splits σ of the anchor map $\rho : E[1] \rightarrow TM$, given over the patch U_a by

$$\sigma \left(\frac{\partial}{\partial x^\mu} \right) := \frac{\partial}{\partial x^\mu} + B_{\mu\nu} dx^\nu , \quad B_{(a)} = \frac{1}{2} B_{\mu\nu} dx^\mu \wedge dx^\nu . \quad (2.25)$$

If we further insert a symmetric tensor, $g_{\mu\nu}$,

$$\sigma \left(\frac{\partial}{\partial x^\mu} \right) := \frac{\partial}{\partial x^\mu} + (g_{\mu\nu} + B_{\mu\nu}) dx^\nu , \quad (2.26)$$

we can capture any splitting, as discussed above. For now, however, we restrict ourselves to splittings of the form (2.25).

In the symplectic NQ-picture, equation (2.25) translates to maps

$$\sigma : TU_a \rightarrow \mathcal{C}_1^\infty(T^*[2]T[1]U_a) , \quad \sigma \left(\frac{\partial}{\partial x_{(a)}^\mu} \right) := \zeta_\mu^{(a)} + B_{\mu\nu}^{(a)} \xi_{(a)}^\nu . \quad (2.27)$$

On the intersection of patches, we then have the relation

$$\zeta_\mu^{(a)} = \zeta_\mu^{(b)} + \partial_{[\mu} A_{\nu]}^{(ab)} \xi^\nu , \quad (2.28)$$

because of the Deligne cocycle relation $B_{(a)} - B_{(b)} = dA_{(ab)}$, cf. (2.6). In terms of sections $X + \alpha$ of $TM \oplus T^*M$, this relation corresponds to the well-known patching relation for the *twisted generalized tangent bundle*,

$$X_{(a)} + \alpha_{(a)} = X_{(b)} + \alpha_{(b)} + \iota_{X_{(b)}} dA_{ab} . \quad (2.29)$$

The coordinate change $\zeta_\mu \rightarrow \zeta_\mu + B_{\mu\nu}\xi^\nu$ induced by the splitting is still compatible with the original symplectic form and homological vector field. (The necessary relations, $Q^2 = 0$ and $\mathcal{L}_Q\omega = 0$ are local conditions, and Q and ω are now simply glued together in a different way.) One can, however, perform a coordinate change in which the B -field contribution to the coordinate transformations is removed but instead one has a change in the homological vector field Q , which is best expressed by a modification of its Hamiltonian \mathcal{Q} ,

$$\mathcal{Q} = \xi^\mu p_\mu + \frac{1}{3!} H_{\mu\nu\kappa} \xi^\mu \xi^\nu \xi^\kappa . \tag{2.30}$$

Here, $H = dB = \frac{1}{3!} H_{\mu\nu\kappa} dx^\mu \wedge dx^\nu \wedge dx^\kappa$. We shall see an explicit example of this transformation later. In both cases, we arrive at the *twisted Courant algebroid* \mathcal{C}_g .

Note that in (2.30), the condition for $\{Q, \mathcal{Q}\} = 0$, which is equivalent to $Q^2 = 0$, is $dH = 0$. Moreover, one can show that the additional contributions from the B -field can be transformed away by a symplectomorphism on \mathcal{C} if and only if H is exact, see e.g. [10]. Similarly, splittings (2.26) can be brought to the form (2.25) by a symplectomorphism. Thus, relevant for the twist is indeed the cohomology class of $[H]$, the Dixmier-Douady class $dd(\mathcal{G})$ of the gerbe \mathcal{G} . We conclude that the isomorphism classes of Courant algebroids over M are determined by $H^3(M)$ and the corresponding characteristic class is called the *Ševera class* of the Courant algebroid.

We thus obtain a one-to-one correspondence between Courant algebroids \mathcal{C}_g and gerbes \mathcal{G} , up to isomorphism classes, mapping the Dixmier-Douady class to the Ševera class. As discussed in [10], the Courant algebroid \mathcal{C}_g comes with an associated Lie 2-algebra which describes the symmetries of the gerbe \mathcal{G} . In this sense, the Courant algebroid generalizes the Atiyah algebroid of principal $U(1)$ -bundles, cf. [31, 32].

In the twisted case, a *generalized metric* is still a half-dimensional subbundle C_+ of E with positive definite induced metric. Over patches U_a , we still have

$$C_\pm|_{U_a} = \{X + (\pm g + B)(X) | X \in \Gamma(TU_a)\}, \tag{2.31}$$

but E and its subbundles C_\pm will have non-trivial transition functions, in general. This leads to the endomorphism G and the generalized metric to be defined only locally:

$$G_{(a)} = \begin{pmatrix} -g^{-1}B_{(a)} & g^{-1} \\ g - B_{(a)}g^{-1}B_{(a)} & B_{(a)}g^{-1} \end{pmatrix}, \quad \mathcal{H}_{(a)} = \begin{pmatrix} g - B_{(a)}g^{-1}B_{(a)} & B_{(a)}g^{-1} \\ -g^{-1}B_{(a)} & g^{-1} \end{pmatrix}. \tag{2.32}$$

Decomposing $\mathcal{H}_{(a)}$ as

$$\mathcal{H}_{(a)} = \begin{pmatrix} \mathbb{1} & B_{(a)} \\ 0 & \mathbb{1} \end{pmatrix} \begin{pmatrix} g & 0 \\ 0 & g^{-1} \end{pmatrix} \begin{pmatrix} \mathbb{1} & B_{(a)} \\ 0 & \mathbb{1} \end{pmatrix}^T, \tag{2.33}$$

we readily verify that

$$\mathcal{H}_{(a)} = \begin{pmatrix} \mathbb{1} & dA_{(ab)} \\ 0 & \mathbb{1} \end{pmatrix} \mathcal{H}_{(b)} \begin{pmatrix} \mathbb{1} & dA_{ab} \\ 0 & \mathbb{1} \end{pmatrix}^T \tag{2.34}$$

due to $B_{(a)} = B_{(b)} + dA_{(ab)}$ and we obtain the transition function for the generalized metric.

2.6 The Courant algebroid of $\mathcal{G}_{j,k}$

We now give the explicit Courant algebroid $\mathcal{C}_{j,k} := \mathcal{C}_{\mathcal{G}_{j,k}}$ for the gerbe $\mathcal{G}_{j,k}$ over the nilmanifold N_j with Dixmier-Douady class k . We use again the periodic coordinates x^i , $i = 1, 2, 3$, with the identification

$$(x^1, x^2, x^3) \sim (x^1, x^2 + 1, x^3) \sim (x^1, x^2, x^3 + 1) \sim (x^1 + 1, x^2, x^3 - jx^2). \quad (2.35)$$

Let us first consider the case $k = 0$. We have much freedom in choosing the coordinates ξ^i and ζ_i of degree 1 on $E[1]$ and it is sensible to consider Darboux coordinates, for which the symplectic form reads as

$$\omega = dx^i \wedge dp_i + d\xi^i \wedge d\zeta_i. \quad (2.36)$$

By considering patches on N_j , we can use the exact Courant algebroid as described above and glue it together over patches. This leads to the coordinate system $(x^i, \tilde{\xi}^i, \tilde{\zeta}_i, \tilde{p}_i)$ with the anchor map and embedding

$$\begin{aligned} \tilde{\zeta}_1 &\mapsto \frac{\partial}{\partial x^1}, & \tilde{\zeta}_2 &\mapsto \frac{\partial}{\partial x^2}, & \tilde{\zeta}_3 &\mapsto \frac{\partial}{\partial x^3}, \\ \tilde{\xi}^1 &\leftrightarrow dx^1, & \tilde{\xi}^2 &\leftrightarrow dx^2, & \tilde{\xi}^3 &\leftrightarrow dx^3. \end{aligned} \quad (2.37)$$

Using equation 3.3 of [13], we then deduce that the identification of coordinates is

$$\begin{aligned} &(x^1, x^2, x^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3) \\ &= (x^1 + 1, x^2, x^3 - jx^2, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3 - j\tilde{\xi}^2, \tilde{\zeta}_1, \tilde{\zeta}_2 + j\tilde{\zeta}_3, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2 + j\tilde{p}_3, \tilde{p}_3). \end{aligned} \quad (2.38)$$

The symplectic form ω and the Hamiltonian \mathcal{Q}_0 are the canonical ones,

$$\omega = dx^i \wedge d\tilde{p}_i + d\tilde{\xi}^i \wedge d\tilde{\zeta}_i \quad \text{and} \quad \mathcal{Q}_0 = \tilde{\xi}^i \tilde{p}_i. \quad (2.39)$$

As necessary, both expressions are indeed global.

Much more convenient, however, is the coordinate system $(x^i, \bar{\xi}^i, \bar{\zeta}_i, \bar{p}_i)$, where we use the global vector fields and dual 1-forms (2.3). The anchor map thus reads as

$$\bar{\zeta}_1 \mapsto \frac{\partial}{\partial x^1}, \quad \bar{\zeta}_2 \mapsto \frac{\partial}{\partial x^2} - jx^1 \frac{\partial}{\partial x^3}, \quad \bar{\zeta}_3 \mapsto \frac{\partial}{\partial x^3} \quad \text{or} \quad \bar{\zeta}_i \mapsto \tau_i^j \frac{\partial}{\partial x^j}, \quad (2.40)$$

where we introduced the vielbein

$$\tau_i^j = \delta_i^j - jx^1 \delta_i^2 \delta_3^j, \quad (2.41)$$

not making any distinction between coordinate and non-coordinate indices. Since $\bar{\xi}^i$ and $\bar{\zeta}_i$ are supposed to be Darboux coordinates, the embedding $T^*M \hookrightarrow E$ identifies

$$\bar{\xi}^1 \leftrightarrow dx^1, \quad \bar{\xi}^2 \leftrightarrow dx^2, \quad \bar{\xi}^3 \leftrightarrow dx^3 + jx^1 dx^2 \quad \text{or} \quad \bar{\xi}^j \leftrightarrow (\tau^{-1})_i^j dx^i. \quad (2.42)$$

Together with $\bar{p}_i = \tilde{p}_i$, we arrive at the following identification of coordinates:

$$\begin{aligned} &(x^1, x^2, x^3, \bar{\xi}^1, \bar{\xi}^2, \bar{\xi}^3, \bar{\zeta}_1, \bar{\zeta}_2, \bar{\zeta}_3, \bar{p}_1, \bar{p}_2, \bar{p}_3) \\ &= (x^1 + 1, x^2, x^3 - jx^2, \bar{\xi}^1, \bar{\xi}^2, \bar{\xi}^3, \bar{\zeta}_1, \bar{\zeta}_2, \bar{\zeta}_3, \bar{p}_1, \bar{p}_2 + j\bar{p}_3, \bar{p}_3). \end{aligned} \quad (2.43)$$

The appropriate global symplectic form ω and Hamiltonian \mathcal{Q}_0 are then

$$\omega = dx^i \wedge d\bar{p}_i + d\bar{\xi}^i \wedge d\bar{\zeta}_i \quad \text{and} \quad \mathcal{Q}_0 = \bar{\xi}^i \bar{p}_i - jx^1 \bar{\xi}^2 \bar{p}_3 = \bar{\xi}^i \tau_i^j \bar{p}_j, \quad (2.44)$$

To include a non-trivial Dixmier-Douady class $dd(\mathcal{G}) = k$ of a gerbe \mathcal{G} , we have to modify the coordinates $\bar{\xi}^i$ and $\bar{\zeta}_i$, taking into account the gluing condition (2.29). From (2.10) and (2.28), we glean that the transition from x^1 to $x^1 + 1$ should read as

$$X^i \bar{\zeta}_i + \alpha_i \bar{\xi}^i \quad \longrightarrow \quad X^i \bar{\zeta}_i + \alpha_i \bar{\xi}^i + k(X^2 \bar{\xi}^3 - X^3 \bar{\xi}^2). \quad (2.45)$$

Therefore, we have the identification

$$(x^1, \bar{\xi}^1, \bar{\xi}^2, \bar{\xi}^3, \bar{\zeta}_1, \bar{\zeta}_2, \bar{\zeta}_3) \sim (x^1 + 1, \bar{\xi}^1, \bar{\xi}^2, \bar{\xi}^3, \bar{\zeta}_1, \bar{\zeta}_2 + k\bar{\xi}^3, \bar{\zeta}_3 - k\bar{\xi}^2). \quad (2.46)$$

We first observe that there is no additional periodicity condition for the momenta from (2.45), but this is due to the fact that the transformation does not depend on the body coordinates. In the general case, also the momenta receive an additional periodicity condition due to the topological non-triviality of the gerbe. Both, ω and \mathcal{Q}_0 remain invariant under the combined (2.43) and (2.45) transformations. Hence, the expressions are globally defined and since $\{\mathcal{Q}_0, \mathcal{Q}_0\} = 0$ as before, they complete the structure of a symplectic NQ-manifold and therefore the structure of a Courant algebroid. Note that we can compute a generalized metric from the natural metric on the nilmanifold (2.4) and the B -field for the gerbe $\mathcal{G}_{j,k}$, (2.9). The result is

$$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & (1 + j^2(x^1)^2)(1 + k^2(x^1)^2) & jx^1(1 + k^2(x^1)^2) & 0 & -jk(x^1)^2 & kx^1(1 + j^2(x^1)^2) \\ 0 & jx^1(1 + k^2(x^1)^2) & 1 + k^2(x^1)^2 & 0 & -kx^1 & jk(x^1)^2 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & -jk(x^1)^2 & -kx^1 & 0 & 1 & -jx^1 \\ 0 & kx^1(1 + j^2(x^1)^2) & jk(x^1)^2 & 0 & -jx^1 & 1 + j^2(x^1)^2 \end{pmatrix}. \quad (2.47)$$

In the transition $x^1 \rightarrow x^1 + 1$, this generalized metric transforms indeed as (2.34).

Our coordinates of degree 1, $\bar{\xi}^i$ and $\bar{\zeta}_i$ are no longer global, as a non-trivial B -field induces the non-trivial gluing of the generalized tangent bundle (2.29). We certainly prefer, however, to work with coordinates $(x^i, \xi^i, \zeta_i, p_i)$, where the degree 1 members are global. It is well-known in Generalized Geometry, that the non-trivial gluing can be absorbed by performing a B -field transformation on the sections of $TM \oplus T^*M$, cf. e.g. [33]. This amounts to

$$(\xi, \zeta) = e^B(\bar{\xi}, \bar{\zeta}) := e^{\{-, B\}}(\bar{\xi}, \bar{\zeta}) = \left(1 - kx^1 \bar{\xi}^3 \frac{\partial}{\partial \bar{\zeta}_2} + kx^1 \bar{\xi}^2 \frac{\partial}{\partial \bar{\zeta}_3}\right)(\bar{\xi}, \bar{\zeta}) \quad (2.48a)$$

on the basis vectors. In terms of coordinates on the symplectic NQ-manifold, we have

$$x^i = \bar{x}^i, \quad \xi^i = \bar{\xi}^i, \quad \zeta_1 = \bar{\zeta}_1, \quad \zeta_2 = \bar{\zeta}_2 - kx^1 \bar{\xi}^3, \quad \zeta_3 = \bar{\zeta}_3 + kx^1 \bar{\xi}^2, \quad (2.48b)$$

which is complemented by

$$p_1 = \bar{p}_1 - k\bar{\xi}^2 \bar{\xi}^3, \quad p_2 = \bar{p}_2, \quad p_3 = \bar{p}_3 \quad (2.48c)$$

to a symplectomorphism, cf. again equation 3.3 of [13]:

$$\omega = dx^i \wedge dp_i + d\xi^i \wedge d\zeta_i . \tag{2.49}$$

For the new coordinate system, we have the identification

$$\begin{aligned} & (x^1, x^2, x^3, \xi^1, \xi^2, \xi^3, \zeta_1, \zeta_2, \zeta_3, p_1, p_2, p_3) \\ & = (x^1 + 1, x^2, x^3 - jx^2, \xi^1, \xi^2, \xi^3, \zeta_1, \zeta_2, \zeta_3, p_1, p_2 + jp_3, p_3) . \end{aligned} \tag{2.50}$$

In the following, we shall use essentially exclusively the coordinates $(x^i, \xi^i, \zeta_i, p_i)$.

The Hamiltonian in the new coordinates reads as

$$\mathcal{Q} = \xi^i p_i - jx^1 \xi^2 p_3 + k \xi^1 \xi^2 \xi^3 \tag{2.51}$$

and we see that the dependence on the Dixmier-Douady class which was previously encoded in the transition function of the coordinates is now moved into the Hamiltonian \mathcal{Q} . As one readily shows, there is no symplectomorphism eliminating k , and the Ševera class k describes the isomorphism classes of the symplectic NQ -manifold $\mathcal{C}_{j,k}$.

As a side remark, we observe that \mathcal{Q} is obtained from \mathcal{Q}_0 by acting with $e^B := e^{\{-, B\}}$ with $B = \frac{1}{2} B_{\mu\nu} \xi^\mu \xi^\nu$: $\mathcal{Q} = e^B \mathcal{Q}_0$. This can be used to relate \mathcal{Q}_0 -closed functions which represent lifts of generalized vectors (i.e. functions of degree 1) to \mathcal{Q} -closed lifts of generalized vectors: if $X \in \mathcal{C}_1^\infty(\mathcal{C}_{j,k})$ is \mathcal{Q}_0 -closed, then $e^B X$ is \mathcal{Q} -closed,

$$\begin{aligned} \{\mathcal{Q}, X\} &= \{e^B \mathcal{Q}_0, e^B X\} \\ &= \{\mathcal{Q}_0, X\} + \left\{ \{\mathcal{Q}_0, B\}, X \right\} + \left\{ \mathcal{Q}_0, \{X, B\} \right\} \\ &= \{\mathcal{Q}_0, X\} + \left\{ \{\mathcal{Q}_0, X\}, B \right\} \\ &= e^B \{\mathcal{Q}_0, X\} . \end{aligned} \tag{2.52}$$

This result is well-known in the classical language: it relates ordinary and twisted cohomology on the generalized tangent bundle [33].

3 Finite global symmetries of Generalized Geometry

One of the major advantages of the above formulation of Generalized Geometry is that it makes the relevant symmetry structures very obvious. As first observed in [16, 17], symplectic Lie n -algebroids always come with an associated Lie n -algebra structure, see also [10] for more details. The symmetry Lie 2-algebra of the gerbe $\mathcal{G}_{j,k}$ is given by the associated Lie 2-algebra of the Courant algebroid $\mathcal{C}_{j,k}$, cf. [10].¹³ The Courant algebroid is to the abelian gerbe roughly what the Atiyah algebroid is to a principal circle bundle, see e.g. [31, 32]. We start with a description of the Lie 2-algebra associated to $\mathcal{C}_{j,k}$, before we interpret it as the appropriate symmetry Lie 2-algebra.

¹³This relation readily extends to n -gerbes.

3.1 Lie 2-algebra associated to $\mathcal{C}_{j,k}$

The Lie 2-algebra associated to $\mathcal{C}_{j,k}$ has an underlying graded vector space $\mathbf{L} = \mathbf{L}_{-1} \oplus \mathbf{L}_0$, where the homogeneous parts

$$\mathbf{L}_{-1} = \mathcal{C}_0^\infty(\mathcal{C}_{j,k}) = \mathcal{C}^\infty(N_j) \quad \text{and} \quad \mathbf{L}_0 = \mathcal{C}_1^\infty(\mathcal{C}_{j,k}) \cong \Gamma(E), \quad (3.1)$$

are the vector spaces of functions on $\mathcal{C}_{j,k}$ of degree 0 and 1, respectively. These correspond to ordinary functions on the base N_j of the Courant algebroid as well as generalized vectors. The graded vector space \mathbf{L} becomes a Lie 2-algebra (a 2-term L_∞ -algebra) if we endow it with the following totally anti-symmetric brackets μ_i of degree $2 - i$:

$$\begin{aligned} \mu_1(f) &= \{\mathcal{Q}, f\}, \\ \mu_2(X + \alpha, Y + \beta) &= \frac{1}{2}(\{\{\mathcal{Q}, X^i \zeta_i + \alpha_i \xi^i\}, Y^\nu \zeta_\nu + \beta_\nu \xi^\nu\} - (X + \alpha) \leftrightarrow (Y + \beta)) \\ \mu_2(X + \alpha, f) &= \frac{1}{2}\{\{\mathcal{Q}, X + \alpha\}, f\}, \\ \mu_3(X + \alpha, Y + \beta, Z + \gamma) &= \frac{1}{3!}(\{\{\{\mathcal{Q}, X + \alpha\}, Y + \beta\}, Z + \gamma\} + \dots) \\ &= \frac{1}{3!}(\{X + \alpha, \mu_2(Y + \beta, Z + \gamma)\} + \text{cycl.}) . \end{aligned} \quad (3.2)$$

The explicit form of (3.2) now depends on the chosen coordinate system since ordinary 1-forms and vector fields are no longer global objects. We first note, that in the holonomic frame $(\tilde{\xi}^i, \tilde{\zeta}_i)$, the expressions are just the ones given in [10], i.e.

$$\begin{aligned} \mu_1(f) &= df, \\ \mu_2(X + \alpha, Y + \beta) &= [X, Y] + \mathcal{L}_X \beta - \mathcal{L}_Y \alpha - \frac{1}{2}d(\iota_X \beta - \iota_Y \alpha) + \iota_X \iota_Y H, \\ \mu_2(X + \alpha, f) &= \mathcal{L}_X f, \\ \mu_3(X + \alpha, Y + \beta, Z + \gamma) &= \frac{1}{3!}(\iota_X \iota_Y d\gamma + \frac{3}{2}\iota_X d\iota_Y \gamma \pm \text{perm.}) + \iota_X \iota_Y \iota_Z H . \end{aligned} \quad (3.3)$$

In the nonholonomic frame (ξ^i, ζ_i) , the expressions look formally the same if we use the vielbein: the homological function takes the form $\mathcal{Q} = \xi^i \tau_i^j p_j + k \xi^1 \xi^2 \xi^3$, so we get a new bracket $[\cdot, \cdot]^\tau$ and differential d^τ given for vector fields $X = X^i \zeta_i$, $Y = Y^i \zeta_i$ and forms $\omega = \frac{1}{k!} \omega_{i_1 \dots i_k} \xi^{i_1} \dots \xi^{i_k}$ by

$$\begin{aligned} [X, Y]^\tau &= (X^i \tau_i^j \partial_j Y^k - Y^i \tau_i^j \partial_j X^k) \zeta_k, \\ d^\tau \omega &= \frac{1}{k!} \tau_{i_0}^j \partial_j \omega_{i_1 \dots i_k} \xi^{i_0} \dots \xi^{i_k}, \end{aligned} \quad (3.4)$$

where ∂_i is short for $\frac{\partial}{\partial x^i}$ and a totally antisymmetrized sum is taken over the indices $i_0 \dots i_k$ in the last expression. Using d^τ and $[\cdot, \cdot]^\tau$ together with the H -flux written in the coordinates ξ^i we arrive at the expressions

$$\begin{aligned} \mu_1(f) &= d^\tau f, \\ \mu_2(X + \alpha, Y + \beta) &= [X, Y]^\tau + \mathcal{L}_X^\tau \beta - \mathcal{L}_Y^\tau \alpha - \frac{1}{2}d^\tau(\iota_X \beta - \iota_Y \alpha) + \iota_X \iota_Y H, \\ \mu_2(X + \alpha, f) &= \mathcal{L}_X^\tau f, \\ \mu_3(X + \alpha, Y + \beta, Z + \gamma) &= \frac{1}{3!}(\iota_X \iota_Y d^\tau \gamma + \frac{3}{2}\iota_X d^\tau \iota_Y \gamma \pm \text{perm.}) + \iota_X \iota_Y \iota_Z H . \end{aligned} \quad (3.5)$$

For our purposes, it will be useful to generalize the underlying graded vector space of \mathbb{L} from global to local functions. Given a cover $\sqcup_a U_a \rightarrow N_j$, this corresponds to extending \mathbb{L} to $\mathbb{L}^{\text{loc}} = \mathbb{L}_{-1}^{\text{loc}} \oplus \mathbb{L}_0^{\text{loc}}$ with

$$\mathbb{L}_{-1}^{\text{loc}} = \sqcup_a \mathcal{C}^\infty(U_a) \quad \text{and} \quad \mathbb{L}_0^{\text{loc}} = \sqcup_a \mathcal{C}_1^\infty(\mathcal{C}_{j,k}|_{U_a}) \cong \sqcup_a \Gamma(E|_{U_a}). \quad (3.6)$$

Since the brackets only involve local expressions (i.e. expressions with finitely many derivatives), the higher products μ_i are not affected by this extension. Since the brackets are given by tensorial expressions, the brackets will be compatible with the gluing conditions.

3.2 Action of the associated Lie algebra and its action Lie algebra

As explained in [10], there is a well-defined general notion of an action of an L_∞ -algebra on a graded vector space and associated L_∞ -algebras of symplectic L_∞ -algebroids naturally come with such an action.

In our case, the Lie 2-algebras \mathbb{L} and \mathbb{L}^{loc} act on the functions and the local functions on the Courant algebroid $\mathcal{C}_{j,k}$ via the Dorfman bracket,¹⁴

$$\begin{aligned} \nu_2(X + \alpha, Y + \beta) &= \{\{Q, X + \alpha\}, Y + \beta\} \\ &= X^i \tau_i^j \partial_j (Y + \beta) + \bar{\xi}^i \tau_i^j (\partial_j X^j) \beta_j - Y^k \tau_k^i \partial_i (X + \alpha) + Y^k \bar{\xi}^i \tau_i^j \partial_j \alpha_k \\ &= [X, Y]^\tau + \iota_X d^\tau \beta + d^\tau (\iota_X \beta) - \iota_Y d^\tau \alpha, \end{aligned} \quad (3.7)$$

where $X + \alpha, Y + \beta \in \mathcal{C}_1^\infty(\mathcal{C}_{j,k})$. This action is readily extended to an action of \mathbb{L}^{loc} on local functions. We can, in fact, regard the Dorfman bracket as a generalized Lie derivative

$$\hat{\mathcal{L}}_{X+\alpha}(Y + \beta) = \{\{Q, X + \alpha\}, Y + \beta\}, \quad (3.8)$$

as is customary in Generalized Geometry and Double Field Theory. We note that this action together with the commutator

$$\begin{aligned} [\hat{\mathcal{L}}_{X_1+\alpha_1}, \hat{\mathcal{L}}_{X_2+\alpha_2}](Y + \beta) &= \hat{\mathcal{L}}_{X_1+\alpha_1} \hat{\mathcal{L}}_{X_2+\alpha_2} Y - \hat{\mathcal{L}}_{X_2+\alpha_2} \hat{\mathcal{L}}_{X_1+\alpha_1} (Y + \beta) \\ &= \hat{\mathcal{L}}_{\mu_2(X_1+\alpha_1, X_2+\alpha_2)}(Y + \beta) \end{aligned} \quad (3.9)$$

forms an ordinary Lie algebra¹⁵ $\text{Lie}(\mathbb{L}^{\text{loc}}, \nu_2)$ with “structure constants” induced by the Courant bracket μ_2 . In particular, the action of the Jacobiator of the Lie 2-algebra is trivial.¹⁶

We can readily identify this Lie algebra as a Lie subalgebra of the symplectomorphisms on $\mathcal{C}_{j,k}$, since the generalized Lie derivative (3.8) is simply the action of a Hamiltonian vector field with Hamiltonian $\{Q, X + \alpha\}$ on the function $Y + \beta$. Moreover, we can restrict

¹⁴We note that the Courant and Dorfman brackets on nilmanifolds were also given, e.g., in [34]. The NQ -manifold picture goes beyond this since it provides also the full Lie 2-algebra of symmetries.

¹⁵The Jacobi identity is trivially satisfied for commutators.

¹⁶This is obvious from a more abstract point of view: the image of μ_1 describes the action of gauge transformations between gauge transformations (level 1) on gauge transformations (level 0). The former always leave the gauge fields invariant, see e.g. the explanation of the BRST complex in [27].

ourselves to the degree-preserving symplectomorphisms on $\mathcal{C}_{j,k}$. Such symplectomorphisms have been discussed, e.g., in [13], and they have general Hamiltonian of degree 2,

$$\mathcal{H} = \mathcal{H}^i p_i + \frac{1}{2} \mathcal{H}_{ij} \xi^i \xi^j + \mathcal{H}_i^j \xi^i \zeta_j + \frac{1}{2} \mathcal{H}^{ij} \zeta_i \zeta_j . \quad (3.10)$$

The case at hand, $\mathcal{H} = \{\mathcal{Q}, X + \alpha\}$, corresponds to

$$\mathcal{H}^i = X^i, \quad \mathcal{H}_{ij} = \partial_i \alpha_j - \partial_j \alpha_i, \quad \mathcal{H}_i^j = \partial_i X^j \quad \text{and} \quad \mathcal{H}^{ij} = 0 . \quad (3.11)$$

To be more specific about the structure of the Lie algebra, let us consider the explicit form of (3.9), first for $H = 0$:

$$\begin{aligned} \hat{\mathcal{L}}_{\mu_2(X_1+\alpha_1, X_2+\alpha_2)}(Y + \beta) &= [[X_1, X_2], Y] + \iota_{[X_1, X_2]} d^\tau \beta + d^\tau \iota_{[X_1, X_2]} \beta - \\ &\quad - \iota_Y d^\tau (\mathcal{L}_{X_1}^\tau \alpha_2 - \mathcal{L}_{X_2}^\tau \alpha_1) . \end{aligned} \quad (3.12)$$

The Lie algebra of actions of $\hat{\mathcal{L}}_X$ is thus clearly

$$\text{Lie}(\mathbb{L}^{\text{loc}}, \nu_2) = \mathfrak{diff}(N_j) \ltimes \Omega_{\text{closed}}^2(N_j), \quad (3.13)$$

where the semidirect product is given by the natural action of the infinitesimal diffeomorphisms on 2-forms. Note that the d^τ -closed forms are locally exact and give rise to the local 1-forms which correspond to part of the local functions $\alpha_{1,2}$ of degree 1 on $\mathcal{C}_{j,k}$.

To understand the case $H \neq 0$, we recall the observation made around (2.48a) where we replaced the local sections $Y + \beta$ by the appropriate B -field corrected expression

$$\begin{aligned} \widetilde{Y + \beta} &:= e^B(Y + \beta) = Y + \beta + \iota_Y B \\ &= e^{\{-, B\}}(Y + \beta) = Y + \beta + \{Y, B\}, \end{aligned} \quad (3.14)$$

where we used $B = \frac{1}{2} B_{\mu\nu} \xi^\mu \xi^\nu$ in the Poisson brackets. Note that also

$$\{\widetilde{X + \alpha}, \widetilde{Y + \beta}\} = \{X + \alpha, Y + \beta\} . \quad (3.15)$$

The modification of \mathcal{Q} due to a non-trivial B -field is merely its covariantization with respect to the e^B action. Denoting the B -field modified Hamiltonian by $\tilde{\mathcal{Q}}$ we have

$$\tilde{\mathcal{Q}} = \mathcal{Q} + dB = \mathcal{Q} + \{\mathcal{Q}, B\} = e^{\{-, B\}} \mathcal{Q} \quad (3.16)$$

as well as

$$\begin{aligned} \{\tilde{\mathcal{Q}}, \widetilde{Y + \beta}\} &= \{e^{\{-, B\}} \mathcal{Q}, e^{\{-, B\}}(Y + \beta)\} \\ &= \{\mathcal{Q} + \{\mathcal{Q}, B\}, Y + \beta + \{Y, B\}\} \\ &= \{\mathcal{Q}, Y + \beta\} + \{\mathcal{Q}, \{Y, B\}\} + \{\{\mathcal{Q}, B\}, Y\} \\ &= \{\mathcal{Q}, Y + \beta\} + \{\{\mathcal{Q}, Y\}, B\} \\ &= e^{\{-, B\}} \{\mathcal{Q}, Y + \beta\} \\ &= \{\tilde{\mathcal{Q}}, \widetilde{Y + \beta}\} \end{aligned} \quad (3.17)$$

and similarly

$$\{\{\mathcal{Q}, \widetilde{X + \alpha}\}, \widetilde{Y + \beta}\} = \{\{\tilde{\mathcal{Q}}, \widetilde{X + \alpha}\}, \widetilde{Y + \beta}\} \quad (3.18)$$

Translating from the action on Y to the action on $e^B Y$ therefore removes the additional terms containing B in the Lie 2-algebra and its action. We thus recover again (3.13).

3.3 Finite symmetries of a gerbe

The symmetries of the gerbe¹⁷ $\mathcal{G}_{j,k}$ are given by diffeomorphisms on the base manifold N_j as well as gauge transformations. The latter are parameterized by local 1-forms α , but only their differential acts non-trivial, shifting the 2-form connective structure

$$B \rightarrow \tilde{B} = B + \delta B \quad \text{with} \quad \delta B = d^\tau \alpha = \{\mathcal{Q}, \alpha\} . \quad (3.19)$$

Therefore, there are gauge transformations between gauge transformations, parameterized by functions f on N_j , and changing the 1-forms α according to

$$\alpha \rightarrow \tilde{\alpha} = \alpha + \delta \alpha \quad \text{with} \quad \delta \alpha = d^\tau f = \{\mathcal{Q}, f\} . \quad (3.20)$$

The graded vector space of both these transformations arranges now into the categorified Lie algebra \mathbb{L}^{loc} .

On the Courant algebroid $\mathcal{C}_{j,k}$ whose Ševera class is the Dixmier-Douady class of the gerbe, this associated Lie 2-algebra also acts on graded functions via the Dorfman bracket ν_2 . This is relevant, as a subset of these functions can be interpreted as extended vector and polyvector fields (as formalized, e.g., in [10]). Together with the commutator, these actions form the Lie algebra (3.13), which is readily integrated to the group

$$\text{Diff}(N_j) \times \Omega_{\text{closed}}^2(N_j) . \quad (3.21)$$

This is also the symmetry group of the Courant algebroid $\mathcal{C}_{j,k}$, the automorphism group of the underlying bundle endowed with the Courant bracket. We note that the fact that the symmetries of a general Courant bracket are diffeomorphisms and closed 2-forms on the base manifold of the corresponding Courant bracket is well-known, cf. e.g. [12].

4 Global Double Field Theory

4.1 Topological T-duality

Performing a T-duality along the fiber direction of the principal circle bundle N_j maps the gerbe $\mathcal{G}_{j,k}$ to the gerbe $\mathcal{G}_{k,j}$ over the nilmanifold N_k . That is, the T-duality interchanges the Chern class j of the principal circle bundle N_j with the Dixmier-Douady class k of the gerbe $\mathcal{G}_{j,k}$ over N_j .

A mathematical picture behind this *topological T-duality* has been given long ago in [24], and the example of the gerbes $\mathcal{G}_{j,k}$ was discussed in full detail. Let us briefly summarize the essential points important for our purposes.

Consider a circle bundle $\pi : P \rightarrow M$ and first Chern class $c_1(P) = [F]$ together with a gerbe \mathcal{G} on P with Dixmier-Douady class $dd(\mathcal{G}) = [H]$. Then we have a closed integral 2-form \hat{F} on M from the integral along the fiber,

$$\hat{F}(x) = \pi_*(H)(x) := \oint_{\pi^{-1}(x)} H \quad (4.1)$$

¹⁷More precisely, the symmetries of the gerbe over local patches, which is all that is needed for T-duality. On double overlaps, we also have gerbe isomorphisms parameterized by Čech 1-cochains.

and $[\hat{F}]$ is the first Chern class of a principal circle bundle $\hat{\pi} : \hat{P} \rightarrow M$. One can now argue that there is a 3-form \hat{H} on \hat{P} with

$$\hat{\pi}_*(\hat{H}) = c_1(P), \tag{4.2}$$

which defines the Dixmier-Douady class of a gerbe $\hat{\mathcal{G}}$ on \hat{P} . Altogether, we obtain the following fibration of spaces:

$$\begin{array}{ccccc}
 & & P \times_M \hat{P} & & \\
 & \swarrow \text{pr} & & \searrow \hat{\text{pr}} & \\
 \mathcal{G} & \longrightarrow & P & & \hat{P} \longleftarrow \hat{\mathcal{G}} \\
 & & \searrow \pi & & \swarrow \hat{\pi} \\
 & & M & &
 \end{array} \tag{4.3}$$

Here, $P \times_M \hat{P}$ is the *correspondence space*

$$P \times_M \hat{P} := \{(x, \hat{x}) \in P \times \hat{P} \mid p(x) = \hat{p}(\hat{x})\}. \tag{4.4}$$

The correspondence space can either be seen as a circle bundle over P with Chern class $\pi^*(c_1(\hat{P}))$ or as a circle bundle over \hat{P} with Chern class $\hat{\pi}^*(c_1(P))$. Given normalized connection 1-forms A and \hat{A} on P and \hat{P} , by which we mean $\pi_*(A) = \hat{\pi}_*(\hat{A}) = 1$, we find that

$$\hat{\text{pr}}^*\hat{H} - \text{pr}^*H = d(\text{pr}^*A \wedge \hat{\text{pr}}^*\hat{A}) =: d\mathcal{F}, \tag{4.5}$$

where \mathcal{F} is a *global* 2-form on $P \times_M P$. Note that the latter implies that the gerbe $\text{pr}^*\mathcal{G} \otimes \hat{\text{pr}}^*\hat{\mathcal{G}}$ is trivial. Using \mathcal{F} , the *T-duality* map was defined in [24] to be the homomorphism $\hat{\text{pr}}_* \circ e^{\mathcal{F}\wedge} \circ \text{pr}^*$ where $\hat{\text{pr}}_* : P \times_M \hat{P} \rightarrow \hat{P}$ is defined analogously to (4.1). It was shown that this map descends to an isomorphism of twisted cohomologies on (P, H) and (\hat{P}, \hat{H}) , respectively. One of the aims of the following sections is to reformulate this in our language in order to have a starting point to investigate T-duality in later works.

4.2 Symplectic pre-NQ-manifold on the correspondence space

Let us now return to our example of pairs of nilmanifolds. Here, topological T-duality interchanges the Dixmier-Douady class k of the gerbe $\mathcal{G}_{j,k}$ over N_j with the first Chern class j of the nilmanifold N_j . We thus specialize and extend the diagram (4.3) to

$$\begin{array}{ccccc}
 & & \mathcal{G}_K & \xleftrightarrow{\quad} & \mathcal{C}_K \\
 & & \swarrow & & \searrow \\
 & & \mathcal{E}_K = \mathcal{E}_{k,j} & & \\
 & & \downarrow & & \\
 & & K := N_j \times_{T^2} N_k & & \\
 & \swarrow \text{pr}_1 & & \searrow \text{pr}_2 & \\
 \mathcal{G}_{j,k} & \longrightarrow & N_j & & N_k \longleftarrow \mathcal{G}_{k,j} \\
 \uparrow \text{dashed} & & \downarrow \pi_j & & \downarrow \pi_k \\
 & & T^2 & &
 \end{array} \tag{4.6}$$

where we also added the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$ to the picture which describe the symmetries of the gerbes $\mathcal{G}_{j,k}$ and $\mathcal{G}_{k,j}$. Note that the latter two gerbes can be pulled back to K and tensored there to the gerbe

$$\mathcal{G}_K := \text{pr}_1^* \mathcal{G}_{j,k} \otimes \text{pr}_2^* \mathcal{G}_{k,j} . \tag{4.7}$$

Restricting this gerbe to $N_j \hookrightarrow K$ and $N_k \hookrightarrow K$ yields the gerbes $\mathcal{G}_{j,k}$ and $\mathcal{G}_{k,j}$, respectively. The gerbe \mathcal{G}_K is thus indeed a suitable correspondence object on K . Its symmetries are described by the Courant algebroid $\mathcal{C}_{\mathcal{G}_K}$ on K whose Ševera class is the Dixmier-Douady class of \mathcal{G}_K .

As explained in the introduction, our goal is to characterize the space \mathcal{E}_K , the appropriate correspondence object on K for Double Field Theory. Since \mathcal{E}_K captures the symmetries and the differential geometry of the Double Field Theory relevant for the T-duality between $\mathcal{G}_{j,k}$ and $\mathcal{G}_{k,j}$, we expect it to be a subspace of $\mathcal{C}_{\mathcal{G}_K}$.

This point of view is further supported by our general discussion of the local case in [10], which also suggests that \mathcal{E}_K should be a *symplectic pre-NQ-manifold of degree 2*. By this we mean a symplectic \mathbb{N} -graded manifold of degree 2 endowed with a vector field Q of degree 1, which satisfies $\mathcal{L}_Q \omega = 0$. In particular, a symplectic pre-NQ-manifold with $Q^2 = 0$ is a symplectic NQ-manifold. By an L_∞ -structure \mathbb{L} on a symplectic pre-NQ-manifold \mathcal{E} , we mean a subset of the algebra of functions $\mathbb{L} \subset \mathcal{C}^\infty(\mathcal{E})$, such that the usual construction of the associated Lie n -algebra on \mathcal{E} closes and forms an L_∞ -algebra. Translated to the physics nomenclature, the algebraic condition of having an L_∞ -structure amounts essentially to (a slight weakening) of the strong section condition and an explicit L_∞ -structure is a solution to the strong section condition.

From our explicit description of N_j , we readily derive an equally explicit description of the correspondence space $K := N_j \times_{T^2} N_k$. The latter is a four-dimensional manifold with coordinates x^μ , $\mu = 1, \dots, 4$, subject to the identification

$$\begin{aligned} (x^1, x^2, x^3, x^4) &\sim (x^1, x^2 + 1, x^3, x^4) \sim (x^1, x^2, x^3 + 1, x^4) \\ &\sim (x^1, x^2, x^3, x^4 + 1) \sim (x^1 + 1, x^2, x^3 - jx^2, x^4 - kx^2) . \end{aligned} \tag{4.8}$$

We now extend K to the Courant algebroid $\mathcal{C}_{\mathcal{G}_K}$ over this space for the gerbe \mathcal{G}_K , which is also straightforward, as it combines the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$. We have coordinates $(x^\mu, \xi^\mu, \zeta_\mu, p_\mu)$, $\mu = 1, \dots, 4$ subject to the identification

$$\begin{aligned} (x^1, \xi^1, \xi^2, \xi^3, \xi^4, \zeta_1, \zeta_2, \zeta_3, \zeta_4, p_1, p_2, p_3, p_4) \\ \sim (x^1 + 1, \xi^1, \xi^2, \xi^3, \xi^4, \zeta_1, \zeta_2, \zeta_3, \zeta_4, p_1, p_2 + jp_3 + kp_4, p_3, p_4) . \end{aligned} \tag{4.9}$$

The symplectic form ω and the Hamiltonian Q of the homological vector field are then

$$\omega = dx^\mu \wedge dp_\mu + d\xi^\mu \wedge d\zeta_\mu, \quad Q = \xi^\mu p_\mu - jx^1 \xi^2 p_3 - kx^1 \xi^2 p_4 + k \xi^1 \xi^2 \xi^3 + j \xi^1 \xi^2 \xi^4 . \tag{4.10}$$

We note the expected symmetry under the exchange of j and k .

To construct the pre-NQ-manifold \mathcal{E}_K , we now linearly combine the fiber coordinates of degree 1 corresponding to the directions involved in the T-duality to new coordinates θ_\pm^α , $\alpha = 3, 4$:

$$\theta_\pm^3 := \frac{1}{2}(\xi^3 \pm \zeta_4), \quad \theta_\pm^4 := \frac{1}{2}(\xi^4 \pm \zeta_3), \tag{4.11}$$

closely following the construction in the local case as outlined in [10]. The pre-NQ-manifold \mathcal{E}_K in which we are interested is then obtained as the locus of $\theta_-^\alpha = 0$, $\alpha = 3, 4$. More formally, $\mathcal{C}^\infty(\mathcal{E}_K) = \mathcal{C}^\infty(\mathcal{C}_{\mathcal{G}_K})/\mathcal{I}$, where \mathcal{I} is the ideal generated by the θ_-^α .

Restricting the symplectic form and the Hamiltonian \mathcal{Q} to \mathcal{E}_K by pull-back along the implied embedding, we obtain

$$\begin{aligned}\omega &= dx^\mu \wedge dp_\mu + d\xi^1 \wedge d\zeta_1 + d\xi^2 \wedge d\zeta_2 + d\theta_+^3 \wedge d\theta_+^4 \\ &= dx^a \wedge dp_a + dx^\alpha \wedge dp_\alpha + d\xi^a \wedge d\zeta_a + \frac{1}{2}\eta_{\alpha\beta}d\theta_+^\alpha \wedge d\theta_+^\beta\end{aligned}\quad (4.12)$$

and

$$\mathcal{Q} = \xi^a p_a + \theta_+^\alpha p_\alpha - jx^1 \xi^2 p_3 - kx^1 \xi^2 p_4 + k\xi^1 \xi^2 \theta_+^3 + j\xi^1 \xi^2 \theta_+^4, \quad (4.13)$$

where we split all coordinates as $x^\mu = (x^a, x^\alpha)$, etc., $a = 1, 2$, $\alpha = 3, 4$ and used the $O(1, 1)$ -metric

$$\eta_{\alpha\beta} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}_{\alpha\beta} \quad (4.14)$$

with inverse $\eta^{\alpha\beta} = \eta_{\alpha\beta}$. The symplectic form ω induces the Poisson bracket

$$\begin{aligned}\{f, g\} &:= \left(\frac{\partial}{\partial p_\mu} f\right) \left(\frac{\partial}{\partial x^\mu} g\right) - \left(\frac{\partial}{\partial x^\mu} f\right) \left(\frac{\partial}{\partial p_\mu} g\right) \\ &\quad - (-1)^{|f|} \left(\frac{\partial}{\partial \zeta_a} f\right) \left(\frac{\partial}{\partial \xi^a} g\right) - (-1)^{|f|} \left(\frac{\partial}{\partial \xi^a} f\right) \left(\frac{\partial}{\partial \zeta_a} g\right) \\ &\quad - (-1)^{|f|} \left(\frac{\partial}{\partial \theta_+^4} f\right) \left(\frac{\partial}{\partial \theta_+^3} g\right) - (-1)^{|f|} \left(\frac{\partial}{\partial \theta_+^3} f\right) \left(\frac{\partial}{\partial \theta_+^4} g\right).\end{aligned}\quad (4.15)$$

We note that ω and \mathcal{Q} render \mathcal{E}_K a pre-NQ-manifold, since $\mathcal{L}_Q \omega = 0$ for the corresponding hamiltonian vector field, and

$$\{\mathcal{Q}, \mathcal{Q}\} = 2p_3 p_4 = \eta^{\alpha\beta} p_\alpha p_\beta. \quad (4.16)$$

This is precisely the relation also found in the local case discussed in [10]. This pre-NQ-manifold describes indeed the target space for double field theory manifestly invariant under the T-duality connecting the nilmanifolds N_j and N_k endowed with gerbes of Dixmier-Douady class k and j , respectively, as we shall see in the following.

We also note that the above constructions are readily extended to general pairs of Courant algebroids $(\mathcal{C}_{j,k}, \mathcal{C}_{m,n})$ between which there is no T-duality. The factorization of $\{\mathcal{Q}, \mathcal{Q}\}$ would here have read as

$$\{\mathcal{Q}, \mathcal{Q}\} = 2(p_3 + (k - m)\xi^1 \xi^2)(p_4 + (n - j)\xi^1 \xi^2). \quad (4.17)$$

We see that this factorization simplifies for T-dual pairs.

4.3 Symmetries, L_∞ -structures and the section condition

The symmetries of Double Field Theory on $K = N_j \times_{T^2} N_k$ are captured by an L_∞ -structure on \mathcal{E}_K as defined in [10]. Recall that an L_∞ -structure on \mathcal{E}_K is a subset \mathbf{L}^{loc} of $\mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K)$ (we restrict immediately to local functions),

$$\mathbf{L}^{\text{loc}} = \mathbf{L}_0^{\text{loc}} \oplus \mathbf{L}_1^{\text{loc}} \subset \mathcal{C}_{\text{loc},1}^\infty(\mathcal{E}_K) \oplus \mathcal{C}_{\text{loc},0}^\infty(\mathcal{E}_K) \quad (4.18a)$$

which becomes a Lie 2-algebra with the following higher brackets:

$$\begin{aligned}
\mu_1(f) &= \{\mathcal{Q}, f\} = \xi^\mu \tau_\mu^\nu \partial_\nu = d^\tau f, \\
\mu_2(X, Y) &= \frac{1}{2}(\{\{\mathcal{Q}, X\}, Y\} - \{\{\mathcal{Q}, Y\}, X\}) = [X, Y]^\tau + \frac{1}{2}(Y^A d^\tau X_A - X^A d^\tau Y_A), \\
\mu_2(X, f) &= \frac{1}{2}\{\{\mathcal{Q}, X\}, f\} = \frac{1}{2}\{X, \{\mathcal{Q}, f\}\} = \frac{1}{2}X^\mu \tau_\mu^\nu \partial_\nu f = \frac{1}{2}\mathcal{L}_X^\tau f, \\
\mu_3(X, Y, Z) &= \frac{1}{3!}(\{\{\{\mathcal{Q}, X\}, Y\}, Z\} \pm \text{perm.}) \\
&= Z_A X^\mu \tau_\mu^\nu \partial_\nu Y^A + X_A Y^\mu \tau_\mu^\nu \partial_\nu Z^A + Y_A Z^\mu \tau_\mu^\nu \partial_\nu X^A.
\end{aligned} \tag{4.18b}$$

Here, $X, Y, Z \in \mathbb{L}_0^{\text{loc}} \subset \mathcal{C}_{\text{loc},1}^\infty(\mathcal{E}_K)$ and $f \in \mathbb{L}_1^{\text{loc}} \subset \mathcal{C}_{\text{loc},0}^\infty(\mathcal{E}_K)$. We also used the index notation

$$x^\mu = (x^a, x^\alpha), \quad x_\alpha = \eta_{\alpha\beta} x^\beta, \quad \eta_{\alpha\beta} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \tag{4.19a}$$

$$(\xi^a, \zeta_a, \theta_+^\alpha) = (\xi^a, \zeta_\mu) = (\xi^\mu, \zeta_a) = \Xi^A$$

so that

$$X_a \xi^a + X^a \zeta_a + X_\alpha \theta_+^\alpha = X_a \xi^a + X^\mu \zeta_\mu = X_\mu \xi^\mu + X^a \zeta_a = X_A \Xi^A \tag{4.19b}$$

and

$$X_A = \eta_{AB} X^B, \quad \eta = \begin{pmatrix} 0 & \mathbb{1}_2 & 0 & 0 \\ \mathbb{1}_2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}. \tag{4.19c}$$

The vielbein τ_μ^ν is the four-dimensional generalisation of the vielbein τ_i^j , cf. (2.41):

$$\tau_\mu^\nu = \delta_\mu^\nu - j x^1 \delta_\mu^2 \delta_3^\nu - k x^1 \delta_\mu^2 \delta_4^\nu, \tag{4.19d}$$

and gives rise to the modified Lie and total derivatives,

$$[X, Y]^\tau = X^\mu \tau_\mu^\nu \partial_\nu Y - Y^\mu \tau_\mu^\nu \partial_\nu X \quad \text{and} \quad d^\tau = \xi^\mu \tau_\mu^\nu \partial_\nu. \tag{4.19e}$$

As shown in Theorem 4.7 of [10], the conditions for (4.18) to be an L_∞ -structure are that the brackets and the Poisson bracket close on \mathbb{L}^{loc} and that further

$$\begin{aligned}
\{Q^2 f, g\} + \{Q^2 g, f\} &= 0, \\
\{Q^2 X, f\} + \{Q^2 f, X\} &= 0, \\
\{\{Q^2 X, Y\}, Z\}_{[X,Y,Z]} &= 0
\end{aligned} \tag{4.20}$$

for all $f, g \in \mathbb{L}_1^{\text{loc}}$ and $X, Y, Z \in \mathbb{L}_0^{\text{loc}}$.

Explicitly, the elements of $\mathbb{L}_0^{\text{loc}}$ parameterize infinitesimal diffeomorphisms on K as well as the gauge transformations of the B -field, which is part of the connective structure of the gerbe \mathcal{G}_K . The elements of $\mathbb{L}_1^{\text{loc}}$ are higher gauge transformations, that is gauge transformations between gauge transformations which encode the redundancy in the gauge transformations parameterized by $\mathbb{L}_0^{\text{loc}}$.

The algebraic conditions (4.20) for having an L_∞ -structure translate to the section condition in the usual Double Field Theory nomenclature. They are necessary for having an appropriate symmetry algebra on the correspondence space. In [10], we showed that in the local case, they correspond in fact to a weakening of the section condition.

4.4 L_∞ -algebra structures: solutions to the section condition

Let us now come to discussing the solutions to the section condition which reduce \mathcal{E}_K to the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$. The factorization $\{\mathcal{Q}, \mathcal{Q}\} = 2p_3p_4$ gives a clear hint that we want to quotient by the Hamiltonian vector fields of p_3 and p_4 ,

$$\frac{\partial}{\partial x^3} = \{p_3, -\} \quad \text{and} \quad \frac{\partial}{\partial x^4} = \{p_4, -\}. \quad (4.21)$$

We want to regard these vector fields as part of integrable distributions which encode foliations whose leaves are the fibers to the projections

$$\mathcal{E}_K \xrightarrow{\hat{p}r_1} \mathcal{C}_{j,k} \quad \text{and} \quad \mathcal{E}_K \xrightarrow{\hat{p}r_2} \mathcal{C}_{k,j}. \quad (4.22)$$

For this, we clearly need an additional pair of vector fields V_3 and V_4 completing the distributions. It is natural to choose the symplectic dual vector fields to the global 1-forms

$$dx^3 + jx^1 dx^2 \quad \text{and} \quad dx^4 + kx^1 dx^2. \quad (4.23)$$

Up to a sign, we then arrive at

$$V_3 = \frac{\partial}{\partial p_3} + jx^1 \frac{\partial}{\partial p_2} \quad \text{and} \quad V_4 = \frac{\partial}{\partial p_4} + kx^1 \frac{\partial}{\partial p_2}, \quad (4.24)$$

which satisfy

$$\iota_{V_3}\omega = -(dx^3 + jx^1 dx^2) \quad \text{and} \quad \iota_{V_4}\omega = -(dx^4 + kx^1 dx^2) \quad (4.25)$$

and are therefore global.¹⁸ We note that both (∂_3, V_3) and (∂_4, V_4) span integrable distributions, giving indeed rise to the projections $\hat{p}r_2$ and $\hat{p}r_1$, respectively. Explicitly, we have the obvious projections

$$\begin{aligned} \hat{p}r_1(x^1, x^2, x^3, x^4, \xi^1, \xi^2, \zeta_1, \zeta_2, \theta_+^3, \theta_+^4, p_1, p_2, p_3, p_4) \\ = (\tilde{x}^1, \tilde{x}^2, \tilde{x}^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3) \\ = (x^1, x^2, x^3, \xi^1, \xi^2, \theta_+^3, \zeta_1, \zeta_2, \theta_+^4, p_1, p_2 - kx^1 p_4, p_3), \\ \hat{p}r_2(x^1, x^2, x^3, x^4, \xi^1, \xi^2, \zeta_1, \zeta_2, \theta_+^3, \theta_+^4, p_1, p_2, p_3, p_4) \\ = (\tilde{x}^1, \tilde{x}^2, \tilde{x}^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3) \\ = (x^1, x^2, x^4, \xi^1, \xi^2, \theta_+^4, \zeta_1, \zeta_2, \theta_+^3, p_1, p_2 - jx^1 p_3, p_4), \end{aligned} \quad (4.26)$$

where $(\tilde{x}^1, \dots, \tilde{p}_3)$ are the coordinates on the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$. Both projections allow for sections, which are the evident embeddings of $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$ into the pre-NQ-manifold \mathcal{E}_K :

$$\begin{aligned} e_1(\tilde{x}^1, \tilde{x}^2, \tilde{x}^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3) \\ = (x^1, x^2, x^3, x^4, \xi^1, \xi^2, \zeta_1, \zeta_2, \theta_+^3, \theta_+^4, p_1, p_2, p_3, p_4) \\ = (\tilde{x}^1, \tilde{x}^2, \tilde{x}^3, -k\tilde{x}^1\tilde{x}^2, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\xi}^3, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3, 0), \\ e_2(\tilde{x}^1, \tilde{x}^2, \tilde{x}^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\xi}^3, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{p}_1, \tilde{p}_2, \tilde{p}_3) \\ = (x^1, x^2, x^3, x^4, \xi^1, \xi^2, \zeta_1, \zeta_2, \theta_+^3, \theta_+^4, p_1, p_2, p_3, p_4) \\ = (\tilde{x}^1, \tilde{x}^2, -j\tilde{x}^1\tilde{x}^2, \tilde{x}^3, \tilde{\xi}^1, \tilde{\xi}^2, \tilde{\zeta}_1, \tilde{\zeta}_2, \tilde{\zeta}_3, \tilde{\xi}^3, \tilde{p}_1, \tilde{p}_2, 0, \tilde{p}_3). \end{aligned} \quad (4.27)$$

¹⁸The symplectic form is global and invertible and so are the two 1-forms.

These embeddings are morphisms of symplectic pre-NQ-manifolds: the symplectic forms and the Hamiltonians of the homological vector fields on the Courant algebroids are pull-backs of those on \mathcal{E}_K along e_1 and e_2 .¹⁹

It is now easy to show that the subsets

$$\mathbb{L}_{\hat{p}r_1} = \{F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K) \mid \partial_4 F = V_4 F = 0\} \quad \text{and} \quad \mathbb{L}_{\hat{p}r_2} = \{F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K) \mid \partial_3 F = V_3 F = 0\} \quad (4.28)$$

indeed form \mathbb{L}_∞ -algebra structures, verifying conditions (4.20). For this, we rewrite

$$Q^2 F = \{\mathcal{Q}, \{\mathcal{Q}, F\}\} = \frac{1}{2} \{\{\mathcal{Q}, \mathcal{Q}\}, F\} \quad (4.29)$$

for $F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K)$. We also note that

$$\{\{\mathcal{Q}, \mathcal{Q}\}, F\} = 2\{p_3 p_4, F\} = \begin{cases} 2p_4 \{p_3, F\} & \text{for } F \in \mathbb{L}_{\hat{p}r_1}, \\ 2p_3 \{p_4, F\} & \text{for } F \in \mathbb{L}_{\hat{p}r_2}. \end{cases} \quad (4.30)$$

We now restrict our discussion to the case $\mathbb{L}_{\hat{p}r_1}$; the case $\mathbb{L}_{\hat{p}r_2}$ follows fully analogously. We have

$$\{Q^2 f, g\} = p_4 \{\{p_3, f\}, g\} = 0 \quad (4.31a)$$

for $f, g \in \mathbb{L}_{\hat{p}r_1,1}$, since $\{p_3, f\}$ and g are functions of degree 0 and the Poisson bracket on \mathcal{E}_K between two such functions vanishes. Similarly,

$$\begin{aligned} \{Q^2 X, f\} &= p_4 \{\{p_3, X\}, f\} = 0, \\ \{Q^2 f, X\} &= p_4 \{\{p_3, f\}, X\} = 0 \end{aligned} \quad (4.31b)$$

for $f \in \mathbb{L}_{\hat{p}r_1,1}$ and $X \in \mathbb{L}_{\hat{p}r_1,0}$ since the Poisson bracket also vanishes between a function of degree 0 and a function of degree 1. Finally, we have

$$\{\{Q^2 X, Y\}, Z\} = p_4 \{\{\{p_3, X\}, Y\}, Z\} = 0 \quad (4.31c)$$

for $X, Y, Z \in \mathbb{L}_{\hat{p}r_1,0}$, since $\{\{p_3, X\}, Y\}$ is a function of degree 0 on \mathcal{E}_K . Altogether, (4.31) are sufficient for equations (4.20) to be satisfied. Note, however, that they are by no means necessary. They are rather special, in that each term in the symmetrizations contained in (4.20) vanishes separately.

A few remarks are in order. First, we have not used $V_4 F = 0$ for $F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K)$ at any point of the proof. Thus the subsets $\{F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K) \mid \partial_4 F = V_4 F = 0\}$ and $\{F \in \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K) \mid \partial_3 F = V_3 F = 0\}$ determining the projected Courant algebroids also form $\mathbb{L}_\infty^{\text{loc}}$ algebras, as it is to be expected from the fact that they represent Courant algebroids. The condition $V_4 F = 0$ seems to be merely necessary for reducing to a symplectic space in the projection $\hat{p}r_1$ of (4.22) and analogously for $V_3 F = 0$. Second, the above discussion straightforwardly extends to general pairs of Courant algebroids $(\mathcal{E}_{j,k}, \mathcal{E}_{m,n})$ between which there is no T-duality. The factorization (4.17) of $\{\mathcal{Q}, \mathcal{Q}\}$ yields two different vector fields,

$$D_3 = \{p_3 + (k - m)\xi^1 \xi^2, -\} \quad \text{and} \quad D_4 = \{p_4 + (n - j)\xi^1 \xi^2, -\}. \quad (4.32)$$

The projection map is correspondingly more complicated.

¹⁹Mathematically, we can thus encode T-duality as a cospan of pre-NQ-manifolds, but we refrain from pushing this thought any further in the current paper.

4.5 T-duality between $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$

We finally want to investigate the meaning of T-duality between (N_j, H) and (N_k, \hat{H}) in our language. Recalling diagram (4.6), we first rewrite the condition (4.5) in our language. Noting that the differential d in (4.5) is the untwisted de Rham differential on the correspondence space K , we split the homological function (4.13) according the untwisted part plus the flux part:

$$\mathcal{Q} = \mathcal{Q}_0 + k\xi^1\xi^2\theta_+^3 + j\xi^1\xi^2\theta_+^4 . \tag{4.33}$$

We further note that the product of the connection forms, i.e. $\hat{\text{pr}}^*A \wedge \hat{\text{pr}}^*\hat{A}$, lifted to \mathcal{E}_K takes the simple form $\mathcal{F} = -\theta_+^3\theta_+^4$. Hence, equation (4.5) takes the form

$$\{\mathcal{Q}_0, \mathcal{F}\} = -k\xi^1\xi^2\theta_+^3 + j\xi^1\xi^2\theta_+^4 . \tag{4.34}$$

Using this condition, it turns out that we are able to transport the homological function on $\mathcal{C}_{j,k}$ to its corresponding object on $\mathcal{C}_{k,j}$ using the projections $\hat{\text{pr}}_i$ of (4.26) and embeddings in (4.27) in composition:

$$\mathcal{Q}_{k,j} = e_2^* \circ \hat{\text{pr}}_1^* \mathcal{Q}_{j,k} , \tag{4.35}$$

where $\mathcal{Q}_{j,k}$ and $\mathcal{Q}_{k,j}$ are the homological functions on $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$, respectively. Evaluating the left hand side of (4.34) gives the relation for the twists:

$$\theta_+^3(p_3 + k\xi^1\xi^2) = \theta_+^4(p_4 + j\xi^1\xi^2) . \tag{4.36}$$

This enables the exchange of the parts differing in the pullbacks of the homological functions:

$$\begin{aligned} \hat{\text{pr}}_1^* \mathcal{Q}_{j,k} &= \xi^1 p_1 + \xi^2 p_2 + \theta_+^3 p_3 - \xi^2 x^1 (k p_4 + j p_3) + k \xi^1 \xi^2 \theta_+^3 \\ &= \xi^1 p_1 + \xi^2 p_2 + \theta_+^4 p_4 - \xi^2 x^1 (k p_4 + j p_3) + j \xi^1 \xi^2 \theta_+^4 = \hat{\text{pr}}_2^* \mathcal{Q}_{k,j} . \end{aligned}$$

It follows from the definition of e_2 that $e_2^* \circ \hat{\text{pr}}_1^* \mathcal{Q}_{j,k} = \mathcal{Q}_{k,j}$. We also have $e_2^* \circ \hat{\text{pr}}_1^* H = k \tilde{\xi}^1 \tilde{\xi}^2 \tilde{\zeta}_3$, which is the “ f -flux” T-dual to the 3-form flux H .

We see that there are two steps for T-duality here: first, the factorization of (4.16) determines projections $\hat{\text{pr}}_i$ from the pre- NQ -manifold \mathcal{E}_K to the respective Courant algebroids. And second, the condition (4.5) allows to relate the homological functions on the NQ -manifolds corresponding to the projected Courant algebroids. This is the reformulation of the T-duality map in our language. Of course we restricted to the very example of the two nilmanifolds with interchanged Chern- and Dixmier-Douady classes, but we expect this to hold true whenever geometric T-duality is possible.

4.6 Differential geometry on \mathcal{E}_K

Let us briefly comment also on the differential geometry on \mathcal{E}_K . The very general framework developed in [10] is readily applied to the case at hand.

The pre- NQ -manifold \mathcal{E}_K by itself describes the symmetries of Double Field Theory, which act on extended tensors. For these, we need to extend the algebra of functions on

\mathcal{E}_K to the free tensor algebra $T(\mathcal{E}_K)$ generated by the local coordinate functions over the ring $\mathcal{C}_{\text{loc},0}^\infty(\mathcal{E}_K)$. The Poisson bracket now uniquely extends to a map

$$\{-, -\} : \mathcal{C}_{\text{loc}}^\infty(\mathcal{E}_K) \times T(\mathcal{E}_K) \rightarrow T(\mathcal{E}_K), \quad (4.37)$$

which encodes in particular the action of elements of an L_∞ -structure \mathbb{L}^{loc} on \mathcal{E}_K on a subset $\mathbb{T}(\mathcal{E}_K) \subset T(\mathcal{E}_K)$. In order for this action to be a reasonable action of a 2-term L_∞ -algebra on a vector space, we require that

$$\{\{Q^2X, Y\} - \{Q^2Y, X\}, t\} = 0 \quad \text{and} \quad \{\{\{Q^2X, Y\} - \{Q^2Y, X\}, QZ\}, t\} = 0 \quad (4.38)$$

for all $t \in \mathbb{T}(\mathcal{E}_K)$ as well as all $X, Y, Z \in \mathbb{L}_0^{\text{loc}}$, cf. [10, Theorem 4.11]. This condition is the analogue for tensors of the condition (4.20) for the parameters of infinitesimal symmetries. We note that both conditions are trivially satisfied if we consider functions F (and the obvious extension to elements of the tensor algebra) with $\partial_4 F = V_4 F = 0$ or $\partial_3 F = V_3 F = 0$. So the above extended tensors on \mathcal{E}_K reduce to the extended tensors on the Courant algebroids $\mathcal{C}_{j,k}$ and $\mathcal{C}_{k,j}$.

Our formalism [10] also knows extended notions of covariant derivative, torsion and Riemann tensor. These merely rely on the available pre-NQ-structures. An extended covariant derivative ∇ on \mathcal{E}_K is defined as a linear map from $\mathcal{X}(\mathcal{E}_K)$ to $\mathcal{C}^\infty(\mathcal{E}_K)$ such that the image ∇_X for $X \in \mathcal{X}(\mathcal{E}_K)$ gives rise to a map

$$\{\nabla_X, -\} : \mathcal{X}(\mathcal{E}_K) \rightarrow \mathcal{X}(\mathcal{E}_K), \quad (4.39)$$

which readily generalizes to extended tensors and satisfies

$$\{\nabla_{fX}, Y\} = f\{\nabla_X, Y\} \quad \text{and} \quad \{\nabla_X, fY\} = \{QX, f\}Y + f\{\nabla_X, Y\} \quad (4.40)$$

for all $f \in \mathcal{C}_0^\infty(\mathcal{E}_K)$ and extended tensors $Y \in \mathbb{T}(\mathcal{E}_K)$. For this notion of extended covariant derivative, we readily write down the extended torsion and extended Riemann tensors,

$$\begin{aligned} \mathcal{T} : \otimes^3 \mathcal{X}(\mathcal{E}_K) &\rightarrow \mathcal{C}_0^\infty(\mathcal{E}_K), \\ \mathcal{T}(X, Y, Z) &:= 3\left(-\left\{X, \{\nabla_Y, Z\}\right\}\right)_{[X, Y, Z]} + \frac{1}{2}(\{X, \{QZ, Y\}\} - \{Z, \{QX, Y\}\}), \\ \mathcal{R} : \otimes^4 \mathcal{X}(\mathcal{E}_K) &\rightarrow \mathcal{C}_0^\infty(\mathcal{E}_K), \\ \mathcal{R}(X, Y, Z, W) &:= \frac{1}{2}\left(\left\{\{\nabla_X, \nabla_Y\} - \nabla_{\mu_2(X, Y)}, Z\right\}, W\right) + (Z \leftrightarrow W) \\ &\quad + \left\{\{\nabla_Z, \nabla_W\} - \nabla_{\{\nabla_Z, W\} - \{\nabla_W, Z\}}, X\right\}, Y\right) + (X \leftrightarrow Y), \end{aligned} \quad (4.41)$$

where $X, Y, Z, W \in \mathcal{X}(\mathcal{E}_K)$, the subset of $\mathbb{T}(\mathcal{E}_K)$ of degree 1. Note that both \mathcal{T} and \mathcal{R} are indeed tensorial up to terms which vanish on L_∞ -structures.

5 Finite symmetries of Double Field Theory

We shall first discuss the finite symmetries of local Double Field Theory in a general context before specializing to the case of the symplectic pre-NQ-manifold \mathcal{E}_K .

5.1 Generalized Lie derivative and infinitesimal symmetries

We start from the description of local Double Field Theory as given in [10]. That is, we consider the symplectic pre-NQ-manifold $\mathcal{E}_2(\mathbb{R}^d)$, which has local coordinates x^M, θ^M, p_M of degrees 0, 1 and 2 and index $M = 1, \dots, 2d$. The symplectic form and the Hamiltonian \mathcal{Q} read as

$$\omega = dx^M \wedge dp_M + \frac{1}{2} \eta_{MN} d\theta^M \wedge d\theta^N \quad \text{and} \quad \mathcal{Q} = \theta^M p_M, \quad (5.1)$$

where η_{MN} is the usual $O(d, d)$ -structure. The Lie 2-algebra structure is given by the same derived bracket as in (4.18b), and L_∞ -structures have to satisfy the constraint (or section condition) (4.20).

The generalized Lie derivative of Double Field Theory is simply the action of the extended vector fields in the Lie 2-algebra of symmetries on extended vector fields and extended tensors. On extended vector fields, it is given by the D-bracket ν_2 , which is the following derived bracket:

$$\begin{aligned} \hat{\mathcal{L}}_X Y &= \nu_2(X, Y) = \{ \{ \mathcal{Q}, X \}, Y \} = \left\{ X^M p_M + \theta^M \theta^N \frac{\partial}{\partial x^M} X_N, Y \right\} \\ &= [X, Y] + \theta^M Y^K \frac{\partial}{\partial x^M} X_K \end{aligned} \quad (5.2)$$

for generalized vectors $X = \theta^M X_M, Y = \theta^M Y_M$ in $\mathbb{L}_0^{\text{loc}} \subset \mathcal{C}_1^\infty(\mathcal{E}_2)$. We now want to impose two further restrictions on our L_∞ -structure \mathbb{L}^{loc} and set of extended tensors \mathbb{T} . First of all, we want that symmetries between symmetries have no effect on the action. This implies that

$$\nu_2(\mu_1(f), Y) = \{ \{ \mathcal{Q}, \{ \mathcal{Q}, f \} \}, Y \} = 0 \quad (5.3)$$

for all $f \in \mathbb{L}_1^{\text{loc}} \subset \mathcal{C}_0^\infty(\mathcal{E}_K)$ and $Y \in \mathbb{T}$. Note that

$$\{ \{ \mathcal{Q}, \{ \mathcal{Q}, f \} \}, Y \} = \{ \{ \frac{1}{2} \{ \mathcal{Q}, \mathcal{Q} \}, f \}, Y \} = \frac{1}{2} \{ \{ (p_M p^M, f), Y \} \} = \left(\frac{\partial}{\partial x^M} f \right) \left(\frac{\partial}{\partial x_M} Y \right), \quad (5.4)$$

which is the usual section condition on $\mathcal{E}_2(M)$, trivially satisfied, for example, for the trivial L_∞ -structure

$$\mathbb{L}_+^{\text{loc}} = \left\{ F \in \mathcal{C}^\infty(\mathcal{E}_2(\mathbb{R}^d)) \mid \frac{\partial}{\partial x_\mu} F = \frac{\partial}{\partial p^\mu} F = 0 \right\}. \quad (5.5)$$

Second, we would like ν_2 to form a Leibniz algebra on the degree 0 elements $\mathbb{L}_0^{\text{loc}}$ of an L_∞ -structure \mathbb{L}^{loc} . As explained in [10], this also requires that

$$\{ \{ \{ \mathcal{Q}, \mathcal{Q} \}, X \}, Y \}, Z \} = 0 \quad (5.6)$$

for all $X, Y, Z \in \mathbb{L}_0^{\text{loc}}$. Explicitly,

$$\begin{aligned} \{ \{ \{ \mathcal{Q}, \mathcal{Q} \}, X \}, Y \}, Z \} &= \{ \{ \{ p_M p^M, X \}, Y \}, Z \} \\ &= \{ (\partial_M X) (\partial^M Y), Z \} + \{ \{ \partial_M X, Y \}, \partial^M Z \}, \end{aligned} \quad (5.7)$$

again an analogue of the section condition trivially satisfied for the L_∞ -structure $\mathbb{L}_+^{\text{loc}}$ introduced above.

Now the action of the generalized Lie derivative, together with the commutator, forms a Lie algebra, just as in the case of the Courant algebroid. The structure constants are encoded in the C-bracket, i.e. the higher product μ_2 ,

$$[\hat{\mathcal{L}}_{X_1}, \hat{\mathcal{L}}_{X_2}]Y = \hat{\mathcal{L}}_{X_1}\hat{\mathcal{L}}_{X_2}Y - \hat{\mathcal{L}}_{X_2}\hat{\mathcal{L}}_{X_1}Y = \hat{\mathcal{L}}_{\mu_2(X_1, X_2)}Y, \quad (5.8)$$

$X_{1,2} \in L_0^{\text{loc}} \subset \mathcal{C}_1^\infty(\mathcal{E}_2)$, which is easily verified using the derived bracket formulation of this equation. Since the action is Hamiltonian and degree preserving, we know that this has to be a Lie subalgebra of the degree preserving symplectomorphisms on $\mathcal{E}_2(\mathbb{R}^d)$ as again familiar from the case of the Courant algebroid. Contrary to that case, however, it is a little bit harder to identify this Lie subalgebra.

We note that the generalized Lie derivative acts as the ordinary Lie derivative on functions f of degree 0:

$$\hat{\mathcal{L}}_X f = \{\{\mathcal{Q}, X\}, f\} = \{X^M p_M + \theta^M \partial_M X, f\} = X^M \partial_M f = \mathcal{L}_X f, \quad (5.9)$$

and the Lie algebra of actions on these functions is simply the Lie algebra of actions of infinitesimal diffeomorphisms on the body of \mathcal{E}_2 , which is homomorphic to the Lie algebra of infinitesimal diffeomorphisms on the body of \mathcal{E}_2 .

On generalized vectors, which are functions Y of degree 1, however, we have an expression differing from \mathcal{L}_X :

$$\begin{aligned} \hat{\mathcal{L}}_X Y &= \{\{\mathcal{Q}, X\}, Y\} = \{X^M p_M + \theta^M \partial_M X, Y\} \\ &= X^M \partial_M Y - Y^M \partial_M X + \theta^M Y_K \partial_M X^K \\ &= [X, Y] + \theta^M Y_K \partial_M X^K. \end{aligned} \quad (5.10)$$

To find the Lie algebra of general actions, we consider the Leibniz algebra formed by ν_2 , which is fully determined by its action on the coordinate functions. We have

$$\begin{aligned} \hat{\mathcal{L}}_X x^M &= \{\{\mathcal{Q}, X\}, x^M\} = X^M, \\ \hat{\mathcal{L}}_X \theta^M &= \{\{\mathcal{Q}, X\}, \theta^M\} = \theta^N \frac{\partial}{\partial x^N} X^M - \frac{\partial}{\partial x^M} X, \\ \hat{\mathcal{L}}_X p_M &= \{\{\mathcal{Q}, X\}, p_M\} = -p_N \frac{\partial}{\partial x^M} X^N - \theta^N \frac{\partial}{\partial x^N} \frac{\partial}{\partial x^M} X. \end{aligned} \quad (5.11)$$

A similar action was discussed already in [13], where also a finite form was given to which we shall return later. The commutator of two actions of the generalized Lie derivative has an interesting form,

$$\begin{aligned} [\hat{\mathcal{L}}_{X_1}, \hat{\mathcal{L}}_{X_2}]F &= \{\{\mathcal{Q}, X_1\}, \{\{\mathcal{Q}, X_2\}, F\}\} - \{\{\mathcal{Q}, X_2\}, \{\{\mathcal{Q}, X_1\}, F\}\} \\ &= \{\{\{\mathcal{Q}, X_1\}, \{\mathcal{Q}, X_2\}\}, F\} = \{\{\mathcal{Q}, \{\{\mathcal{Q}, X_1\}, X_2\}\}, F\} \\ &= \{\{\mathcal{Q}, [X_1, X_2] + \theta^M X_2^K \partial_M X_{1K}\}, F\} \\ &= \{[X_1, X_2]^M p_M + \theta^M \partial_M [X_1, X_2] - \theta^M \theta^N \partial_N (X_2^K \partial_M X_{1K}), F\} \\ &= \mathcal{L}_{[X_1, X_2]}F + R(X_1, X_2)F, \end{aligned} \quad (5.12)$$

where F is an element of an L_∞ -structure and thus $|F| \leq 1$. Here, we used the variants of the section condition

$$\{\{\{\mathcal{Q}, \mathcal{Q}\}, X_1\}, X_2\}, F\} = 0 \quad \text{and} \quad \{p^M X_2^K \partial_M X_{1K}, F\} = 0, \quad (5.13)$$

as well as the definition of the ordinary Lie derivative $\mathcal{L}_X = d\iota_X + \iota_X d$ via the identification

$$dX = \theta^M \{p_M, X\} \quad \text{and} \quad \iota_{X_1} X_2 = \{X_1, X_2\}. \quad (5.14)$$

Note that $R(X_1, X_2)F$ is a contraction of F with a local and antisymmetric matrix. So we conclude that the action of the generalized Lie derivative forms the Lie algebra of diffeomorphisms whose action on extended tensors is modified by a local $\mathfrak{o}(d, d)$ -transformation. Note that the same conclusion was already reached in [18].

5.2 Finite symmetries

A more general form of (5.11) has already been integrated in [13] and the finite transformations are given by

$$\begin{aligned} x^M &\mapsto \tilde{x}^M = \tilde{x}^M(x), \\ \theta^M &\mapsto \tilde{\theta}^M = T^M_N(x)\theta^N, \\ p_M &\mapsto \tilde{p}_M = \frac{\partial x^N}{\partial \tilde{x}^M} p_N + \frac{1}{2}\theta^N \frac{\partial T^K_N(x)}{\partial \tilde{x}^M} \eta_{KL} T^L_P(x)\theta^P, \end{aligned} \quad (5.15)$$

where the first map is a diffeomorphism and $T^M_N(x)$ is a local $\mathfrak{O}(d, d)$ -transformation, because we are integrating an infinitesimal degree-preserving symplectomorphism and η_{MN} is part of the symplectic form, cf. (5.1).

From the infinitesimal case, we know that the choice of diffeomorphism, together with the metric η_{MN} , induces the local $\mathfrak{O}(d, d)$ -transformation. A clear description of the action of the generalized Lie derivative is found by recalling the interpretation of the Dorfman bracket from Generalized Geometry,

$$\begin{aligned} \hat{\mathcal{L}}_{X+\alpha}^0(Y + \beta) &:= \{\{\mathcal{Q}_0, X + \alpha\}, Y + \beta\} \\ &= [X, Y] + \mathcal{L}_X \beta - \iota_Y d\alpha. \end{aligned} \quad (5.16)$$

for vector fields X, Y and 1-forms α, β , where \mathcal{Q}_0 is the Hamiltonian of the homological vector field *without* any B -field corrections (i.e. the effects of the B -field are in the transition functions). The generalized Lie derivative gives an action of the semidirect product of the Lie algebra of diffeomorphisms and local 1-forms on generalized vectors in Generalized Geometry. There is, however, a second action within ordinary differential geometry, which is given by the first two terms in (5.16). In this picture, the last term describes the additional term in $e^B(Y + \beta)$ due to the gauge transformation

$$\begin{aligned} B &\mapsto B' = B + \delta B = B + d\alpha + \hat{\mathcal{L}}_{X+\alpha} B \\ &= B + d\alpha + \mathcal{L}_X B \\ &= B + \{\mathcal{Q}_0, \alpha\} + \{\{\mathcal{Q}_0, X + \alpha\}, B\}. \end{aligned} \quad (5.17)$$

Explicitly, we have

$$\begin{aligned}
 Y + \beta &\xrightarrow{\hat{\mathcal{L}}_X^0} Y' + \beta' = Y + \beta + \hat{\mathcal{L}}_{X+\alpha}(Y + \beta) \\
 &= Y + \beta + [X, Y] + \mathcal{L}_X \beta - \iota_Y d\alpha,
 \end{aligned}
 \tag{5.18}$$

but after a B -field coordinate change, $e^B(Y + \beta)$ is mapped to $e^{B'}(Y' + \beta')$ as follows:

$$\begin{aligned}
 e^{B'}(Y' + \beta') &= Y + \beta + \{\{\mathcal{Q}_0, X + \alpha\}, Y + \beta\} + \{Y + \beta + \{\{\mathcal{Q}_0, X + \alpha\}, Y + \beta\}, B\} + \\
 &\quad + \{Y + \beta, \{\mathcal{Q}_0, \alpha\}\} + \{\{\mathcal{Q}_0, X + \alpha\}, B\} \\
 &= Y + \beta + \{Y + \beta, B\} + \{\{\mathcal{Q}_0, X\}, Y + \beta\} + \{\{\mathcal{Q}_0, X + \alpha\}, \{Y + \beta, B\}\} \\
 &= (1 + \mathcal{L}_X)e^B(Y + \beta).
 \end{aligned}
 \tag{5.19}$$

Thus, $e^B(Y + \beta)$ transforms as an ordinary section of $TM \oplus T^*M$, namely with the action of diffeomorphisms.

In Double Field Theory, we have a similar split of contributions as in (5.16):

$$\hat{\mathcal{L}}_X^0 Y = X^M \partial_M Y - Y^M \partial_M X + \theta^M Y^K \partial_M X_K
 \tag{5.20}$$

for $X, Y \in \mathcal{C}_1^\infty(\mathcal{E}_2)$. Here, the first term describes the action of infinitesimal diffeomorphisms on the generalized vectors Y regarded as functions and the second and third term describe the terms of the form $e^{\delta B} Y$ for

$$\delta B = \frac{1}{2} \partial_M X_N \theta^M \theta^N.
 \tag{5.21}$$

We thus interpret the generalized Lie derivative as a transformation

$$Y \xrightarrow{\hat{\mathcal{L}}_X} Y' = Y + \mathcal{L}_X Y - \iota_Y dX,
 \tag{5.22}$$

which amounts to the diffeomorphism

$$e^B Y = Y \xrightarrow{\hat{\mathcal{L}}_X} e^{B'} Y' = Y + X^M \partial_M Y.
 \tag{5.23}$$

This transformation is readily checked in the case $B = 0$ and $B' = \delta B$, and it is the analogue of the above situation in Generalized Geometry.

Regarding B -transformed doubled vectors as scalars on the doubled space may be unusual, but our interpretation is supported by the following points. First, we saw that the action of the generalized Lie derivative on functions $f \in \mathcal{C}_0^\infty(\mathcal{E}_2)$ generates the diffeomorphisms group on the base space of \mathcal{E}_2 . Second, we also know that the action of the generalized Lie derivative on arbitrary functions on \mathcal{E}_2 induces a subgroup of the symplectomorphisms on \mathcal{E}_2 . The only remaining candidate is thus the group of diffeomorphisms. Third, this fits with the strong section condition, $\partial_M X \partial^M Y = 0$ on arbitrary tensors X, Y , which always seemed incompatible with covariance from a differential geometric point of view. Fourth, doubled vectors combine vector fields and 1-forms, which transform in opposite ways under diffeomorphisms. The scalar transformation law can be regarded as a consistent “compromise.”

The situation for an arbitrary initial Kalb-Ramond field B is more involved. First of all, e^B does not truncate,

$$\begin{aligned} e^B Y &= Y + \{Y, B\} + \frac{1}{2} \{\{Y, B\}, B\} + \dots, \\ &= \theta_M (Y^M + Y^N B_N^M + \frac{1}{2} Y^K B_K^N B_N^M + \dots) \\ &= \theta_M Y^N (\delta_N^M + B_N^M + \frac{1}{2} B_N^K B_K^M + \dots), \end{aligned} \tag{5.24}$$

but this expression converges: it is simply the usual matrix exponential multiplying Y . The problem is rather that in Double Field Theory, the analogue of the transformation e^B in Generalized Geometry which moves the B -field modification from the transition function to the homological vector field (and thus into the generalized Lie derivative) may now produce noncommutative or nonassociative spaces, which cannot be described without generalizing our framework.

5.3 Specialization to \mathcal{E}_K

Let us briefly specify the above discussion to \mathcal{E}_K ; the necessary adaptations are rather straightforward. The generalized Lie derivative reads as

$$\begin{aligned} \hat{\mathcal{L}}_X Y &= \{\{Q, X\}, Y\} = [X, Y]^\tau + Y_A d^\tau X^A + \\ &+ j((Y^1 X^2 - X^1 Y^2)\theta_+^4 + (Y^2 X_4 - X^2 Y_4)\xi^1 + (X^1 Y_4 - Y^1 X_4)\xi^2) \\ &+ k((Y^1 X^2 - X^1 Y^2)\theta_+^3 + k(Y^2 X_3 - X^2 Y_3)\xi^1 + k(X^1 Y_3 - Y^1 X_3)\xi^2), \end{aligned} \tag{5.25}$$

where we used again the notation introduced in (4.19) and $X, Y \in \mathcal{C}_{loc,1}^\infty(\mathcal{E}_K)$.

The relevant section conditions (5.3) and (5.6) translate to

$$\begin{aligned} \{\{Q, \{Q, f\}\}, Y\} &= \{\{\frac{1}{2}\{Q, Q\}, f\}, Y\} \\ &= \{\{2p_3 p_4, f\}, Y\} \\ &= (\partial_3 f)(\partial_4 Y) + (\partial_4 Y)(\partial_3 f) = (\partial_\alpha f)(\partial^\alpha Y) = 0, \end{aligned} \tag{5.26}$$

and

$$\begin{aligned} \{\{\{\{Q, Q\}, X\}, Y\}, Z\} &= \{\{\{p_3 p_4, X\}, Y\}, Z\} \\ &= \{(\partial_\alpha X)(\partial^\alpha Y), Z\} + \{\{\partial_\alpha X, Y\}, \partial^\alpha Z\} = 0. \end{aligned} \tag{5.27}$$

Both conditions are trivially satisfied for the L_∞ -structures $\mathbb{L}_{\hat{pr}_1}^{loc}$ and $\mathbb{L}_{\hat{pr}_2}^{loc}$ defined in (4.28).

The generalized Lie derivative now acts as the semi-direct product of diffeomorphisms on the correspondence space K and closed 2-forms along the directions x^1 and x^2 not involved in the T-duality.

6 Discussion

We would like to stress that nilmanifolds as base spaces for a discussion of global Double Field Theory were chosen merely for pedagogical reasons. If we have a T-duality relation between two spaces X_1 and X_2 fibered over a common base manifold M and carrying B -fields or, equivalently, gerbes \mathcal{G}_1 and \mathcal{G}_2 with connective structure, we always have a

correspondence space K carrying a gerbe \mathcal{G}_K which is the tensor product of the pullbacks of \mathcal{G}_1 and \mathcal{G}_2 . This gerbe comes with a Courant algebroid \mathcal{C}_K , which can be restricted to a pre-NQ-manifold \mathcal{E}_K capturing the global geometry of Double Field Theory:

$$\begin{array}{ccccc}
 & & \mathcal{G}_K & & \\
 & & \swarrow & \leftarrow & \mathcal{E}_K & \leftarrow & \mathcal{C}_K \\
 & & \downarrow & & \downarrow & & \downarrow \\
 \mathcal{G}_1 & \rightarrow & X_1 & \xrightarrow{\text{pr}_1} & K := X_1 \times_M X_2 & \xrightarrow{\text{pr}_2} & X_2 & \leftarrow & \mathcal{G}_2 \\
 \uparrow & & \uparrow & & \uparrow & & \uparrow & & \uparrow \\
 \mathcal{C}_1 & & \mathcal{G}_1 & & \mathcal{E}_K & & \mathcal{G}_2 & & \mathcal{C}_2 \\
 \downarrow & & \downarrow & & \downarrow & & \downarrow & & \downarrow \\
 & & \mathcal{G}_1 & & & & \mathcal{G}_2 & &
 \end{array}
 \tag{6.1}$$

In this way, any geometric T-duality²⁰ is readily captured in our framework as the special case in which the pull-backs of the Dixmier-Douady classes of the gerbes to K agree.

The extension to Exceptional Field Theory (EFT), however, is more problematic. The issue here is that a clear analogue of the local picture developed in [10] is still missing. Note that the analogues of Courant algebroids capturing the symmetries of the C -field are mostly clear. They are higher versions of Vinogradov algebroids containing shifted copies to accommodate both M2- and M5-brane windings, cf. [35, 36]. The situation over the correspondence space, however, is much less clear. In particular, we do not have (yet) the higher analogue of \mathcal{E}_K , not even in the local case.

Acknowledgments

This work has been partially supported by the European Cooperation in Science and Technology (COST) Action MP1405 QSPACE. The research of AD was supported by OP RDE project No. CZ.02.2.69/0.0/0.0/16_027/0008495, International Mobility of Researchers at Charles University.

Open Access. This article is distributed under the terms of the Creative Commons Attribution License (CC-BY 4.0), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.

References

[1] G. Aldazabal, D. Marques and C. Núñez, *Double field theory: A pedagogical review*, *Class. Quant. Grav.* **30** (2013) 163001 [[arXiv:1305.1907](#)] [[INSPIRE](#)].
 [2] D.S. Berman and D.C. Thompson, *Duality symmetric string and M-theory*, *Phys. Rept.* **566** (2014) 1 [[arXiv:1306.2643](#)] [[INSPIRE](#)].

²⁰Not involving noncommutative or nonassociative spaces.

- [3] O. Hohm, D. Lüst and B. Zwiebach, *The spacetime of double field theory: Review, remarks and outlook*, *Fortsch. Phys.* **61** (2013) 926 [[arXiv:1309.2977](#)] [[INSPIRE](#)].
- [4] M. Cederwall, *The geometry behind double geometry*, *JHEP* **09** (2014) 070 [[arXiv:1402.2513](#)] [[INSPIRE](#)].
- [5] M. Cederwall, *T-duality and non-geometric solutions from double geometry*, *Fortsch. Phys.* **62** (2014) 942 [[arXiv:1409.4463](#)] [[INSPIRE](#)].
- [6] R. Blumenhagen, F. Hassler and D. Lüst, *Double field theory on group manifolds*, *JHEP* **02** (2015) 001 [[arXiv:1410.6374](#)] [[INSPIRE](#)].
- [7] R. Blumenhagen, P. du Bosque, F. Hassler and D. Lüst, *Generalized metric formulation of double field theory on group manifolds*, *JHEP* **08** (2015) 056 [[arXiv:1502.02428](#)] [[INSPIRE](#)].
- [8] F. Hassler, *The topology of Double Field Theory*, *JHEP* **04** (2018) 128 [[arXiv:1611.07978](#)] [[INSPIRE](#)].
- [9] L. Freidel, F.J. Rudolph and D. Svoboda, *Generalised kinematics for double field theory*, *JHEP* **11** (2017) 175 [[arXiv:1706.07089](#)] [[INSPIRE](#)].
- [10] A. Deser and C. Sämann, *Extended Riemannian geometry I: Local double field theory*, [[arXiv:1611.02772](#)] [[INSPIRE](#)].
- [11] A. Deser, M.A. Heller and C. Sämann, *Extended Riemannian geometry II: Local heterotic double field theory*, *JHEP* **04** (2018) 106 [[arXiv:1711.03308](#)] [[INSPIRE](#)].
- [12] M. Gualtieri, *Generalized complex geometry*, Ph.D. Thesis, Oxford University, Oxford U.K. (2003) [[math/0401221](#)] [[INSPIRE](#)].
- [13] D. Roytenberg, *On the structure of graded symplectic supermanifolds and Courant algebroids*, in *Contemporary Mathematics. Vol. 315: Quantization, Poisson Brackets and Beyond*, T. Voronov eds., AMS Press, Providence U.S.A. (2002) [[math/0203110](#)] [[INSPIRE](#)].
- [14] A. Deser and J. Stasheff, *Even symplectic supermanifolds and double field theory*, *Commun. Math. Phys.* **339** (2015) 1003 [[arXiv:1406.3601](#)] [[INSPIRE](#)].
- [15] A. Deser and C. Sämann, *Derived brackets and symmetries in Generalized Geometry and Double Field Theory*, *PoS(CORFU2017)* 141 [[arXiv:1803.01659](#)] [[INSPIRE](#)].
- [16] D. Fiorenza and M. Manetti, *L_∞ structures on mapping cones*, *Algebra Number Theory* **1** (2007) 301 [[math/0601312](#)].
- [17] E. Getzler, *Higher derived brackets*, [[arXiv:1010.5859](#)].
- [18] D.S. Berman, M. Cederwall and M.J. Perry, *Global aspects of double geometry*, *JHEP* **09** (2014) 066 [[arXiv:1401.1311](#)] [[INSPIRE](#)].
- [19] O. Hohm and B. Zwiebach, *Large gauge transformations in double field theory*, *JHEP* **02** (2013) 075 [[arXiv:1207.4198](#)] [[INSPIRE](#)].
- [20] C.M. Hull, *Finite gauge transformations and geometry in double field theory*, *JHEP* **04** (2015) 109 [[arXiv:1406.7794](#)] [[INSPIRE](#)].
- [21] J.-H. Park, *Comments on double field theory and diffeomorphisms*, *JHEP* **06** (2013) 098 [[arXiv:1304.5946](#)] [[INSPIRE](#)].
- [22] G. Papadopoulos, *Seeking the balance: Patching double and exceptional field theories*, *JHEP* **10** (2014) 089 [[arXiv:1402.2586](#)] [[INSPIRE](#)].

- [23] M.A. Heller, N. Ikeda and S. Watamura, *Unified picture of non-geometric fluxes and T-duality in double field theory via graded symplectic manifolds*, *JHEP* **02** (2017) 078 [[arXiv:1611.08346](#)] [[INSPIRE](#)].
- [24] P. Bouwknegt, J. Evslin and V. Mathai, *T duality: Topology change from H flux*, *Commun. Math. Phys.* **249** (2004) 383 [[hep-th/0306062](#)] [[INSPIRE](#)].
- [25] G.R. Cavalcanti and M. Gualtieri, *Generalized complex geometry and T-duality*, in *A Celebration of the Mathematical Legacy of Raoul Bott (CRM Proceedings & Lecture Notes)*, AMS Press, Providence U.S.A. (2010), pg. 341 [[arXiv:1106.1747](#)] [[INSPIRE](#)].
- [26] O. Hohm and B. Zwiebach, *L_∞ algebras and field theory*, *Fortsch. Phys.* **65** (2017) 1700014 [[arXiv:1701.08824](#)] [[INSPIRE](#)].
- [27] B. Jurčo, L. Raspollini, C. Sämann and M. Wolf, *L_∞ -algebras of classical field theories and the Batalin-Vilkovisky formalism*, [arXiv:1809.09899](#) [[INSPIRE](#)].
- [28] C.M. Hull and R.A. Reid-Edwards, *Non-geometric backgrounds, doubled geometry and generalised T-duality*, *JHEP* **09** (2009) 014 [[arXiv:0902.4032](#)] [[INSPIRE](#)].
- [29] N.J. Hitchin, *Lectures on special Lagrangian submanifolds*, *AMS/IP Stud. Adv. Math.* **23** (2001) 151 [[math/9907034](#)] [[INSPIRE](#)].
- [30] S. Bunk, C. Sämann and R.J. Szabo, *The 2-Hilbert space of a prequantum bundle gerbe*, *Rev. Math. Phys.* **30** (2017) 1850001 [[arXiv:1608.08455](#)] [[INSPIRE](#)].
- [31] P. Bressler and A. Chervov, *Courant algebroids*, *J. Math. Sci.* **128** (2005) 3030 [[hep-th/0212195](#)] [[INSPIRE](#)].
- [32] B.L. Collier, *Infinitesimal symmetries of Dixmier-Douady gerbes*, [arXiv:1108.1525](#).
- [33] N. Hitchin, *Brackets, forms and invariant functionals*, [math/0508618](#) [[INSPIRE](#)].
- [34] A. Chatzistavrakidis, L. Jonke and O. Lechtenfeld, *Dirac structures on nilmanifolds and coexistence of fluxes*, *Nucl. Phys. B* **883** (2014) 59 [[arXiv:1311.4878](#)] [[INSPIRE](#)].
- [35] D. Baraglia, *Leibniz algebroids, twistings and exceptional generalized geometry*, *J. Geom. Phys.* **62** (2012) 903 [[arXiv:1101.0856](#)] [[INSPIRE](#)].
- [36] A.S. Arvanitakis, *Brane Wess-Zumino terms from AKSZ and exceptional generalised geometry as an L_∞ -algebroid*, [arXiv:1804.07303](#) [[INSPIRE](#)].